

LASERS

Applications:

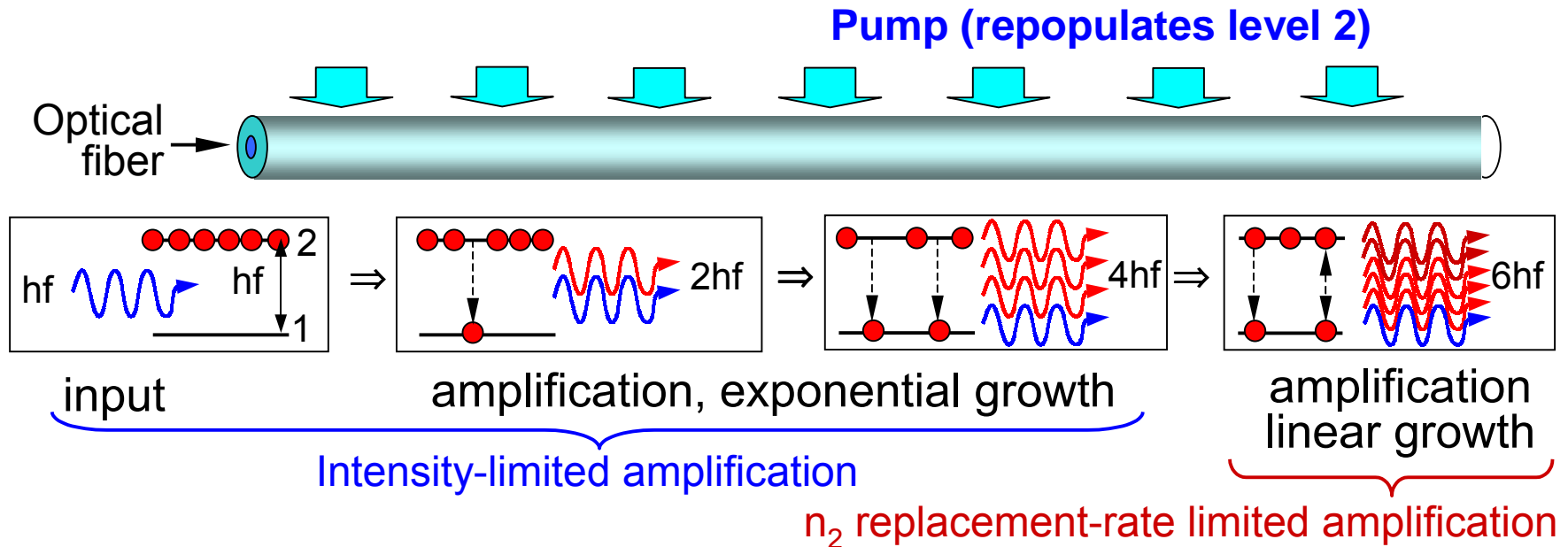
- Amplification: Broad-band communications links
(e.g. EFDA; avoids down-conversion)
- Oscillator: Frequency/distance reference, local oscillators, illuminators,
sources for fiber communications, CD/DVD players
- Focused power: Laser machining, weapons, laser fusion (pellet compression).
Peak $> 10^{15} \text{W}$ (10- μm spot $\Rightarrow \bar{E} \approx 10^{14} \text{Vm}^{-1}$ vs 10^6 in H atom)
Average $> 1 \text{kw}$; high intensity because $I \propto |\sum_i \bar{E}_i|^2$

Basic Principles:

- Quantum states characterize atoms and molecules in gases, impurities in solids,
and electrons and holes in semiconductors
- A transition to a lower state emits a photon coherent with the triggering photon
 \Rightarrow amplification (or, with internal reflection, oscillation)
- Amplification/lasing requires upper state population to exceed lower state

BASIC LASER AMPLIFIER PHYSICS

Basic Amplification Process:

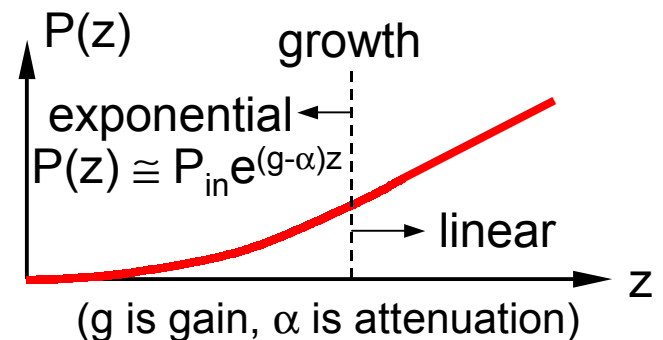


[Each ● is a separate atom or molecule; need $n_2 > n_1$ for amplification]

