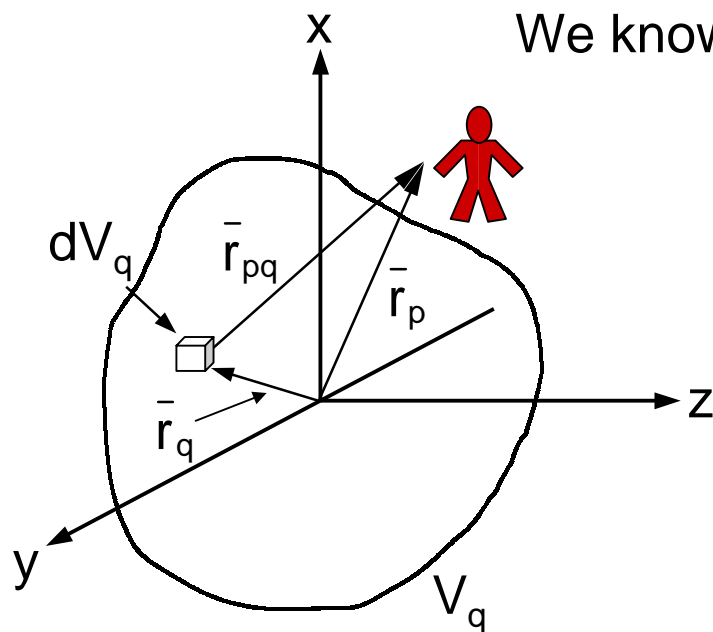


RADIATION BY CURRENTS

How are waves generated? How do we build antennas?



We know: $\nabla \cdot \bar{\mathbf{B}} = 0$, Therefore $\bar{\mathbf{B}} = \nabla \times \bar{\mathbf{A}}$

$$\bar{\mathbf{A}} = (\mu_0/4\pi) \int_{V_q} \mathbf{J}_q(t - r_{pq}/c) r_{pq}^{-1} dV_q$$

$$\nabla \times \bar{\mathbf{H}} = \bar{\mathbf{J}} + \partial \bar{\mathbf{D}}/\partial t, \text{ so } \bar{\mathbf{E}} = +\epsilon^{-1} \int (\nabla \times \bar{\mathbf{H}}) dt$$

Algorithm is:

$$\bar{\mathbf{J}}(t, \bar{\mathbf{r}}) \rightarrow \bar{\mathbf{A}} \rightarrow \bar{\mathbf{H}} \rightarrow \bar{\mathbf{E}}$$

Simplify:

Since radiation = $f(\omega)$, switch to s.s.s. (sinusoidal steady state)

Sinusoidal Steady State:

$$\begin{aligned} x(t) &= \text{Re} \{ \underline{x} e^{j\omega t} \} \\ &= \text{Re} \{ [\text{Re} \{ \underline{x} \} + j \text{Im} \{ \underline{x} \}] [\cos \omega t + j \sin \omega t] \} \\ &= \text{Re} \{ \underline{x} \} \cos \omega t - \text{Im} \{ \underline{x} \} \sin \omega t \end{aligned}$$

EXPRESSIONS FOR RADIATION

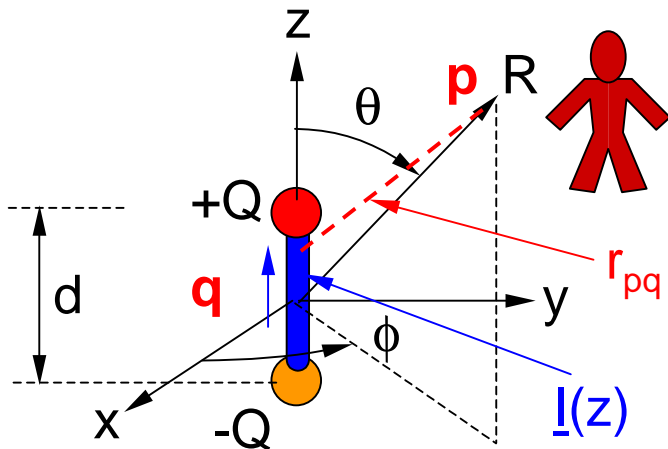
Sinusoidal Steady State Expressions:

