

Ex. 1

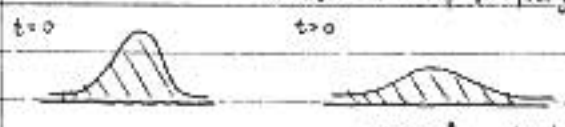
$$f(x) = \begin{cases} 1, & x > 0 \\ 0, & x < 0 \end{cases}$$

$$u(x,t) = \frac{1}{\sqrt{4\pi vt}} \int_{-\infty}^{\infty} d\xi e^{-\frac{(x-\xi)^2}{4vt}}$$

(after a change of variables)

x) We have discontinuous Initial data? Can this discontinuity propagate?
 No, the integral gives us a smooth function (u and derivatives)
 => discontinuity data can NOT propagate via diff. system.

x) No sense of finite speed of propagation in the relation. Solution appears to "propagate" everywhere instantaneously.
 ↳ don't really propagate, spread



spreads, but the area is the same

because diffusion equation comes from conservation law

$$\begin{cases} \frac{\partial q}{\partial t} + \frac{\partial q}{\partial x} = 0 \\ q = -v \nabla u \end{cases}$$

March 8, 2004 Lecture 10

Opt. Reading Debnath 1.1-1.9

Rev. Session FRI, 4:15-5:45

Quiz 1: MON, March 15

Diffusion eqn: IVP $\begin{cases} u_t = v u_{xx}, & -\infty < x < \infty, t > 0 \\ u(x, t=0) = f(x) \\ u(x,t) \rightarrow 0 \text{ as } |x| \rightarrow \infty \text{ "sufficiently fast"} \end{cases}$

FT: solution $u(x,t) = \int_{-\infty}^{\infty} dx' f(x') \frac{1}{\sqrt{4\pi vt}} e^{-\frac{(x-x')^2}{4vt}}$ ($u \rightarrow 0, u_x \rightarrow 0, u_{xx} \rightarrow 0$ etc.)

$G(x, x'; t, t') = \frac{1}{\sqrt{4\pi v(t-t')}} e^{-\frac{(x-x')^2}{4v(t-t')}}$, $t > t'$, Green's function (Green's function is also useful for non-linear equations)

Feature: Area is conserved $\int_{-\infty}^{\infty} dx u(x,t) = \text{const}$ Why?


© Diffusion eqn. comes from a conservation law, $g_t + q_x = 0$ ($g = u, q = -v u_x$)

$\rho_t + q_x = 0 \Rightarrow u_t = \nu u_{xx}$ (the structure of the PDE)

$\rho = u, q = -\nu u_x$

2. Condition for u as $|x| \rightarrow \infty$: $u_x \rightarrow 0$ as $|x| \rightarrow \infty$

this means that the flux $q = -\nu u_x \rightarrow 0$ as $|x| \rightarrow \infty$

 total mass inside the box $\int u dx$
 $\frac{d}{dt} (\text{total mass}) = \text{flux going out}$ but $q = 0$ as $|x| \rightarrow \infty$
 \Rightarrow mass conserved

Notion of conservation of total mass in diffusion equation:

$$\int_{-\infty}^{+\infty} dx u_t = \nu \int_{-\infty}^{+\infty} u_{xx} dx$$

$$\frac{d}{dt} \int_{-\infty}^{+\infty} u dx = \nu u_x \Big|_{-\infty}^{+\infty}$$

$$= \nu [u_x(+\infty, t) - u_x(-\infty, t)] = 0$$

$\Rightarrow \int_{-\infty}^{+\infty} dx u(x, t) = \text{const}$

The assumptions we used here

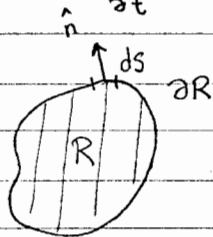
- (i) the structure of the PDE
- (ii) flux goes to zero at infinity