Amplification frequency f [Hz]:

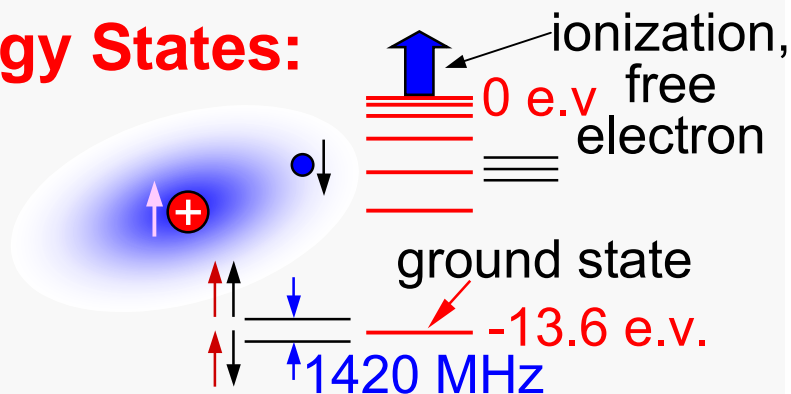
$$E_2 - E_1 = hf \text{ [J]}$$

$$h = 6.625 \times 10^{-34} \text{ [Js]}$$



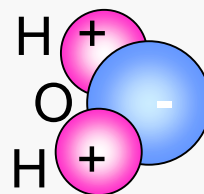
ENERGY STATES AND POPULATIONS

Energy States:



Hydrogen atom (Galactic arms)

States:



electronic (visible, UV)
vibrational (visible)
bending (IR)
rotational (microwave)

Water vapor H₂O

(e.g. water vapor masers around stars)
(electric dipole transitions)

Also: Chromium atoms in lattice (e.g. ruby), Erbium atoms in glass

Level Populations—Kinetic Temperature T_k :

Thermal equilibrium dominated by collisions

⇒ Boltzmann distribution:

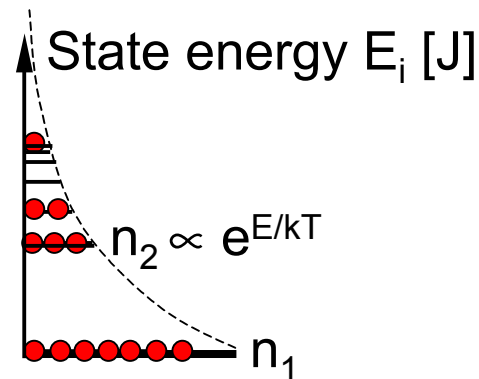
$$\frac{n_i}{n_j} = e^{-(E_i - E_j)/kT_k}$$

If thermal equilibrium dominated by radiation, $T_{i,j}$

⇒ Boltzmann distribution:

$$\frac{n_i}{n_j} = e^{-(E_i - E_j)/kT_{\text{rad}}}$$

$n_2 > n_1$ if $T_{\text{rad}} < 0$



⇒ $n_i \rightarrow n_j$ if $T_{\text{rad}} \rightarrow \infty$

“A” AND “B” COEFFICIENTS

Rate Equation:

Assume: Two-level system, $E_2 > E_1$, and $n_i =$ atoms m^{-1} in state i

Then: $dn_2/dt = - \underbrace{An_2}_{\text{Spontaneous emission}} - \underbrace{B(n_2 - n_1)}_{\text{Induced emission}} [m^{-1}s^{-1}]$ (collisionless system)

Spontaneous emission Induced emission

Spontaneous emission between states i and j :

$$A_{ij} = \omega^3 |D_{ij}|^2 (2/3h\epsilon c^3) [s^{-1}]$$

D_{ij} [Cm] is a quantum mechanical dipole moment (electric or magnetic)

Decay time $\tau_A = A^{-1}$

Note: $\tau_A \propto \omega^{-3}$, so “visible” τ 's are very short, microwave τ 's are long

Stimulated emission and absorption:

B coefficient:

$$B_{ij} = F\sigma_{ij} \propto F g_{ij}(f) A_{ij}/\omega^3$$

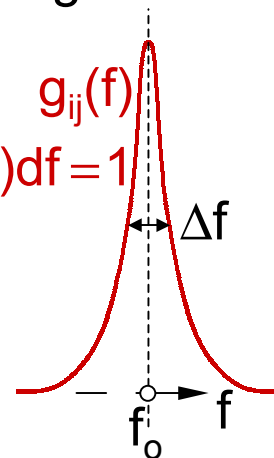
$$\int_{-\infty}^{\infty} g_{ij}(f) df = 1$$

Photon flux density F :

$$F = |\underline{E}|^2 / 2\eta_0 hf \quad [\text{photons } m^{-2} s^{-1}]$$

Lorentzian line shape $g_{ij}(f)$:

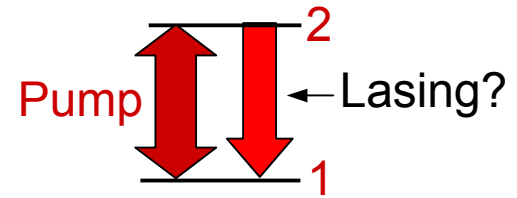
$$g_{ij}(f) = [2/\pi(\Delta f)] / [1 + 4(f - f_0)^2 / (\Delta f)^2]$$



PUMPING LASERS

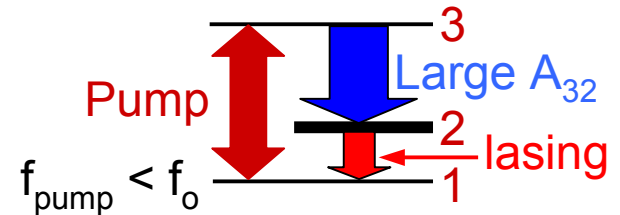
Two-Level Lasers:

No degree of pumping can yield $n_2 > n_1$
 (early 2-level lasers spatially isolated n_2 group)

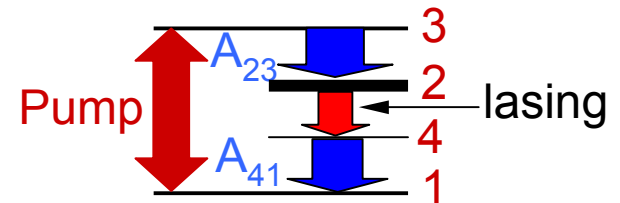


Three-Level Lasers:

Pump levels 1,3 so $n_1 \cong n_3$
 Large A_{32} populates 2 so $n_2 \gg n_1 \cong n_3 \cong 0$



More levels sometimes used, e.g. to utilize quantum states with larger A's



Laser Power Efficiency (P_{out}/P_{in}):

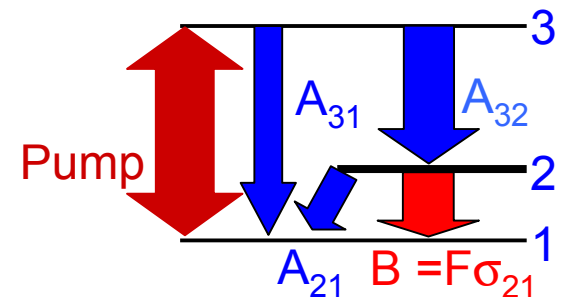
Intrinsic efficiency: $\eta_i = f_L/f_p$ ($P \propto nhf$ [W]) < 1

B/A efficiency: $\eta_B = B_{21}/(A_{21}+B_{21}) < 1$

A/A efficiency: $\eta_A = A_{32}/(A_{31}+A_{32}) < 1$

Total efficiency: $\eta = \eta_i \eta_B \eta_A$

Recall $A \propto \omega^3$; if $B \gg A \propto \omega^3$, then x-ray lasers need very high B (pump values) ($B \neq \sim f(\omega)$)



LINE SHAPE

Lorentzian Line Shape and Broadening Mechanisms:

Lorentzian line shape:

$$g_{ij}(f) = [2/\pi(\Delta f)] / [1 + 4(f - f_0)^2/(\Delta f)^2]$$

$g_{ij}(f)$ has unity integral:

$$\int_{-\infty}^{\infty} g_{ij}(f) df = 1$$

$$g_{ij}(f)/g_0 = 0.5 \text{ for } |f - f_0| = \Delta f/2$$

Broadening mechanisms:

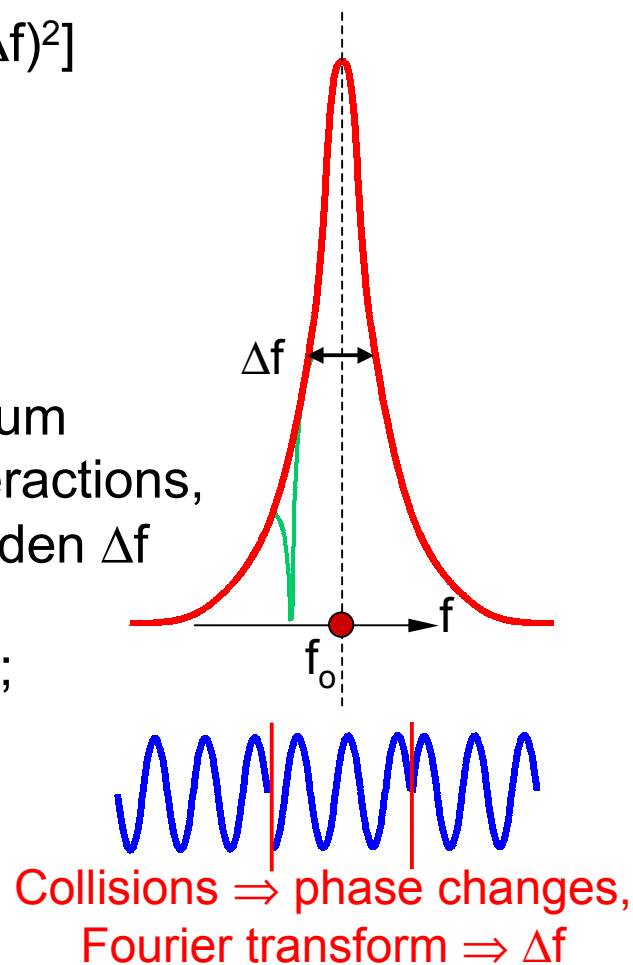
$\Delta f_0 > 1/\tau_A = A \cong 10$ MHz (minimum linewidth); collisions, lattice interactions, Fields (\vec{E}, \vec{B}), & Doppler all broaden Δf

Homogeneous broadening:

E.g., each atom has $\Delta f \cong 4$ THz; a single frequency can drain G; e.g. EDFA's, most solid-state and semiconductor lasers

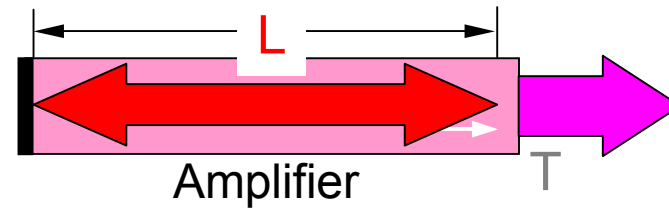
Inhomogeneous broadening:

Each narrow-band atom shifted differently, e.g. HeNe; \Rightarrow hole burning



LASER OSCILLATORS

Laser Oscillation:



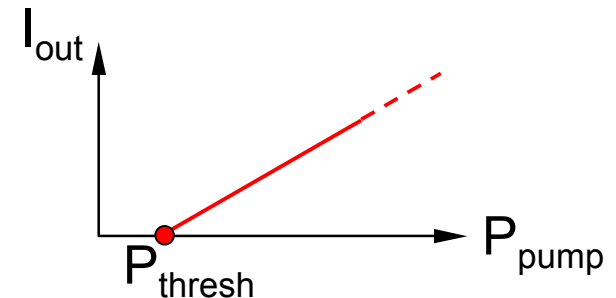
Oscillator: Assume length L , perfect mirrors at both ends;
Closed lossless amplifier must oscillate and saturate

Threshold: Gain m^{-1} must exceed loss (threshold condition)
Gain \propto pump power P_p , therefore $P_p >$ threshold too

Mirrors: Assume one mirror has power transmission coefficient $T > 0$
Gain \cong Loss: $P_+(1 - T)e^{2(g-\alpha)L} \geq P_+$
 \Rightarrow Two-pass gain $e^{2(g-\alpha)L} \geq 1/(1 - T) > 1$ for oscillation ($g > \alpha$)

