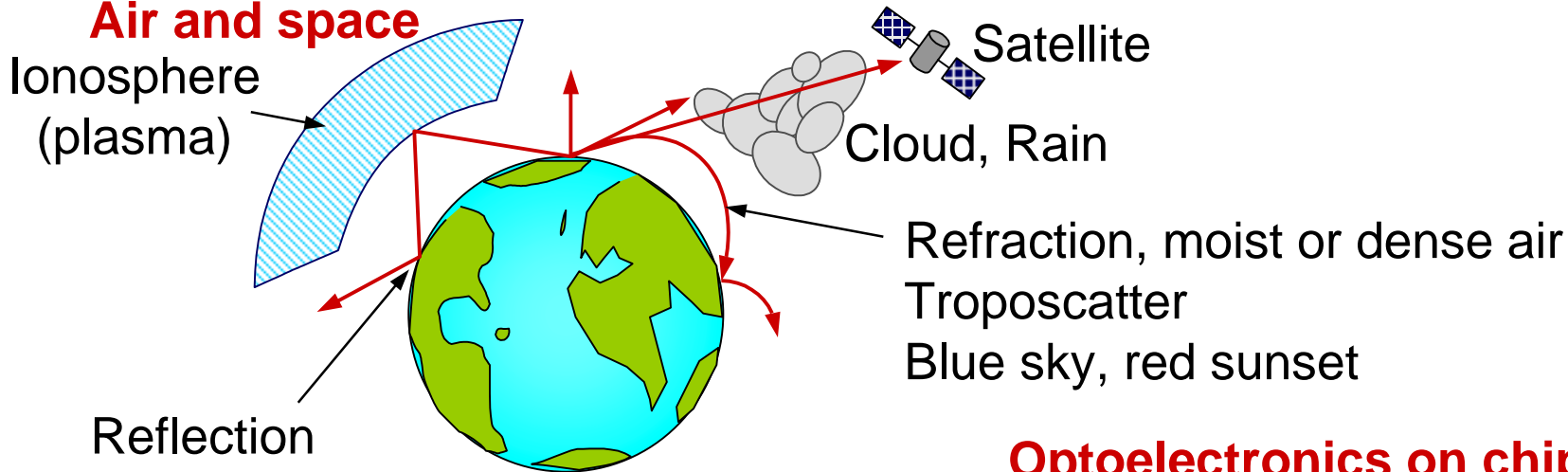


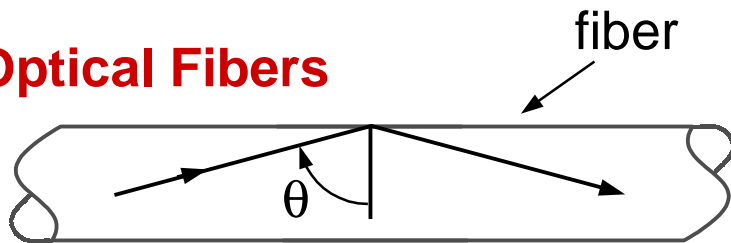
WAVES IN MEDIA

Significance to Communications

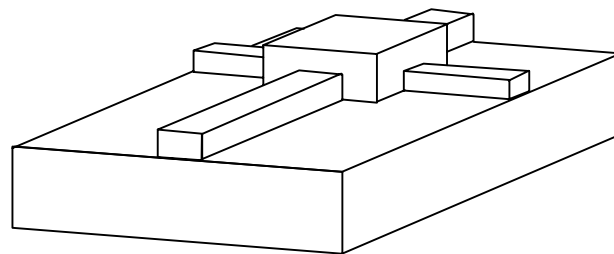
Air and space



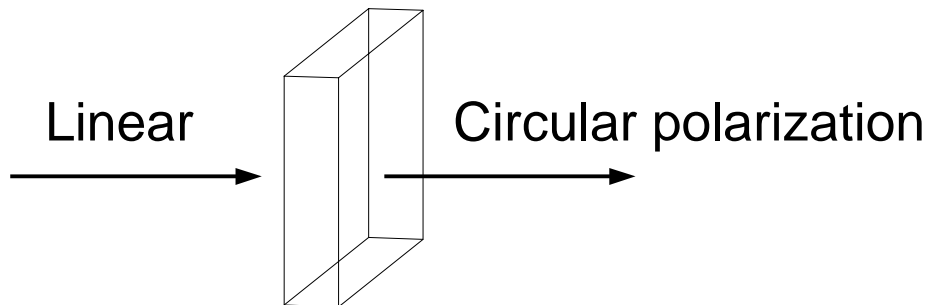
Optical Fibers



Optoelectronics on chips

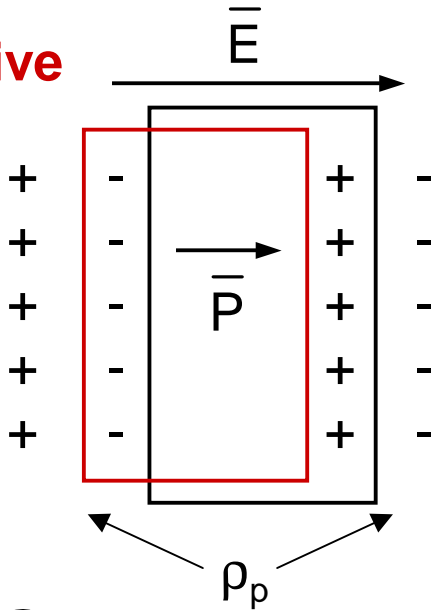


Polarization-based optoelectronic devices



WAVES IN MEDIA

Constitutive Relations



Vacuum:

$$\bar{\mathbf{D}} = \epsilon_0 \bar{\mathbf{E}} \quad \nabla \cdot \bar{\mathbf{D}} = \rho_f$$

ρ_f = free charge density

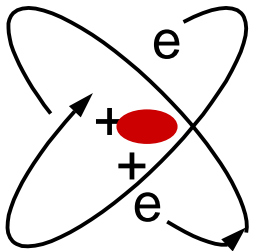
Dielectric Materials:

$$\bar{\mathbf{D}} = \epsilon \bar{\mathbf{E}} = \epsilon_0 \bar{\mathbf{E}} + \bar{\mathbf{P}}$$

$$\nabla \cdot \epsilon_0 \bar{\mathbf{E}} = \rho_f + \rho_p$$