Let's see what happens if we use the general conservation law

$\frac{\partial \rho}{\partial t} + \frac{\partial q}{\partial x} = 0$ (1D conservation law)

Ex. In n -dimensions, the form of the conservation law is

$\frac{\partial \rho}{\partial t} + \nabla \cdot \vec{q} = 0$



- area R , boundary ∂R
- the conservation law is true inside R

Let's integrate over the area:

$\rho(\vec{r}, t), \vec{q}(\vec{r}, t)$

from divergence theorem

$\int_R \frac{\partial \rho}{\partial t} d\vec{r} + \int_R \nabla \cdot \vec{q} d\vec{r} = 0 \Rightarrow \frac{d}{dt} \int_R \rho d\vec{r} + \int_{\partial R} ds \hat{n} \cdot \vec{q} = 0$

$$\Rightarrow \underbrace{\frac{d}{dt} \int_R g d\vec{r}}_{\text{total mass in region } R} = - \int_{\partial R} ds \hat{n} \cdot \vec{q}$$

Specifically, let's assume $\vec{q} = 0 \mid_{\vec{r} \in \partial R}$

$\int_R e d\vec{r} = \text{const}$ - the diffusion equation is a special case (1D)

In n dimensions: \vec{r}
 \downarrow
 $x \in \mathbb{R}^n, t > 0$

$$\begin{cases} u_t = \nu \nabla^2 u \\ \text{IVP: } u(\vec{r}, t) = f(\vec{r}) \end{cases}$$

Solution (by Fourier transform in all space-variables)

$$u(\vec{r}, t) = \int_{\mathbb{R}^n} d\vec{r}' f(\vec{r}') \frac{1}{(4\pi\nu t)^{n/2}} e^{-\frac{(\vec{r}-\vec{r}')^2}{4\nu t}}$$

Next: Optimal Control Theory, Burger equation: nonlinear PDE can be converted by a nonlinear transformation to the diffusion equation.

(diffusion equation: in FT space, we have an ODE)

Linear PDE's

Nonlinear PDE's convertible to (Diffusion Equation)

$$\text{PDE: } u_t - a \nabla^2 u + b |\nabla u|^2 = 0 \quad \begin{cases} \text{n dimensions + time } t > 0 \\ \vec{r} = (x_1, \dots, x_n) \quad a, b = \text{const} \end{cases}$$

↑
nonlinear term

Question: Is there a change $w = \phi(u)$ that makes this transformation

Nonlinear PDE \Rightarrow a Linear PDE

Transformation: $w = \phi(u)$

- $w_t = \phi'(u) u_t$
- $\nabla w = \phi'(u) \nabla u$
- $\nabla^2 w = \nabla \cdot (\nabla w) = \nabla \cdot (\phi'(u) \nabla u) = \phi'(u) \nabla^2 u + \nabla(\phi'(u)) \cdot \nabla u$
 $= \phi'(u) \nabla^2 u + \phi''(u) |\nabla u|^2$

Apply the PDE:

$$w_t = \phi'(u) u_t = \phi'(u) [a \nabla^2 u - b |\nabla u|^2]$$

$$= a [\nabla^2 w - \phi''(u) |\nabla u|^2] - b \phi'(u) |\nabla u|^2$$

$$= a \nabla^2 w - \underbrace{(a \phi''(u) + b \phi'(u))}_{\text{we set this to zero (definition of } \phi(u) \text{)}} |\nabla u|^2$$

we set this to zero (definition of $\phi(u)$)
 \Rightarrow we will have a linear equation for w

Transformation:

$$a \phi''(u) + b \phi'(u) = 0 \quad (\text{linear ODE, constant coefficients})$$

solve for ϕ , any particular solution will do

$$\phi(u) = e^{-\frac{bu}{a}} = w \quad \begin{matrix} \text{Cole-Hopf transformation} \\ \text{(non-linear transformation)} \end{matrix}$$

