Massachusetts Institute of Technology Department of Electrical Engineering and Computer Science 6.641 Electromagnetic Fields, Forces, and Motion, Spring 2005

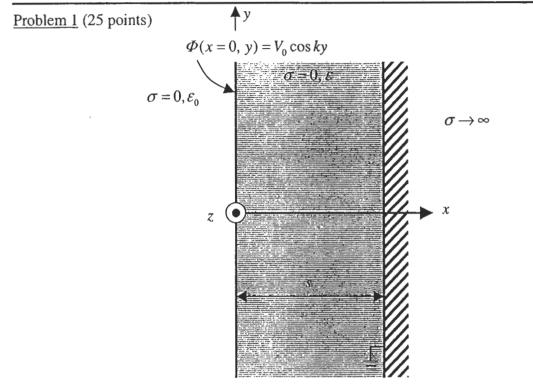
Final Review Packet, 4/26/05

Final Exam Spring 2004. 1:30-4:30PM

6.641 Formula Sheet Attached in the study materials section.

You are also allowed to use the formula sheets that you prepared for Quiz 1, Quiz 2,

and an additional 8 1/2"x 11" formula sheet (both sides) that you have prepared for the Final.



A potential sheet of infinite extent in the y and z directions is placed at x = 0 and has potential distribution $\Phi(x=0, y) = V_0 \cos ky$. Free space with no conductivity ($\sigma = 0$) and permittivity ε_0 is present for x < 0 while for 0 < x < s a perfectly insulating dielectric ($\sigma = 0$) with permittivity ε is present. The region for x > s is a grounded perfect conductor at zero potential.

a) What are the potential distributions for x < 0 and 0 < x < s?

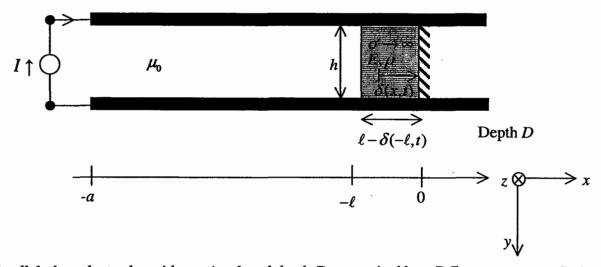
b) What are the surface charge densities at x = 0, $\sigma_f(x = 0, y)$, and at x = s, $\sigma_f(x = s, y)$? $E_x = -\frac{2\pi}{3x} = \begin{cases} -\lambda_e \ V_o \ coskey \ e^{\lambda_e x} \ x = 0 \end{cases} \qquad \nabla_f(x = 0, y), \text{ and at } x = s, \sigma_f(x = s, y)?$ $= \begin{bmatrix} -\lambda_e \ V_o \ coskey \ e^{\lambda_e x} \ x = 0 \end{cases} \qquad \nabla_f(x = 0, y), \text{ and at } x = s, \sigma_f(x = s, y)?$ $= \begin{bmatrix} -\lambda_e \ V_o \ coskey \ e^{\lambda_e x} \ x = 0 \end{cases} \qquad \nabla_f(x = 0, y), \text{ and at } x = s, \sigma_f(x = s, y)?$ $= \begin{bmatrix} -\lambda_e \ V_o \ coskey \ e^{\lambda_e x} \ x = 0 \end{cases} \qquad \nabla_f(x = 0, y), \text{ and at } x = s, \sigma_f(x = s, y)?$ $= \begin{bmatrix} -\lambda_e \ V_o \ coskey \ e^{\lambda_e x} \ x = 0 \end{cases} \qquad \nabla_f(x = 0, y), \text{ and at } x = s, \sigma_f(x = s, y)?$ $= \begin{bmatrix} -\lambda_e \ V_o \ coskey \ e^{\lambda_e x} \ x = 0 \end{cases} \qquad \nabla_f(x = 0, y) = -E_x(x = s, y)$ c) What is the force, magnitude and direction, on a section of the perfect conductor at x = s

that extends over the region
$$0 < y < \frac{\pi}{k}$$
 and $0 < z < D$?