$$\nabla \cdot \bar{\mathbf{P}} = -\rho_p \quad \text{polarization charge density}$$

$\bar{\mathbf{P}}$ = "Polarization Vector"



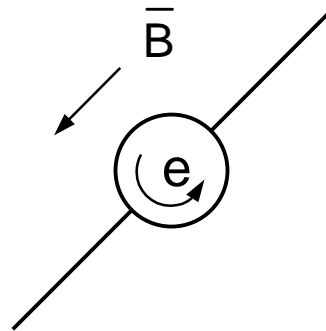
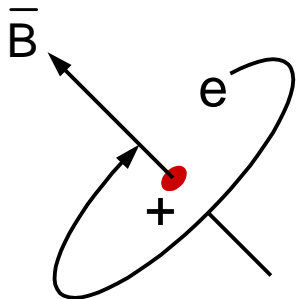
Magnetic Materials:

$$\nabla \cdot \bar{\mathbf{B}} = 0$$

$$\bar{\mathbf{B}} = \mu_0 \bar{\mathbf{H}} \quad \text{in vacuum}$$

$$\bar{\mathbf{B}} = \mu \bar{\mathbf{H}} = \mu_0 (\bar{\mathbf{H}} + \bar{\mathbf{M}})$$

$\bar{\mathbf{M}}$ = "Magnetization Vector"



TYPES OF MEDIA

Properties are a function of:

Field direction

Position

Time: $\neq f(t)$

$\neq f(\text{history})$

Frequency

$\bar{\mathbf{E}}$ or $\bar{\mathbf{H}}$

Temperature

Pressure

Designation:

Anisotropic $\bar{\mathbf{D}} = \bar{\bar{\epsilon}} \bar{\mathbf{E}}$, $\bar{\mathbf{B}} = \bar{\bar{\mu}} \bar{\mathbf{H}}$

