

Lecture 2 - Uniform Excitation - Outline

- **Announcements**
Handouts - Lecture Outline, Review, and Summary
- **Review** (of Lecture 1 and Recitation 2)
Carrier concentrations in TE given the doping level
What happens above and below room temperature?
- **Uniform excitation of uniform samples: drift** (introduced in Rec. 2)
Drift motion: carrier velocity versus field
Mobility, Drift currents, Conductivity
Impact of temperature on mobility, conductivity
Integrated circuit resistors
- **Uniform excitation: optical generation**
Generation/recombination in TE
Uniform optical generation - external excitation
Population excesses, p' and n' , and their transients
Low level injection; minority carrier lifetime
- **Uniform excitation - applied field and optical generation**
Photoconductivity, photoconductors

Variation of carrier concentration with temperature

(Note: for convenience we assume an n-type sample)

- **Around R.T.**
Full ionization
Extrinsic doping

$$N_{d\Box}^+ \approx N_d, N_{a\Box}^- \approx N_a$$

$$(N_{d\Box}^+ - N_{a\Box}^-) \gg n_i$$

$$n_{o\Box} \approx (N_{d\Box}^+ - N_{a\Box}^-), p_{o\Box} = n_i^2 / n_{o\Box}$$

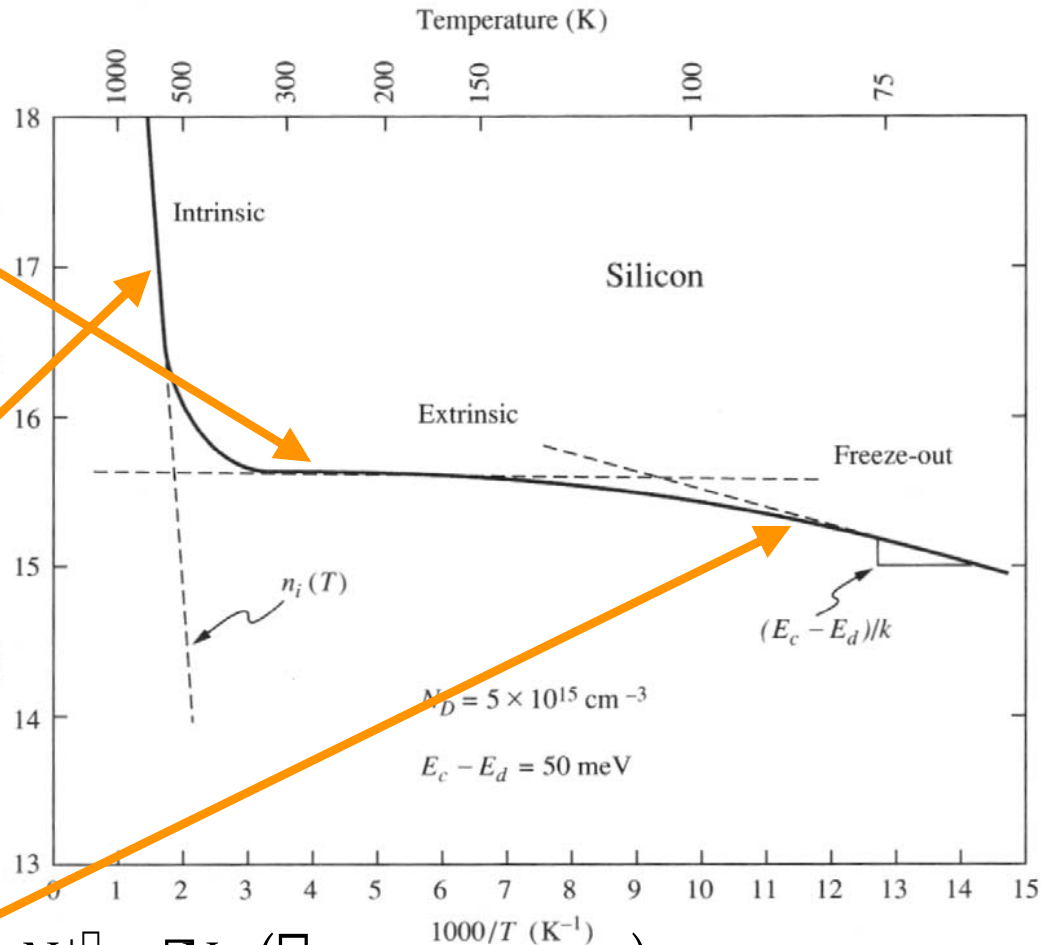
- **At very high T**
Full ionization
Intrinsic behavior