$$\bar{\mathbf{B}} = \nabla \times \bar{\mathbf{A}}$$

$$\bar{\mathbf{A}} = (\mu_0/4\pi) \int_{V_q} \bar{\mathbf{J}}_q e^{-j\bar{\mathbf{k}} \cdot \bar{\mathbf{r}}_{pq}} r_{pq}^{-1} dv_q$$

$$\nabla \times \bar{\mathbf{H}} = \bar{\mathbf{J}} + j\omega \bar{\mathbf{D}}, \text{ so } \bar{\mathbf{E}} = (\nabla \times \bar{\mathbf{H}})/j\omega\epsilon$$

Assume short current element \mathbf{I} , $d \ll \lambda/2\pi$:



Let $d \ll \lambda/2\pi$,

Then $e^{-j\bar{\mathbf{k}} \cdot \bar{\mathbf{r}}} \cong e^{-jkR}$ ($k = 2\pi/\lambda$)

Note: $\int_{V_q} \bar{\mathbf{J}} dv_q = \int_{V_q} \mathbf{I}(z) dz = I_0 d$

Therefore: $\bar{\mathbf{A}} = (\mu_0/4\pi) \int_{V_q} \bar{\mathbf{J}}_q e^{-j\bar{\mathbf{k}} \cdot \bar{\mathbf{r}}} r_{pq}^{-1} dv_q$

$$= (\mu_0/4\pi R) e^{-jkR} \int_{V_q} \mathbf{I}(z) dz$$

$$\bar{\mathbf{A}} = \hat{\mathbf{z}} \mu_0 I_0 d e^{-jkR} / 4\pi R$$

EXPRESSIONS FOR RADIATION

Finding radiated $\bar{\mathbf{H}}$:

We saw: $\bar{\mathbf{A}}(r) = \hat{\mathbf{z}}\mu_0 I_0 d e^{-jkr} / 4\pi r$ (r = distance to origin)

Then: $\bar{\mathbf{H}} = \nabla \times \bar{\mathbf{A}} / \mu_0$

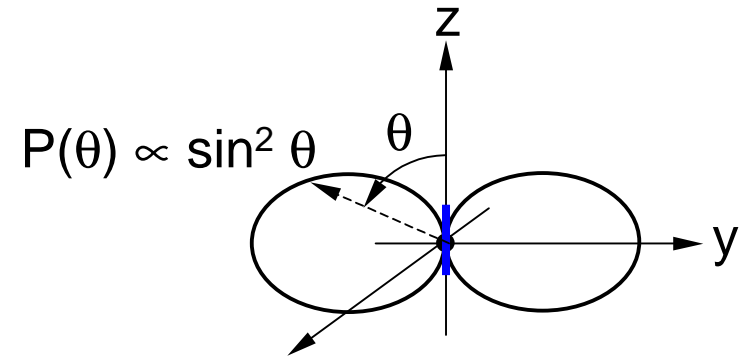
Where: $\nabla \times \bar{\mathbf{A}} = (r^2 \sin \theta)^{-1} \begin{vmatrix} \hat{\mathbf{r}} & r\hat{\theta} & r \sin \theta \hat{\phi} \\ \partial/\partial r & \partial/\partial \theta & \partial/\partial \phi \\ A_r & rA_\theta & r \sin \theta A_\phi \end{vmatrix}$

Therefore: $\bar{\mathbf{H}} = (\hat{\phi} / \mu_0 r) [\partial(rA_\theta) / \partial r - \partial A_r / \partial \theta]$
 $= \hat{\phi} (j I_0 d e^{-jkr} / 2\lambda r) [1 + (jkr)^{-1}] \sin \theta$

EXPRESSIONS FOR RADIATION

Finding radiated $\underline{\bar{E}}$:

Recall: $\underline{\bar{E}} = (\nabla \times \underline{\bar{H}}) / j\omega\epsilon$