Inhomogeneous

Stationary

Amnesic

Dispersive

Non-linear

Temperature dependent

Compressive

ANISOTROPIC DIELECTRICS

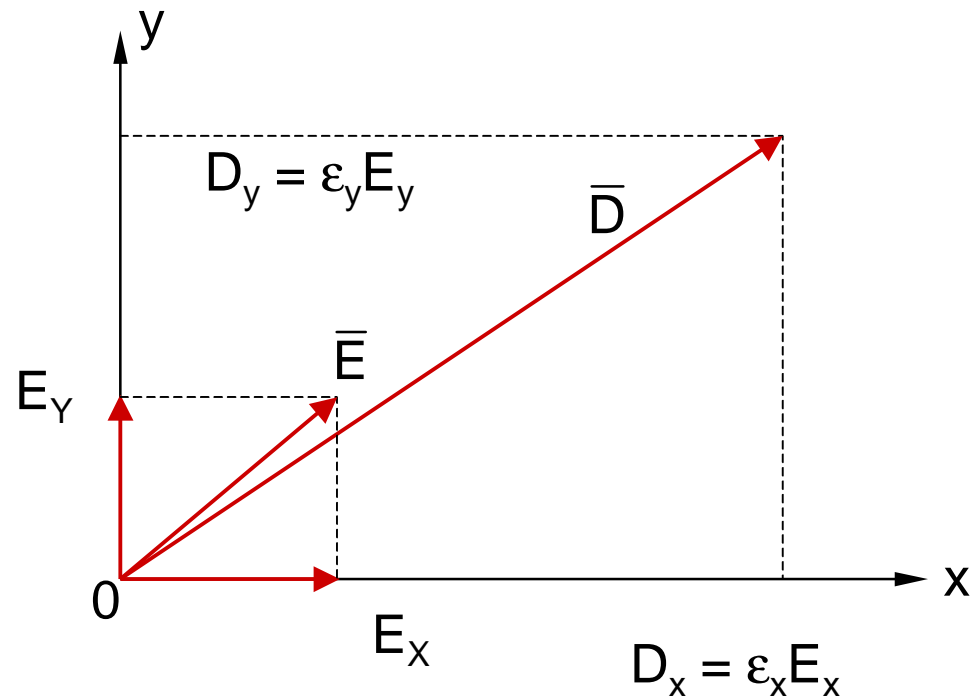
$$\underline{\bar{D}} = \underline{\bar{\epsilon}} \underline{\bar{E}}$$

$$D_x = \epsilon_{xx} E_x + \epsilon_{xy} E_y + \epsilon_{xz} E_z$$

$$D_y = \epsilon_{yx} E_x + \epsilon_{yy} E_y + \epsilon_{yz} E_z$$

$$D_z = \epsilon_{zx} E_x + \epsilon_{zy} E_y + \epsilon_{zz} E_z$$

$$\text{Let } \underline{\bar{\epsilon}} = \begin{bmatrix} \epsilon_x & 0 & 0 \\ 0 & \epsilon_y & 0 \\ 0 & 0 & \epsilon_z \end{bmatrix}$$



x, y, z are “principal axes”

Note: $\underline{\bar{D}} // \underline{\bar{E}}$ iff $\underline{\bar{E}} // \hat{x}, \hat{y},$ or \hat{z} for $\epsilon_x \neq \epsilon_y \neq \epsilon_z$

Real $\underline{\bar{\epsilon}}, \underline{\bar{\mu}} \Rightarrow$ Lossless medium

HOW TO MAKE ANISOTROPIC MATERIALS

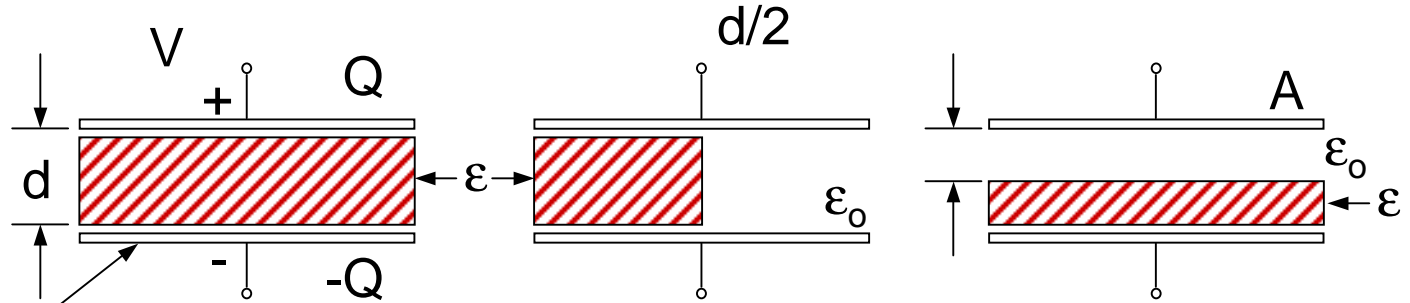
Consider:

$$\epsilon \gg \epsilon_0$$

(capacitors)

$$Q = CV$$

Area A (m^2)



$$C = \frac{\epsilon_{\text{eff}} A}{d}$$

$$\epsilon_{\text{eff}} = \epsilon$$

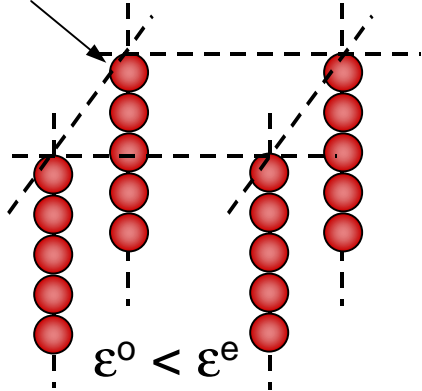
$$C \cong \frac{\epsilon (A/2)}{d}$$

$$\epsilon_{\text{eff}} \cong \epsilon/2$$

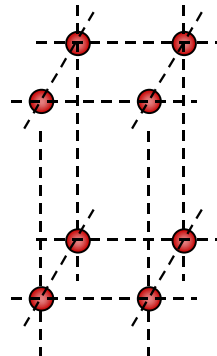
$$C \cong \frac{\epsilon_0 A}{(d/2)}$$

$$\epsilon_{\text{eff}} \cong 2\epsilon_0$$

atom, molecule



“uniaxial medium”

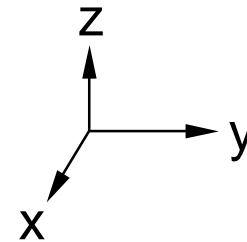


$$\epsilon_x = \epsilon_y = \epsilon^0$$

$$\epsilon_z = \epsilon^e$$

“ordinary”

“extraordinary”



WAVE BEHAVIOR IN UNIAXIAL MEDIUM

Assume wave in $+\hat{z}$ direction, $\sigma = 0$

Derive wave equation:

$$\begin{aligned} \nabla \times \underline{\underline{E}} &= -j\omega \underline{\underline{B}} & \nabla \cdot \underline{\underline{D}} &= \rho_f = 0 & \underline{\underline{D}} &= \underline{\underline{\epsilon}} \underline{\underline{E}} & \underline{\underline{\epsilon}} &= \begin{bmatrix} \epsilon^e & 0 & 0 \\ 0 & \epsilon & 0 \\ 0 & 0 & \epsilon \end{bmatrix}, & \underline{\underline{\mu}} &= \mu \\ \nabla \times \underline{\underline{H}} &= j\omega \underline{\underline{D}} & \nabla \cdot \underline{\underline{B}} &= 0 \end{aligned}$$

Therefore $\nabla \times (\nabla \times \underline{\underline{E}}) = \nabla (\nabla \cdot \underline{\underline{E}}) - \nabla^2 \underline{\underline{E}} = -j\omega \mu \nabla \times \underline{\underline{H}} = \omega^2 \underline{\underline{\mu}} \underline{\underline{\epsilon}} \underline{\underline{E}}$

Does $\nabla \cdot \underline{\underline{E}} = 0$ here? Yes, (let's skip proof) can test final solution

Therefore $\nabla^2 \underline{\underline{E}} + \omega^2 \underline{\underline{\mu}} \underline{\underline{\epsilon}} \underline{\underline{E}} = 0 \Rightarrow$ 3 equations (x,y,z components)

$$\underbrace{\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right]}_{\nabla^2} [\hat{x} \underline{\underline{E}}_x + \hat{y} \underline{\underline{E}}_y + \hat{z} \underline{\underline{E}}_z] + \omega^2 \underline{\underline{\mu}} \underline{\underline{\epsilon}} \underline{\underline{E}} = 0$$

Assume $= 0$ (UPW in z direction)

This leads to 2 decoupled equations for x and y polarization

BIREFRINGENT MEDIA

Decoupled wave equations:

$$\left[\frac{\partial^2}{\partial z^2} + \underbrace{\omega^2 \mu \epsilon^e}_{\triangleq (k^e)^2} \right] \underline{E}_x = 0, \quad k^e = \omega \sqrt{\mu \epsilon^e}, \quad \text{(x-pol equation)}$$
$$\left[\frac{\partial^2}{\partial z^2} + \underbrace{\omega^2 \mu \epsilon}_{\triangleq (k^o)^2} \right] \underline{E}_y = 0, \quad k^o = \omega \sqrt{\mu \epsilon} \quad \text{(y-pol equation)}$$

$$\text{Where } \underline{E}_x \propto e^{-jk^e z} = e^{-j(\omega/v^e)z} \Rightarrow \begin{cases} v^e = 1/\sqrt{\mu \epsilon^e} \\ v^o = 1/\sqrt{\mu \epsilon} \end{cases}$$

