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6.641 Electromagnetic Fields, Forces, and Motion  
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## Quiz 1 - Solutions 2003

## Problem 1

A

**Question:** What is the potential distribution,  $\Phi(r=0, z)$ , along the  $z$  axis for  $z > b$ ?

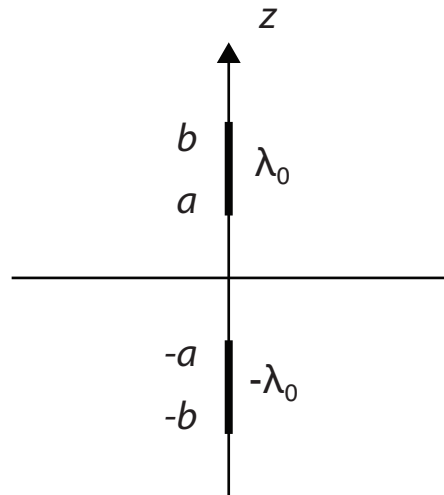


Figure 1: Line charge  $\lambda_0$  above  $z=0$  ground plane and image line charge  $-\lambda_0$  (Image by MIT OpenCourseWare.)

Image line charge  $-\lambda_0$  extends from  $-a < z < -b$ . Take charge elements  $\pm\lambda_0 dz'$  at  $z = z'$  and  $z = -z'$ . For  $z > b$ ,  $z > z'$ .

**Solution:**

$$\begin{aligned}
 d\Phi(r=0, z) &= \frac{\lambda_0}{4\pi\epsilon_0} \left[ \frac{1}{z-z'} - \frac{1}{z+z'} \right] dz' \\
 \Phi(r=0, z) &= \frac{\lambda_0}{4\pi\epsilon_0} \int_a^b \left[ \frac{1}{z-z'} - \frac{1}{z+z'} \right] dz' \\
 &= \frac{\lambda_0}{4\pi\epsilon_0} \left[ -\ln(z'-z) - \ln(z'+z) \right] \Big|_{z'=a}^b \\
 &= \frac{\lambda_0}{4\pi\epsilon_0} \left[ -\ln(b-z) - \ln(b+z) + \ln(a-z) + \ln(a+z) \right] \\
 &= \frac{\lambda_0}{4\pi\epsilon_0} \ln \left[ \frac{(a-z)(a+z)}{(b-z)(b+z)} \right] = \frac{\lambda_0}{4\pi\epsilon_0} \ln \left[ \frac{a^2 - z^2}{b^2 - z^2} \right]
 \end{aligned}$$

**B**

**Question:** What is the electric field,  $\vec{E}(r=0, z)$ , along the  $z$  axis for  $z > b$ ?

**Solution:**

$$\begin{aligned} E_z(r=0, z) &= -\frac{\partial\Phi(r=0, z)}{\partial z} = \frac{\lambda_0}{4\pi\epsilon_0} \left[ -\frac{1}{b-z} + \frac{1}{b+z} + \frac{1}{a-z} - \frac{1}{a+z} \right] \\ &= \frac{\lambda_0 z}{2\pi\epsilon_0} \frac{(b^2 - a^2)}{(a^2 - z^2)(b^2 - z^2)} \end{aligned}$$

**C**

**Question:** What is the free surface charge density,  $\sigma_S(r, z=0)$ , as a function of radial distance  $r$  on the  $z=0$  ground plane? **Hint:** One or more of these integrals may be useful for solving this problem.

$$\begin{aligned} \int \frac{x dx}{[x^2 + c^2]^{\frac{3}{2}}} &= -\frac{1}{\sqrt{x^2 + c^2}}; \int \frac{dx}{[x^2 + c^2]^{\frac{3}{2}}} = -\frac{x}{c^2\sqrt{x^2 + c^2}} \\ \int \frac{x dx}{\sqrt{x^2 + c^2}} &= \sqrt{x^2 + c^2}; \int \frac{dx}{x\sqrt{x^2 + c^2}} = -\frac{1}{c} \ln \left[ \frac{c + \sqrt{x^2 + c^2}}{x} \right] \end{aligned}$$