Output Limit: $P_{out} = \eta P_{pump}$. $P_+ = P_{out} / T$, so $P_+ / P_{out} \rightarrow \infty$ as $T \rightarrow 0$

Q-switching: Set $T \cong 0$ until P_+ peaks,
then set $T \cong 1$; yields very large
“Q-switched pulse”



LASER RESONANCES

Oscillator Resonant Frequencies f :

Resonances when $\{m\lambda_m/2 = L\}$ (mirrors approx. short circuits)

$$\Rightarrow \lambda_m = 2L/m, f_m = cm/2LN \quad (N = \text{index of refraction})$$

$$f_{i+1} - f_i = c/2LN \cong 10^8 \text{ Hz (100 MHz) for 1-meter fiber;}$$

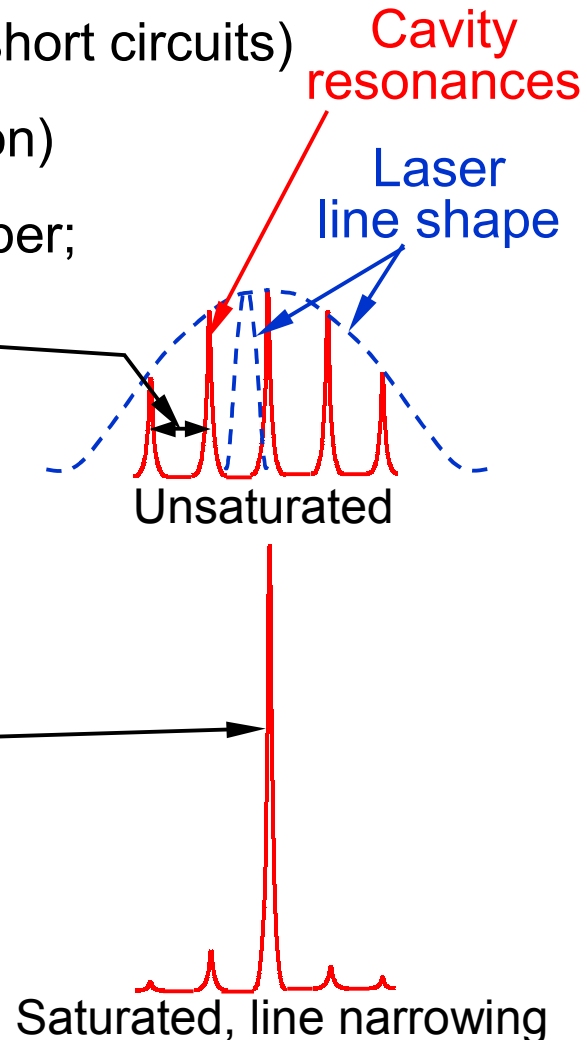
$$\cong 50 \text{ GHz line spacing for 0.5-mm diodes}$$

Laser Output Spectrum:

In saturation, $\text{Gain}(f) \cong \text{frequency independent}$

If laser linewidth $\Delta f >$ line spacing,
dominant line wins if exponential growth