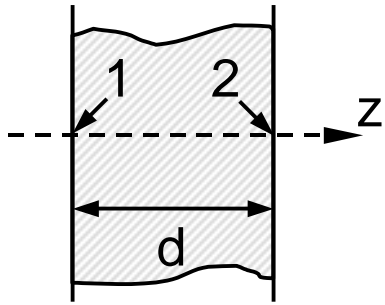
Thus the x- and y-polarized waves propagate independently at different velocities

If $v^e < v^o$ then $v^e \rightarrow$ “slow-axis velocity”

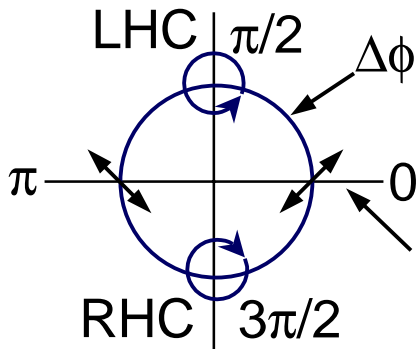
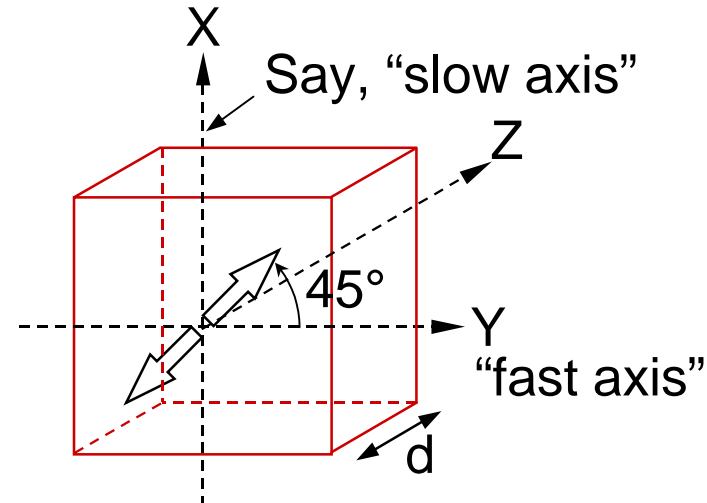
BIREFRINGENT MEDIA

Example:

$$\bar{E}_1 = E_0 (\hat{x} + \hat{y}) \quad 45^\circ \text{ linear pol. input}$$

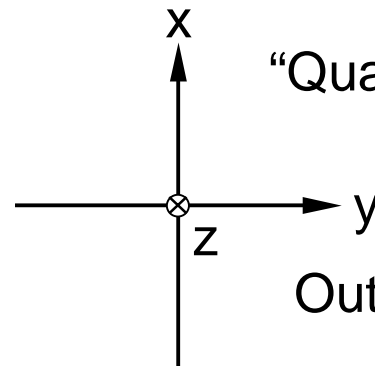
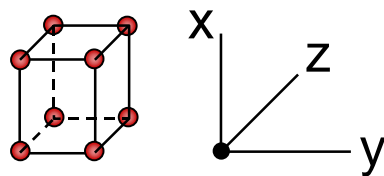


$$\bar{E}_2 = E_0 \underbrace{\hat{x}e^{-jk^e d} + \hat{y}e^{-jk^o d}}_{\text{What pol. ?}}$$



$$\Delta\phi \triangleq \phi^e - \phi^o = (k^e - k^o)d$$

Linear pol.



"Quarter wave plate"

$$\text{Output: } d \ni \Delta\phi = \pi/2$$

Demo; Polaroids

1) $\updownarrow \oplus \updownarrow \Rightarrow$

4) $\updownarrow \boxtimes_S \leftrightarrow \Rightarrow$

2) $\updownarrow \oplus \leftrightarrow \Rightarrow$

5) $\updownarrow \text{MICA} \leftrightarrow \Rightarrow$

3) $\updownarrow \boxplus_F \leftrightarrow \Rightarrow$

6) GEARS