Therefore: $\underline{\bar{E}} = (j \hat{I} d e^{-jkr} / 2\lambda r) \eta_0$

$$\left\{ \hat{\theta} \left[1 + (jkr)^{-1} + (jkr)^{-2} \right] \sin \theta + \hat{r} \left[(jkr)^{-1} + (jkr)^{-2} \right] 2 \cos \theta \right\}$$

Radiated power:

$$\hat{r} P(\theta, \phi, r) [\text{Wm}^{-2}] = 0.5 R_e \{ \underline{\bar{S}} \} = 0.5 R_e \{ \underline{\bar{E}} \times \underline{\bar{H}}^* \}$$

So: $\hat{r} P(\theta, \phi, r) \cong \hat{r} (\eta_0 / 2) |I d / 2\lambda r|^2 \sin^2 \theta [\text{Wm}^{-2}]$ in far field, $kr \gg 1$

Where: $\underline{\bar{H}} = \hat{\phi} \left[j I d e^{-jkr} / 2\lambda r \right] \left[1 + (jkr)^{-1} \right] \sin \theta$

$\eta_0 = (\mu_0 / \epsilon_0)^{0.5}$ free-space impedance (377 ohms)

RADIATION: FAR FIELDS

Recall the Radiated Fields:

$$\bar{\mathbf{H}} = \hat{\phi} \left(jI_0 d e^{-jkr} / 2\lambda r \right) \left[1 + (jkr)^{-1} \right] \sin\theta$$

$$\bar{\mathbf{E}} = \left(jI_0 d e^{-jkr} / 2\lambda r \right) (\mu_0 / \epsilon_0)^{0.5}$$

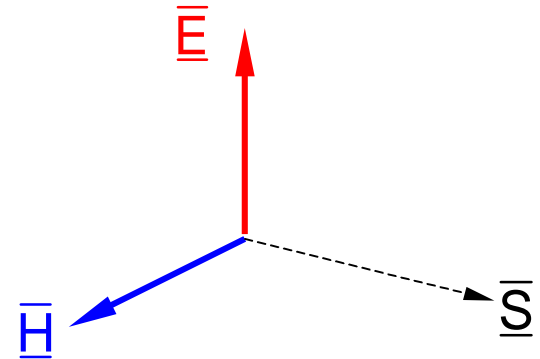
$$\left\{ \hat{\theta} \left[1 + (jkr)^{-1} + (jkr)^{-2} \right] \sin\theta + \hat{r} \left[(jkr)^{-1} + (jkr)^{-2} \right] 2\cos\theta \right\}$$

Far-Field Limit ($kr \gg 1$):

$$(r \gg \lambda / 2\pi)$$

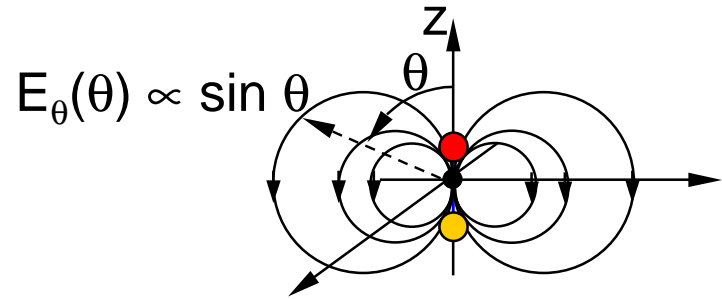
$$\bar{\mathbf{E}}_{ff} \cong \hat{\theta} \sin\theta \left(jI_0 d e^{-jkr} / 2\lambda r \right) (\mu_0 / \epsilon_0)^{0.5}$$

$$\bar{\mathbf{H}}_{ff} \cong \hat{\phi} \sin\theta \left(jI_0 d e^{-jkr} / 2\lambda r \right)$$