**Solution:**

$$z = 0$$

$$dE_z = \frac{-\lambda_0 dz'}{4\pi\epsilon_0 [z'^2 + r^2]} 2 \cos \theta; \cos \theta = \frac{z'}{\sqrt{z'^2 + r^2}}$$

$$\begin{aligned} dE_z &= \frac{-\lambda_0 z' dz'}{2\pi\epsilon_0 [z'^2 + r^2]^{\frac{3}{2}}} \\ E_z &= \frac{-\lambda_0}{2\pi\epsilon_0} \int_a^b \frac{z' dz'}{[z'^2 + r^2]^{\frac{3}{2}}} \\ &= \frac{\lambda_0}{2\pi\epsilon_0} \frac{1}{\sqrt{z'^2 + r^2}} \Big|_{z'=a}^b \\ &= \frac{\lambda_0}{2\pi\epsilon_0} \left[ \frac{1}{\sqrt{b^2 + r^2}} - \frac{1}{\sqrt{a^2 + r^2}} \right] \\ \sigma_s(z=0) &= \epsilon_0 E_z(z=0_+) = \frac{\lambda_0}{2\pi} \left[ \frac{1}{\sqrt{b^2 + r^2}} - \frac{1}{\sqrt{a^2 + r^2}} \right] \end{aligned}$$

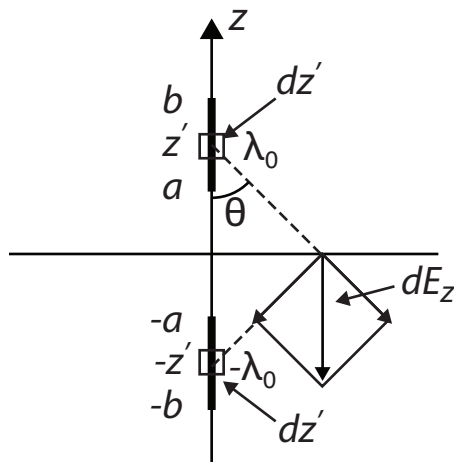


Figure 2: By symmetry each charge element and its image causes a  $-z$  directed electric field. (Image by MIT OpenCourseWare.)

## Problem 2

A

**Question:** What is the volume charge distribution for  $0 < x < s$  as a function of time?

**Solution:**

$$\rho_f(t) = \rho_0 e^{-\frac{t}{\tau}}, \tau = \frac{\epsilon}{\sigma}$$

B

**Question:** What is the electric field  $E_x(x, t)$ ?

**Solution:**

$$\begin{aligned} \nabla \cdot \vec{E} &= \frac{\partial E_x}{\partial x} = \frac{\rho_f(t)}{\epsilon} \Rightarrow E_x = \frac{\rho_f(t)x}{\epsilon} + c(t) \\ \int_0^s E_x dx &= \frac{\rho_f(t)s^2}{2\epsilon} + c(t)s = 0 \\ c(t) &= -\frac{\rho_f(t)s}{2\epsilon} \\ E_x &= \frac{\rho_f(t)}{\epsilon} \left( x - \frac{s}{2} \right) \end{aligned}$$

C

**Question:** What are the surface charge densities as a function of time at  $x = 0$  and  $x = s$ ?

**Solution:**

$$\sigma_s(x=0) = \epsilon E_x(x=0) = \frac{-\rho_f(t)s}{2}$$

$$\sigma_s(x=s) = -\epsilon E_x(x=s) = \frac{-\rho_f(t)s}{2}$$

**D**

**Question:** What is the current  $i(t)$  flowing through the short circuit?

**Solution:**

$$\begin{aligned} \frac{i(t)}{A} &= \sigma E_x(x=s) + \epsilon \left. \frac{\partial E_x}{\partial t} \right|_{x=s} \\ &= \frac{\sigma \rho_f(t)s}{2\epsilon} + \frac{\cancel{s} \partial \rho_f}{2\cancel{s} \partial t} \\ &= \frac{\sigma s}{2\epsilon} \rho_0 e^{-\frac{t}{\tau}} + \frac{s}{2} \left( -\frac{1}{\tau} \right) \rho_0 e^{-\frac{t}{\tau}} \\ &= \frac{\sigma s}{2\epsilon} \rho_0 e^{-\frac{t}{\tau}} - \frac{s\sigma}{2\epsilon} \rho_0 e^{-\frac{t}{\tau}} \\ &= 0 \end{aligned}$$

**Problem 3****A**

**Question:** Consider first the system shown in Figure 3.1 (see questions). In this system, a surface current exists only in the plane  $z=0$  and is everywhere surrounded by free space. For  $x > 0$ , the surface current is given by  $\vec{K} = K_0 \hat{i}_y$ . For  $x < 0$ , the surface current is given by  $\vec{K} = -K_0 \hat{i}_y$ . Which components of the magnetic field  $\vec{H}$  do you expect to be non-zero in this system? On which coordinates ( $x$  and/or  $y$  and/or  $z$ ) do you expect the non-zero components to depend? Explain briefly.