Hint:
$$\int \cos^2 y \, dy = y/2 + (\sin 2y)/4$$

 $\frac{F_x}{ama} = \frac{1}{2} \sqrt{\frac{1}{2}} E_x|_{x=5} = -\frac{1}{2} \frac{1}{2} \frac{1}{2$

May, 17 2004

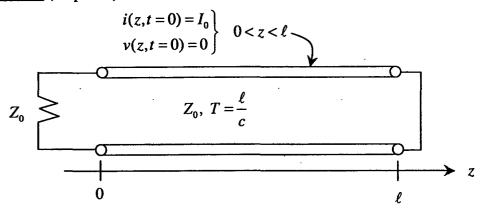


Parallel plate electrodes with spacing h and depth D are excited by a DC current source I. An elastic rod surrounded by free space has mass density ρ , modulus of elasticity E, equilibrium length ℓ and has infinite ohmic conductivity σ . The elastic rod end at x = 0 is fixed while the deflections of the rod are described as $\delta(x,t)$ and are assumed small $|\delta(x,t)| \ll \ell$. The rod width $\ell - \delta(-\ell,t)$ changes as I is changed because of the magnetic force. The DC current flows as a surface current on the $x = -(\ell - \delta(-\ell, t))$ end of the perfectly conducting rod.

- a) Calculate H_z in the free space region $-a < x < -(\ell \delta(-\ell, t))$. Neglect fringing field effects and assume $h \ll a$ and $h \ll D$. $H_z = \frac{T}{D}$
- b) Using the Maxwell Stress Tensor calculate the magnetic force per unit area on the $x = -(\ell - \delta(-\ell, t))$ end of the rod. $T_{xx} = \frac{1}{2} \Lambda_0 \left[H_x^2 - H_y^2 - H_z^2 \right] \left[\begin{array}{c} = -\frac{\mu_0}{2} & \frac{1}{2}^2 \\ x = -(\ell - \delta(-\ell, t)) \end{array} \right]$ $f_{x} = -T_{xx} \left[\begin{array}{c} = -\frac{\mu_0}{2} & \frac{1}{2}^2 \\ x = -(\ell - \delta(-\ell_1, t)) \end{array} \right]$ $f_{x} = -T_{xx} \left[\begin{array}{c} = -\frac{\mu_0}{2} & \frac{1}{2}^2 \\ x = -(\ell - \delta(-\ell_1, t)) \end{array} \right]$ $f_{x} = -T_{xx} \left[\begin{array}{c} = -\frac{\mu_0}{2} & \frac{1}{2}^2 \\ x = -(\ell - \delta(-\ell_1, t)) \end{array} \right]$ $f_{x} = -T_{xx} \left[\begin{array}{c} = -\mu_0 & \frac{1}{2}^2 \\ x = -(\ell - \delta(-\ell_1, t)) \end{array} \right]$ $f_{x} = -\frac{\mu_0 + 1}{2} \\ f_{x} = -\frac{\mu_0$

$$\begin{aligned} \frac{\partial S}{\partial t^{2}} &= \mathcal{E} \frac{\partial S}{\partial x^{2}} \quad S'(x,t) = \frac{1}{2e} \left[\hat{S}(x) e^{i k t t} \right] \\ - \frac{\partial U}{\partial t^{2}} \hat{S}(x) &= \mathcal{E} \frac{\partial^{2} S}{\partial x^{2}} \quad \frac{\partial^{2} S}{\partial x^{2}} \quad + \frac{\partial^{2} S}{\partial x^{2}} \quad + \frac{\partial^{2} S}{\partial x^{2}} = 0 \quad \frac{\partial^{2} = \omega^{2} \rho}{\partial x^{2}} \rho \left[\mathcal{E} \right] \\ & \omega_{n} = \int \frac{\mathcal{E}}{\mathcal{E}} \frac{T}{2k} (2n+k), \quad h = 0, 1, 2 \\ \hat{S}(x) &= A \sin kx + B \cosh kx \\ \hat{S}(x=0) &= B = 0 \\ \frac{\partial S}{\partial x} \left[\begin{array}{c} z &= 0 \\ z &= -kA \cosh k \end{array} \right] = \sum \left[ke \left[(2n+1) \frac{T}{2}, \quad h > 0, 1, 2, 1 \right] \right] \\ & 2 \end{aligned} \end{aligned}$$

 (\mathcal{O})