RADIATION: NEAR FIELDS

$kr \ll 1 \Rightarrow$ Static Dipole Fields:



$$\bar{\mathbf{E}} = (jI d e^{-jkr} / 2\lambda r) (\mu_0 / \epsilon_0)^{0.5} \left\{ \hat{\theta} \left[1 + (jkr)^{-1} + (jkr)^{-2} \right] \sin \theta + \hat{r} \left[(jkr)^{-1} + (jkr)^{-2} \right] 2 \cos \theta \right\}$$

$$\bar{\mathbf{E}} \cong (jI d e^{-jkr} / 2\lambda r) (\mu_0 / \epsilon_0)^{0.5} \left\{ \hat{\theta} (jkr)^{-2} \sin \theta + \hat{r} (jkr)^{-2} 2 \cos \theta \right\} \quad \text{near field}$$

$$\bar{\mathbf{H}} = \hat{\phi} (jI d e^{-jkr} / 4\pi r) \left[1 + (jkr)^{-1} \right] \sin \theta$$

$$\bar{\mathbf{H}} \cong \hat{\phi} (I d e^{-jkr} / 4\pi r^2) \sin \theta \quad \text{near field}$$

Imaginary Power:

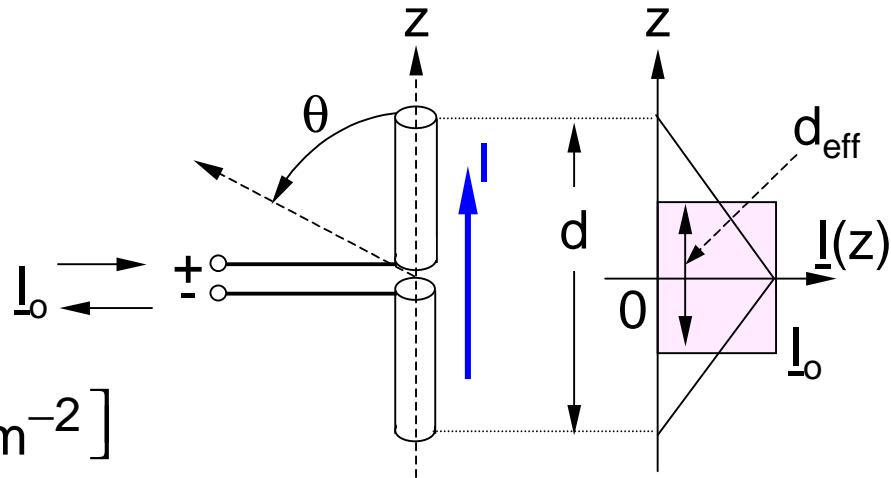
$$\bar{\mathbf{S}} = \bar{\mathbf{E}} \times \bar{\mathbf{H}}^* = -j \{ \bullet \} \quad (\text{average net stored energy is electric})$$

SHORT DIPOLE ANTENNAS

Currents on Short Dipole:

$$\int_{V_q} \bar{J} dv_q = \int_{V_q} \bar{I}(z) dz$$

$$\int_{V_q} \bar{J} dv_q = \int_{V_q} \bar{I}(z) dz = I_0 d_{\text{eff}} \cong I_0 d/2$$



Total Power Radiated:

$$P(\theta, \phi, r) = (\eta_0/2) |I_0 d_{\text{eff}} / 2\lambda r|^2 \sin^2 \theta \text{ [Wm}^{-2}\text{]}$$

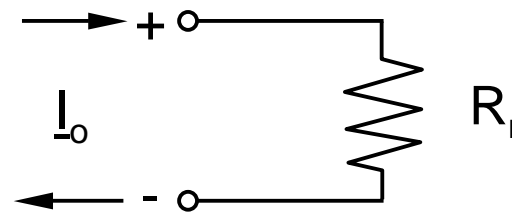
$$P_R = \int_0^{\pi} \int_0^{2\pi} P(\theta, \phi, r) (r d\phi \sin \theta) r d\theta = |k I_0 d_{\text{eff}}|^2 \eta_0 / 12\pi \text{ [W]}$$