$$N_{d\Box}^+ \approx N_d, N_{a\Box}^- \approx N_a$$

$$n_i \gg (N_{d\Box}^+ - N_{a\Box}^-)$$

$$n_{o\Box} \approx p_{o\Box} \approx n_i$$

- **At very low T**
Incomplete ionization
Extrinsic doping, but with carrier freeze-out



$$N_{d\Box}^+ \ll N_{d\Box} \text{ (assuming n-type)}$$

$$(N_{d\Box}^+ - N_{a\Box}^-) \gg n_i$$

$$n_{o\Box} \approx (N_{d\Box}^+ - N_{a\Box}^-) \ll (N_{d\Box} - N_{a\Box}), p_{o\Box} = n_i^2 / n_{o\Box}$$

Uniform material with uniform excitations

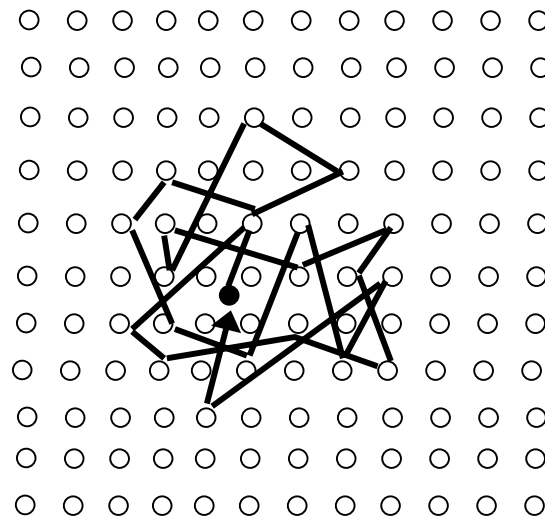
(pushing semiconductors out of thermal equilibrium)

A. Uniform Electric Field, E_x

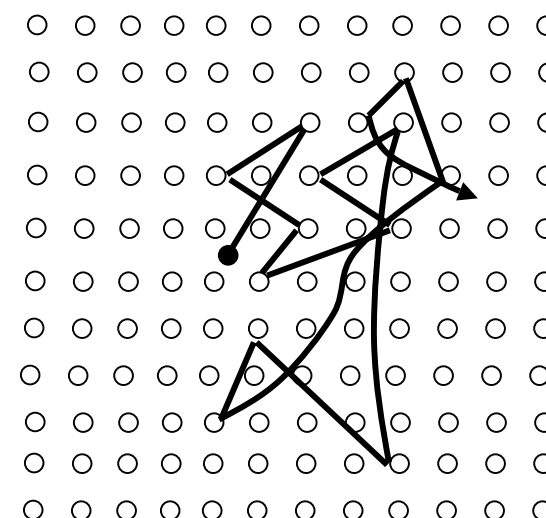
Drift motion:

Holes and electrons acquire a constant net velocity, s_x , proportional to the electric field:

$$\overline{s_{ex}} = -\mu_{e\Box} E_x, \quad \overline{s_{hx}} = \mu_h E_x \quad \mathcal{E}_x \leftarrow$$