**Solution:**  $\vec{J}$  is only in the  $\hat{y}$  direction so therefore  $\vec{A}$  is only in the  $\hat{y}$  direction. The system is symmetric in  $y$ , so only  $x$  and  $z$  dependencies are expected. Therefore,  $\vec{A} = A_y(x, z) \hat{y}$ . Following  $\vec{B} = \mu_0 \vec{H} = \nabla \times \vec{A}$ ,  $\vec{B} = \hat{z} \frac{\partial A_y}{\partial x} - \hat{x} \frac{\partial A_y}{\partial z}$ . In summary,  $\vec{B} = B_x(x, z) \hat{x} + B_z(x, z) \hat{z}$ .

**B**

**Question:** Derive an integral expression for the magnetic field  $\vec{H}$  but do not evaluate it.

**Solution:**

$$\vec{H}(\vec{r}) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\vec{K}(\vec{r}') \times (\vec{r} - \vec{r}')}{4\pi |\vec{r} - \vec{r}'|^3} dx' dy'$$

$$\vec{K}(\vec{r}') = \begin{cases} K_0 \hat{y} & \text{for } x' > 0, \\ -K_0 \hat{y} & \text{for } x' < 0 \end{cases}$$

$$\begin{aligned}\bar{r} &= x\hat{x} + z\hat{z} \quad (y \text{ dependence is not necessary}) \\ \bar{r}' &= x'\hat{x} + y'\hat{y} \quad (\text{integration only in the } x - y \text{ plane})\end{aligned}$$

$$\begin{aligned}\bar{H}(\bar{r}) &= \int_0^\infty \int_{-\infty}^\infty \frac{K_0 \hat{y} \times ((x-x')\hat{x} - y'\hat{y} + z\hat{z})}{4\pi [(x-x')^2 + (y')^2 + (z)^2]^{\frac{3}{2}}} dy' dx' - \int_{-\infty}^0 \int_{-\infty}^\infty \frac{K_0 \hat{y} \times ((x-x')\hat{x} - y'\hat{y} + z\hat{z})}{4\pi [(x-x')^2 + (y')^2 + (z)^2]^{\frac{3}{2}}} dy' dx' \\ \bar{H}(x, z) &= \int_0^\infty \int_{-\infty}^\infty \frac{K_0 z \hat{x} - K_0 (x-x') \hat{z}}{2\pi [(x-x')^2 + (y')^2 + (z)^2]^{\frac{3}{2}}} dy' dx'\end{aligned}$$

### C

**Question:** Consider now the system shown in Figure 3.2 (see questions). In this system, a surface current exists only in the plane  $z = d$  above a perfect conductor that occupies the infinite half space  $z < 0$ . Free space occupies the infinite half-space  $z > 0$ . For  $x > L$ , the surface current is given by  $\vec{K} = K_0 \hat{i}_y$ . For  $x < -L$ , the surface current is given by  $\vec{K} = -K_0 \hat{i}_y$ . Writing the answer to Part (b) as  $\bar{H}^*(x, y, z)$ , which of the possible answers given below is the solution  $\bar{H}(x, y, z)$  for the configuration shown in Figure 3.2?

- i)  $\bar{H}(x, y, z) = [\bar{H}^*(x - L, y, z - d) + \bar{H}^*(x + L, y, z - d) + \bar{H}^*(x - L, y, z + d) + \bar{H}^*(x + L, y, z + d)]$
- ii)  $\bar{H}(x, y, z) = \frac{1}{2} [\bar{H}^*(x - L, y, z - d) + \bar{H}^*(x + L, y, z - d) - \bar{H}^*(x - L, y, z + d) - \bar{H}^*(x + L, y, z + d)]$
- iii)  $\bar{H}(x, y, z) = \frac{1}{2} [\bar{H}^*(x - L, y, z - d) - \bar{H}^*(x + L, y, z + d)]$
- iv)  $\bar{H}(x, y, z) = [\bar{H}^*(x - L, y, z + d) + \bar{H}^*(x + L, y, z - d)]$

See questions for Figure 3.2

**Solution:** Let  $K^*$  be the distribution of the surface current in the  $\hat{y}$  direction at  $z = 0$  in Parts (A) and (B). Then, the distribution of the surface current in the  $\hat{y}$  direction at  $z = D$  in this part is given by  $K = \frac{1}{2} (K^*(x - L, y) + K^*(x + L, y))$ . Here, superposition is used to create the new surface current distribution. The resulting magnetic fields are similarly superimposed, after translation by  $D$  in the  $\hat{z}$  direction. A second set of superpositions is used to match the boundary condition at the perfect conductor following the method of images. Hence,

$$\bar{H} = \frac{1}{2} \left[ \underbrace{\bar{H}^*(x - L, z - D) + \bar{H}^*(x + L, z - D)}_{\text{Translated and superimposed original solutions}} - \underbrace{\bar{H}^*(x - L, z + D) - \bar{H}^*(x + L, z + D)}_{\text{Image solutions}} \right]$$

where  $\bar{H}^*$  is the magnetic field found in Part B.