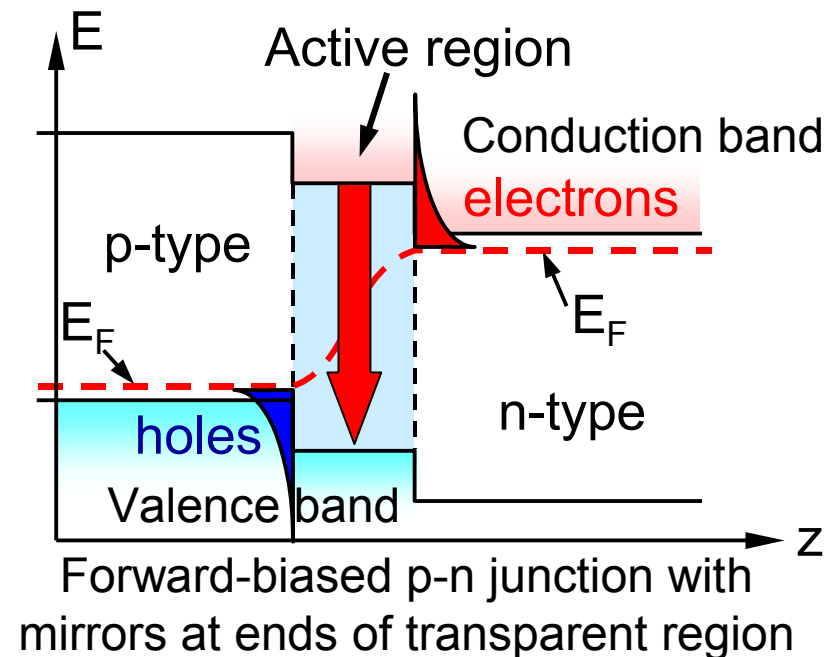
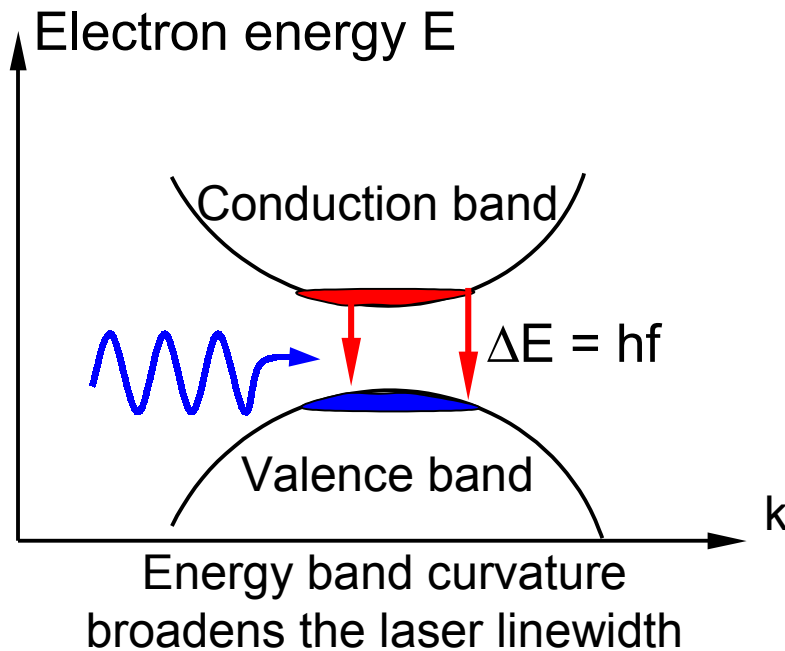
If linewidth $\Delta f <$ line spacing,
must tune cavity length to f_0



EXAMPLES OF LASERS

Electrically Pumped Solid-State Lasers:

Forward biased GaAs p-n junction injects carriers into conduction band
Compact (grain of sand), ~50 percent efficiency, $>100 \text{ W/cm}^2$ for arrays,
 1 mW/micron^2 for diodes (1-1000 mW typical)

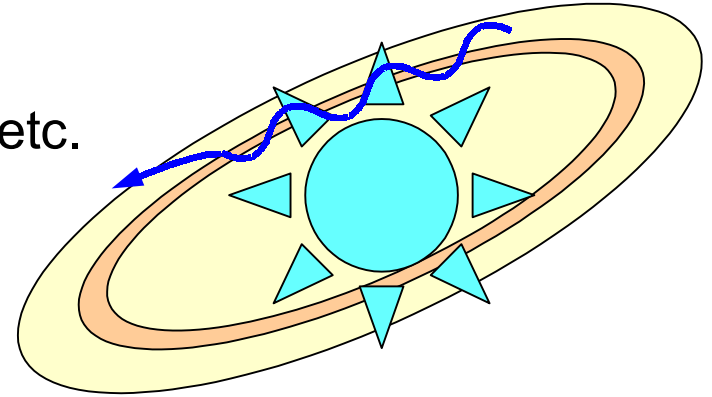


EXAMPLES OF LASERS (2)

Astrophysical Masers:

Stellar pumped: Water vapor, OH, CO, etc.

Interstellar collisions: OH, etc.



Gas Lasers:

Ammonia (23 GHz): “Pumped” by diverting molecules in ground state

CO₂, HeNe: Pumped by electrical discharges that form energetic plasmas

Chemical: Chemical combustion yields upper-state excess

Externally Pumped Solid-State Lasers

Ruby: Pumped by flash lamps, etc.

EDFA: Pumped by semiconductor lasers