No field



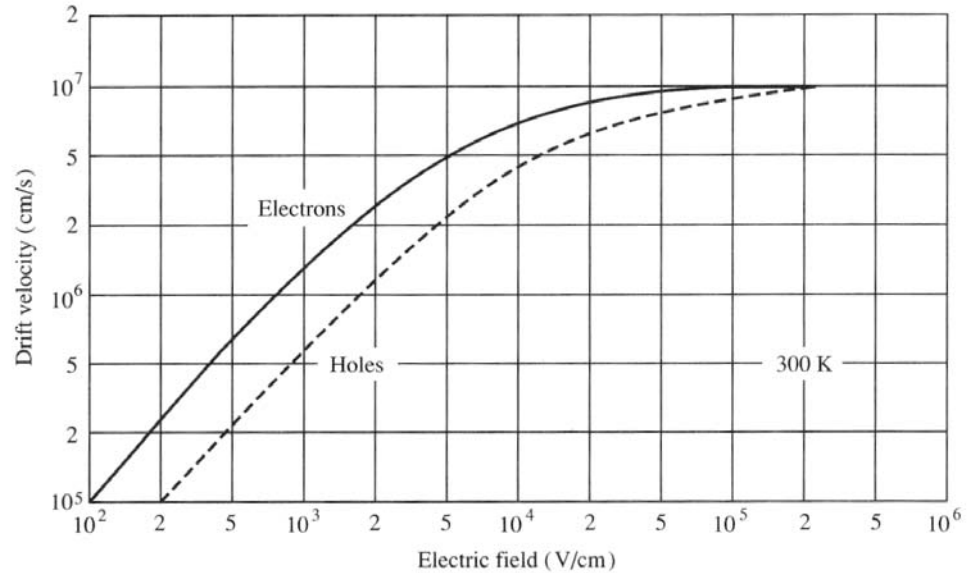
E-field applied

At low and moderate $|E|$, the mobility, μ , is constant.

At high $|E|$ the velocity saturates and μ decreases with increasing $|E|$.

Velocity saturation

The breakdown of Ohm's law at large electric fields.



See Fig 5.7, from Neamen, D. *Electronic Circuit Analysis and Design*. McGraw-Hill Higher Education, 1996.

Above: Velocity vs. field plot at R.T. for holes and electrons in Si (log-log plot). (Fonstad, Fig. 3.2)

Left: Velocity-field curves for Si, Ge, and GaAs at R.T. (log-log plot). (Neaman, Fig. 5.7)

Variation of mobility with temperature and doping levels

See Neamen, D. *Electronic Circuit Analysis and Design*. McGraw-Hill Higher Education, 1996.

See Fig 5.3, from Neamen, D. *Electronic Circuit Analysis and Design*. McGraw-Hill Higher Education, 1996.

Above: μ_e vs T in Si at several doping levels

Left: μ vs doping for Si, Ge, and GaAs at R.T. (Neaman, Fig. 5.3)

Variation of mobility with temperature and doping levels in Si near room temperature

See Fig 5.2, from Neamen, D. *Electronic Circuit Analysis and Design*. McGraw-Hill Higher Education, 1996.

Uniform material with uniform excitations

(pushing semiconductors out of thermal equilibrium)

A. Uniform Electric Field, E_x

Drift motion:

Holes and electrons acquire a constant net velocity, s_x , proportional to the electric field:

$$\overline{s_{ex}} = -\mu_e E_x, \quad \overline{s_{hx}} = \mu_h E_x$$

At low and moderate $|E|$, the mobility, μ , is constant.

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Drift currents:

Net velocities imply net charge flows, which imply currents:

$$J_{ex}^{dr} = -q n_o \overline{s_{ex}} = q \mu_e n_o E_x \quad J_{hx}^{dr} = q p_o \overline{s_{hx}} = q \mu_h p_o E_x$$

Note: Even though the semiconductor is no longer in thermal equilibrium the hole and electron populations still have their thermal equilibrium values.

Conductivity, σ_o :

Ohm's law on a microscale states that the drift current density is linearly proportional to the electric field:

$$J_{x\Box}^{dr\Box} = \sigma_o E_{x\Box}$$

The total drift current is the sum of the hole and electron drift currents. Using our early expressions we find:

$$J_x^{dr\Box} = J_{ex\Box}^{dr} + J_{hx\Box}^{dr} = q\mu_e n_o E_{x\Box} + q\mu_h p_o E_x = q(\mu_e n_{o\Box} + \mu_h p_o) E_{x\Box}$$

From this we see obtain our expression for the conductivity:

$$\sigma_{o\Box} = q(\mu_e n_{o\Box} + \mu_h p_o) \quad [\text{S/cm}]$$

Majority vs. minority carriers:

Drift and conductivity are dominated by the most numerous, or "majority," carriers:

$$\text{n-type} \quad n_{o\Box} \gg p_{o\Box} \Rightarrow \sigma_o \approx q\mu_e n_{o\Box}$$

$$\text{p-type} \quad p_{o\Box} \gg n_{o\Box} \Rightarrow \sigma_{o\Box} \approx q\mu_h p_o$$

Resistance, R, and resistivity, ρ_o :

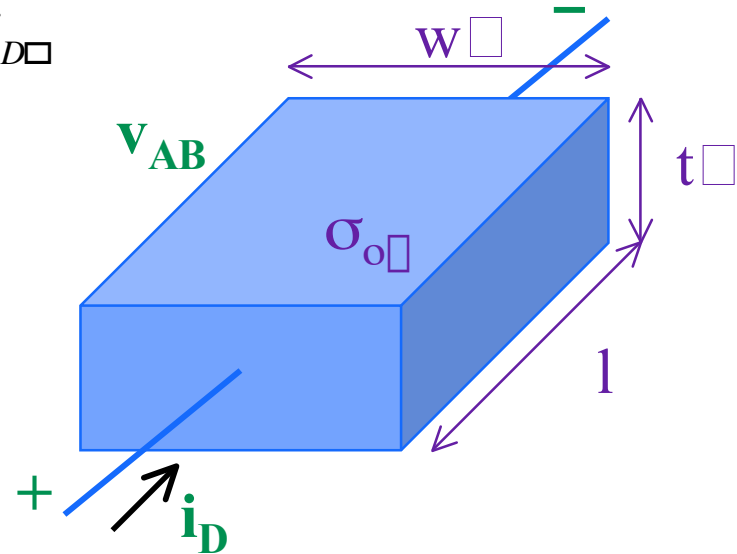
Ohm's law on a macroscopic scale says that the current and voltage are linearly related:

$$v_{ab} = R i_D$$

The question is, "What is R?"

We have: $J_x^{dr} = \sigma_o E_x$

with $E_x = \frac{v_{AB}}{l}$ and $J_x^{dr} = \frac{i_D}{w \cdot t}$



Combining these we find:

$$\frac{i_D}{w \cdot t} = \sigma_o \frac{v_{AB}}{l}$$

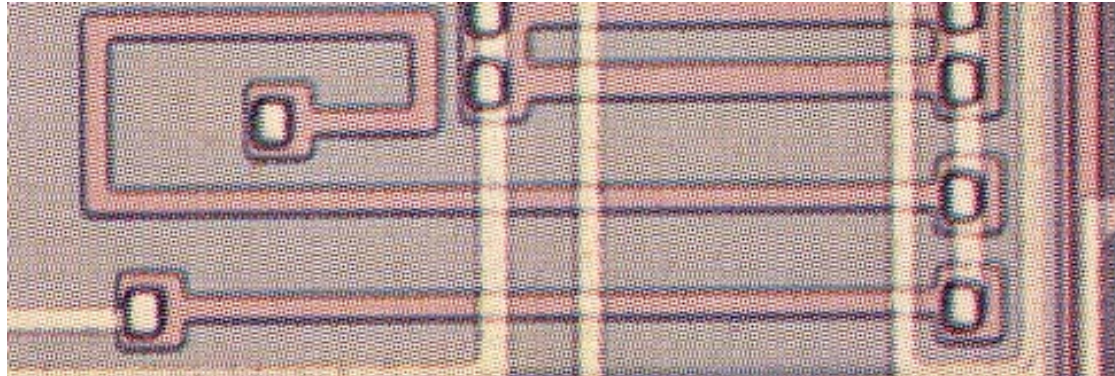
which yields: $v_{AB} = \frac{l}{w \cdot t \sigma_o} i_D = R i_D$

where $R \equiv \frac{l}{w \cdot t \sigma_o} = \frac{l}{w \cdot t} \rho_o$

Note: Resistivity, ρ_o , is defined as the inverse of the conductivity:

$$\rho_o \equiv \frac{1}{\sigma_o} \quad [\text{Ohm} \cdot \text{cm}]$$

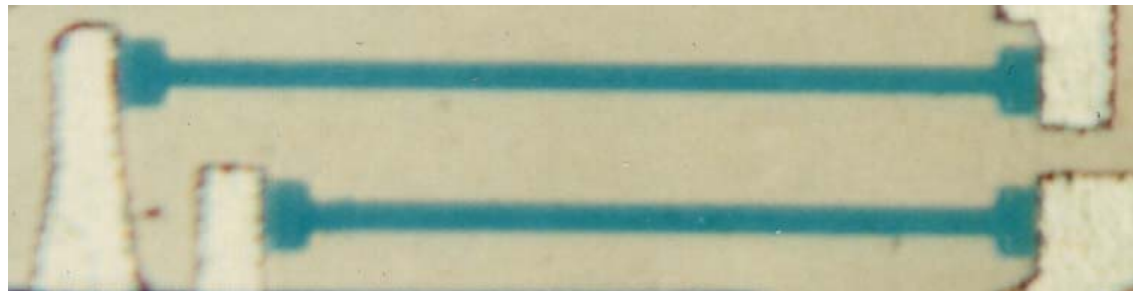
Integrated resistors Our first device!!



Courtesy of Fairchild Semiconductor. Used with permission.

(Courtesy of Fairchild Semiconductor Corporation. Fairchild Semiconductor, A Solid State of Progress, 1964.)

Diffused resistors: High sheet resistance semiconductor patterns (pink) with low resistance Al (white) "wires" contacting each end.



Thin-film resistors: High sheet resistance tantalum films (green) with low resistance Al (white) "wires" contacting each end.

Uniform material with uniform excitations

(pushing semiconductors out of thermal equilibrium)

B. Uniform Optical Generation, $g_L(t)$

Generation, G, and recombination, R:

If $G > R$, the populations are increasing

$$G > R \Rightarrow \frac{dn}{dt} = \frac{dp}{dt} > 0$$

If $G < R$, the populations are decreasing

$$G < R \Rightarrow \frac{dn}{dt} = \frac{dp}{dt} < 0$$

In general

$$\frac{dn}{dt} = \frac{dp}{dt} = G - R$$

Optical generation, $g_L(t)$:

In thermal equilibrium, $G_o = R_o = n_o p_o r(T)$

With $g_L(t)$:

$$G = G_o + g_L(t)$$
$$R = n p r(T) = (n_o + n')(p_o + p') r(T)$$

Lecture 2 - Uniform Excitation - Summary

- **Review**

Frozen out if too cold; intrinsic-like if too hot

- **Uniform excitation of uniform samples: drift**

n, p unchanged; carriers attain constant velocity, s (viscous flow)

$s = \mu E$ at low to moderate fields; s saturates at high fields

$$J_{e,\text{drift}} = q \mu_e n_o E \text{ [A/cm}^2\text{]}, J_{h,\text{drift}} = q \mu_h p_o E \text{ [A/cm}^2\text{]}$$

$$J_{\text{drift}} = J_{e,\text{drift}} + J_{h,\text{drift}} = q (\mu_e n_o + \mu_h p_o) E = \sigma_o E \quad \text{the majority rules}$$

Mobility, μ , decreases as temperature goes up from RT

- **Uniform excitation: optical generation**

In TE, $G_o(T) = R_o(T) = n_o p r(T)$

Uniform illumination adds uniform generation term, $g_L(t)$

Populations increase: $n_o \rightarrow [n_o + n', p_o \rightarrow [p_o + p', \text{ and } n' \rightleftharpoons p'$

$$dn'/dt = dp'/dt = G_o(T) + g_L(t) - R(T) = g_L(t) - [np - n_o p_o] r(T)$$

focus is on minority $\approx g_L(t) - n'/\tau_{\text{min}}$ if LLI holds, with $\tau_{\text{min}} \equiv [p_o r(T)]^{-1}$

- **Uniform excitation: both optical and electrical**

$$\begin{aligned} \text{Photoconductivity: } \sigma_o &\rightarrow [\sigma_o + \sigma' = q [\mu_e (n_o + n') + \mu_h (p_o + p')]] \\ &= \sigma_o + q (\mu_e + \mu_h) 2 p' \end{aligned}$$

Photoconductors: an important class of light detectors