PDE for w is the diffusion equation: $w_t = a \nabla^2 w$

$$\text{IVP: } \begin{cases} w_t - a \nabla^2 w + b |\nabla w|^2 = 0, & \vec{r} \in \mathbb{R}^n, t > 0 \\ w(\vec{r}, t=0) = g(\vec{r}) \end{cases}$$

$$\downarrow \begin{cases} w_t = a \nabla^2 w \\ w(\vec{r}, t=0) = e^{-\frac{bg(\vec{r})}{a}} \end{cases} \Rightarrow w(\vec{r}, t) = \int_{\mathbb{R}^n} d\vec{r}' e^{-\frac{bg(\vec{r}')}{a}} \frac{1}{(4\pi ut)^{n/2}} e^{-\frac{(\vec{r}-\vec{r}')^2}{4ut}}$$

$$u = -\frac{a}{b} \ln w = -\frac{a}{b} \ln \left\{ \int_{\mathbb{R}^n} d\vec{r}' e^{-\frac{bg(\vec{r}')}{a}} \frac{1}{(4\pi ut)^{n/2}} e^{-\frac{(\vec{r}-\vec{r}')^2}{4ut}} \right\}$$

What happens at infinity?

w is a solution of the linear diffusion equation and an assumption is that $w \rightarrow 0$ at $|\vec{x}| \rightarrow \infty$

$$w = e^{-\frac{b}{a} u} \Rightarrow \text{Solution for } u \text{ holds if } |u| \rightarrow \infty, |\vec{r}| \rightarrow \infty$$

($+\infty$ or $-\infty$ depends on the sign b/a)

⚠ Be aware that the transformation converts not only the equation but also the conditions.
The difficulty has been moved from the nonlinear PDE to the diffusion equation: solve it with different conditions at infinity.

Example: Burger's equation

c = c(x,t) c_t + cc_x = \nu c_{xx} , c(x,0) = g(x)

Recall:

conservation law g_t + q_x = 0 => g_t + \underbrace{Q'(g)}_{c(x,t)} g_x = \nu g_{xx}

assumed the flux q = Q(g) - \nu g_x

Next, we derived an equation for c

when Q(g) is a quadratic function Q(g) = dg^2 + \beta g + \gamma

This is the Burger's equation.

c_t + cc_x = \nu c_{xx} (1)

\rightarrow c_t + (\frac{1}{2}c^2)_x = \nu c_{xx}

introduce u : c = u_x

u(x,t) = \int_{-\infty}^x dx' c(x',t) assuming that

c \to 0 as |x| \to \infty

sufficiently fast

\Rightarrow (u_t)_x + (\frac{1}{2}u_x^2)_x = \nu (u_{xx})_x

integrate \int dx u_t + \frac{1}{2}u_x^2 = \nu u_{xx} + K(t)

take limit x \to -\infty , u(x,t) = \int_{-\infty}^x dx' c(x',t) = 0 => K(t) = 0

u_t - \nu u_{xx} + \frac{1}{2}(u_x)^2 = 0 (2) a = \nu b = \frac{1}{2} in 1D

Solution for u : we need initial conditions

IC for u : q = u_x , u = \int_{-\infty}^x dx' g(x') = h(x)

Solution for u by Cole-Hopf transformation

u(x,t) = -2\nu \ln \left\{ \frac{1}{\sqrt{4\nu t}} \int_{-\infty}^{+\infty} dx' e^{-\frac{h(x')}{2\nu}} e^{-\frac{(x-x')^2}{4\nu t}} \right\}

c(x,t) = +2\nu \times \left\{ \int_{-\infty}^{+\infty} dx' \frac{(x-x')}{2\nu t} e^{-\frac{h(x')}{2\nu}} e^{-\frac{(x-x')^2}{4\nu t}} \right\} \cdot \left\{ \int_{-\infty}^{+\infty} dx' e^{-\frac{h(x')}{2\nu}} e^{-\frac{(x-x')^2}{4\nu t}} \right\}^{-1}