An electrical transmission line of length ℓ has characteristic impedance Z_0 . Electromagnetic waves can travel on the line at speed c, so that the time to travel one-way over the line length ℓ is $T = \ell/c$. The line is matched at z = 0 and is short circuited at $z = \ell$. At time t = 0, a lightning bolt strikes the entire line so that there is a uniform current along the line but with zero voltage:

$$\begin{array}{l} i(z,t=0) = I_0 \\ v(z,t=0) = 0 \end{array} \} \quad 0 < z < \ell$$

Since the voltage and current obey the telegrapher's relations:

$$\frac{\partial v}{\partial z} = -L\frac{\partial i}{\partial t}, \quad c = \frac{1}{\sqrt{LC}}$$
$$\frac{\partial i}{\partial z} = -C\frac{\partial v}{\partial t}, \quad Z_0 = \sqrt{L/C}$$

the voltage and current along the line are related as

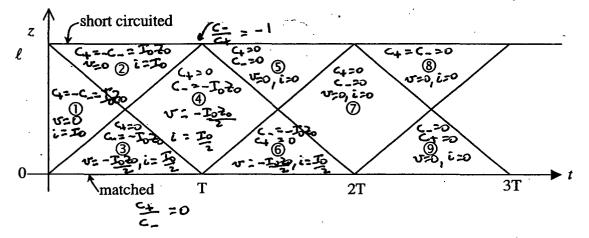
$$v + iZ_0 = c_+ \text{ on } \frac{dz}{dt} = c$$

$$v - iZ_0 = c_- \text{ on } \frac{dz}{dt} = -c$$

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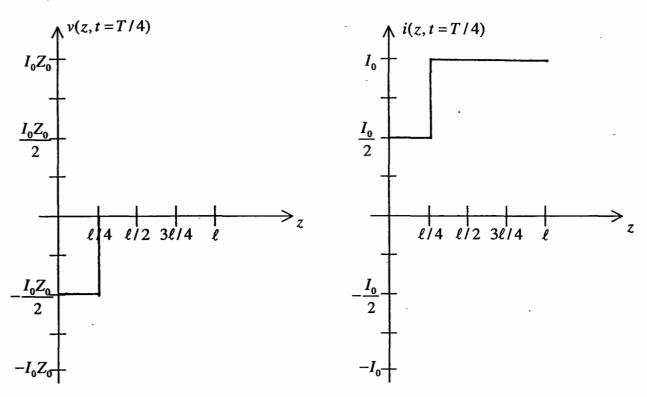
$$v - iZ_0 = c_- \text{ on } \frac{dz}{dt} = -c$$

a) The solutions for v(z,t) and i(z,t) can be found using the method of characteristics within each region shown below. Within regions 1-9 give the values of c_+, c_-, v and iZ_0 .



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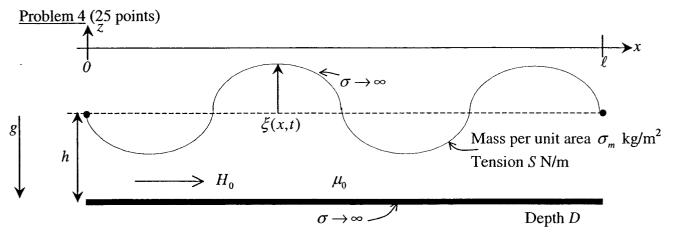
b) Plot v(z, t = T/4) and i(z, t = T/4).



c) How long a time does it take for the transmission line to have v(z, t) = 0 and i(z, t) = 0 everywhere for $0 < z < \ell$ for all further time? $\geq T$

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A perfectly conducting membrane of depth D with mass per unit area σ_m and tension S is a distance h above a rigid perfect conductor. The membrane and rigid conductor are in free space and support currents such that when the membrane is flat, $\xi(x,t) = 0$, the static uniform magnetic field intensity is H_0 . As the membrane deforms, the flux through the region between membrane and rigid conductor is conserved. The system is in a downward gravity field with gravitational acceleration $\overline{g} = -g\overline{i_z}$. The membrane deflection has no dependence on y and is fixed at its two ends at x = 0 and $x = \ell$.

a) Assuming that $\xi(x,t) \ll h$ and that the only significant magnetic field component is x directed, how is $H_x(x,t)$ approximately related to $\xi(x,t)$ to linear terms in $\xi(x,t)$?