Radiation Resistance R_r :

$$P_R = |I_0|^2 R_r / 2 \text{ [W]}$$

$$R_r = (kd)^2 \eta_0 / 6\pi \text{ ohms}$$

$$= (2\pi\eta_0/3) (d_{\text{eff}}/\lambda)^2$$

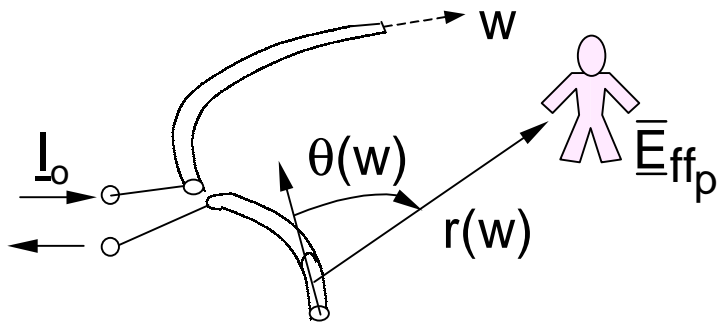


Equivalent
resistance

Example: 1-MHz AM signal, $\lambda = 300 \text{ m}$
 $d_{\text{eff}} = 1 \text{ m} \Rightarrow R_r \cong 0.1 \text{ ohm}$ (badly mismatched)

WIRE ANTENNAS

Sum of Short Dipole Elements:



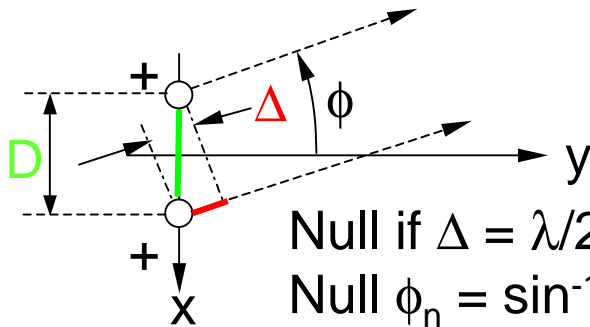
$$\bar{E}_{ff} d_{eff} \cong \hat{\theta} \sin \theta (j I_0 d_{eff} e^{-jk r} / 2 \lambda r) \eta_0 d_{eff}$$

$$\bar{E}_{ff_p} = \int_W \bar{E}_{ff_p}(w) dw$$

where $I, r,$ and θ vary with w

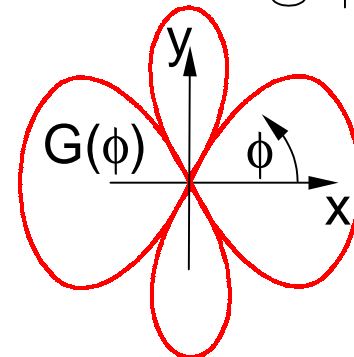
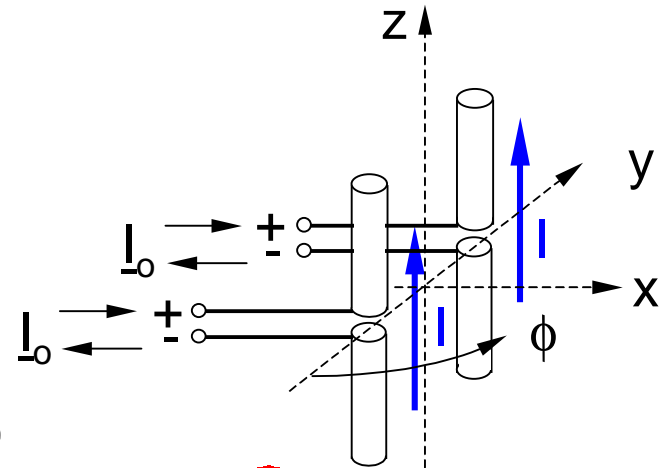
Superposition applies

Radiation from a Pair of Dipoles:



Null if $\Delta = \lambda/2 + n\lambda$ ($n = 0, 1, 2, \dots$)

Null $\phi_n = \sin^{-1} \Delta/D = \sin^{-1} \lambda/2D$



Example, $D = \lambda$: $\phi_n = \sin^{-1} [(\lambda/2)/\lambda] = 60^\circ$