$$H_{x}(h+3) = H_{0}h \Rightarrow H_{x} = \frac{H_{0}h}{h+3} = \frac{H_{0}}{1+3}h \approx H_{0}(1-3/h)$$

b) Using the Maxwell Stress tensor and the result of part (a), to linear terms in small displacement $\xi(x,t)$, what is the z directed magnetic force per unit area, F_z , on the

$$\begin{array}{l} \text{membrane}?\\ T_{22} = \frac{A_{00}}{2} \left(\frac{1}{42} - \frac{1}{4} \frac{2}{x} - \frac{1}{4} \frac{2}{y} \right) = -\frac{A_{00}}{2} \frac{1}{40} \left(1 - \frac{2}{3} \frac{2}{y} \right)^{2} \approx -\frac{A_{00}}{2} \frac{1}{40} \left(1 - \frac{2}{3} \frac{2}{y} \right) \\ F_{2} = -T_{22} = \frac{A_{00}}{2} \frac{1}{40} \left(1 - \frac{2}{3} \frac{2}{y} \right) \end{array}$$

c) To linear terms in small displacement $\xi(x,t)$, express the membrane equation of motion in the form

$$a\frac{\partial^{2}\xi}{\partial t^{2}} = b\frac{\partial^{2}\xi}{\partial x^{2}} + c\xi + d$$
What are a, b, c, and d?
$$a\frac{\partial^{2}\xi}{\partial t^{2}} = b\frac{\partial^{2}\xi}{\partial x^{2}} + c\xi + d$$

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d) What value of H_0 is needed so that in static equilibrium, the membrane has no sag,

$$\xi(x,t) = 0. \qquad \text{is of } \xi(x,t) = 0. \qquad \text{is of } \xi(x,$$

Continue to next page for parts (e)-(g)

(5)

e) About the equilibrium of part (d), what is the $\omega - k$ dispersion relation for membrane deflections of the form

 $\xi(x,t) = \operatorname{Re}[\hat{\xi}e^{j(\alpha t - kx)}]?$

Solve for k as a function of ω and system parameters.

$$J_{m}\omega^{2} = -Sk^{2} - \mu_{0}H_{0}^{2}$$

$$k^{2} = J_{m}\omega^{2} - \mu_{0}H_{0}^{2}$$

$$k^{2} = \pm \left(J_{m}\omega^{2} - \mu_{0}H_{0}^{2}\right)^{2}$$

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f) Using all the values of k found in part (e), find a superposition of solutions of the form of $\xi(x,t)$ given in (e), that satisfy the zero deflection boundary conditions at the ends of the membrane at x = 0 and $x = \ell$. What are the allowed values for k?

$$\begin{split} \vec{s}(x,t) &= Re \left[e^{j\omega t} \left[\vec{s}_{1} e^{-jk_{0}x} + \vec{s}_{2} e^{+jk_{0}x} \right] \right] \\ \vec{s}(x;0,t) &= 0 = Re \left[e^{j\omega t} \left[\vec{s}_{1} + \vec{s}_{2} \right] \right] \\ \vec{s}(x;0,t) &= Re \left[e^{j\omega t} \left[e^{-jk_{0}x} - e^{jk_{0}x} \right] \right] \\ \vec{s}(x;t) &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ \vec{s}(x;t) &= 0 = Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ \vec{s}(x;t) &= 0 = Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ \vec{s}(x;t) &= 0 = Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{1}(-2i) \sin k_{0}x \right] \\ &= Re \left[e^{j\omega t} \vec{s}_{$$

g) Is this system always stable or under what conditions can it be unstable? When stable, what are the natural frequencies and if unstable what are the growth rates of the instability?

$$\omega_{n}^{2} = \frac{S}{\sqrt{m}} k^{2} + \frac{\kappa_{0}H_{0}^{2}}{\sqrt{m}} + \frac{\kappa_{0}H_{0}^{2}}{\sqrt{m}} + \frac{\kappa_{0}H_{0}^{2}}{\sqrt{m}} - \frac{1}{2} \frac{1}{2}$$

$$\omega_{n} = \left[\frac{S}{\sqrt{m}} \left(\frac{\pi_{1}}{2}\right)^{2} + \frac{\kappa_{0}H_{0}^{2}}{\sqrt{m}} - \frac{1}{2}\right]$$

(6)