MIT Course 8.033, Fall 2005, Relativistic dynamics Max Tegmark Last revised October 25 2005

Topics

- Formula summary
- Momentum & energy
- Acceleration & force (optional)
- Transformation of force (optional)
- Transformation of acceleration (optional)

Dynamics toolbox: formula summary

• Mass-energy unification:

$$E = mc^2 = m_0 \gamma c^2$$

• Momentum 4-vector:

$$\mathbf{P} \equiv m_0 \mathbf{U} = \begin{pmatrix} p_x \\ p_y \\ p_z \\ E/c \end{pmatrix}$$

• Energy formula:

$$E = \sqrt{(m_0 c^2)^2 + (cp)^2}$$

• Velocity formula:

$$\beta = \frac{cp}{E}$$

- Optional material:
 - Acceleration 4-vector:

$$\mathbf{A} \equiv \frac{d\mathbf{U}}{d\tau} = \gamma_u^2 \begin{pmatrix} \mathbf{a} \\ 0 \end{pmatrix} + \gamma_u^4 \frac{\mathbf{u} \cdot \mathbf{a}}{c^2} \begin{pmatrix} \mathbf{u} \\ c \end{pmatrix}$$

- Force 4-vector:

$$\mathbb{F} \equiv \frac{d}{d\tau} \mathbf{P} = \gamma_u \begin{pmatrix} \mathbf{F} \\ P/c \end{pmatrix} = m_0 \mathbf{A}$$

- Power:

$$P = \dot{E} = \mathbf{u} \cdot \mathbf{F} = m_0 \gamma_u^3 \, \mathbf{u} \cdot \mathbf{a}$$

- Force 3-vector:

$$\frac{\mathbf{F}}{m_0 \gamma_u} = \mathbf{a} + \gamma_u^2 \frac{\mathbf{u} \cdot \mathbf{a}}{c^2} \mathbf{u} = \begin{cases} \gamma_u^2 \mathbf{a} & (\mathbf{u} \& \mathbf{a} \text{ parallel}) \\ \mathbf{a} & (\mathbf{u} \& \mathbf{a} \text{ perpendicular}) \end{cases}$$

- Acceleration 3-vector:

$$m\mathbf{a} = \mathbf{F} - \frac{P\mathbf{u}}{c^2}$$

– Force transformation:

$$F'_{x} = \frac{F_{x} - \frac{v}{c^{2}}P}{1 - \frac{u_{x}v}{c^{2}}},$$

$$F'_{y} = \frac{F_{y}}{\gamma\left(1 - \frac{u_{x}v}{c^{2}}\right)},$$

$$F'_{z} = \frac{F_{z}}{\gamma\left(1 - \frac{u_{x}v}{c^{2}}\right)},$$

$$P' = \frac{P - vF_{x}}{1 - \frac{u_{x}v}{c^{2}}}$$

Momentum & energy toolbox:

• Relativistic mass:

 $m = \gamma m_0$

• Mass-energy unification:

$$E = mc^2$$

• Momentum 4-vector (momentum-energy unification):

$$\mathbf{P} \equiv m_0 \mathbf{U} = m_0 \frac{d\mathbf{X}}{d\tau} = m_0 \gamma_u \begin{pmatrix} u_x \\ u_y \\ u_z \\ c \end{pmatrix} = m \begin{pmatrix} u_x \\ u_y \\ u_z \\ c \end{pmatrix} = \begin{pmatrix} p_x \\ p_y \\ p_z \\ E/c \end{pmatrix},$$

(Use upper case \mathbf{X} , \mathbf{U} and \mathbf{P} for the 4-vectors to avoid confusion with the \mathbf{x} , \mathbf{u} and \mathbf{p} 3-vectors.)

• Handy velocity formula follows straight from this:

$$\beta = \frac{cp}{E}$$

• Rest energy:

$$E_0 = m_0 c^2$$

is total energy of particle in the frame where it is at rest

• Kinetic enery:

$$K = E - E_0 = mc^2 - m_0 c^2 = m_0 c^2 (\gamma - 1) = \frac{1}{2} m_0 u^2 + O\left(\frac{u^4}{c}\right)$$

• Rest mass invariant:

$$m_0 = \frac{1}{c}\sqrt{-\mathbf{P}^t\boldsymbol{\eta}\mathbf{P}} = \frac{1}{c^2}\sqrt{E^2 - c^2p^2},$$

giving the handy relations

$$E = \sqrt{(m_0 c^2)^2 + (cp)^2},$$
$$p \equiv |\mathbf{p}| = \sqrt{\frac{E^2}{c^2} - (m_0 c)^2}.$$

• Low-speed limit $|\beta| \ll 1$:

$$E \approx m_0 c^2 + \frac{1}{2} m_0 u^2,$$
$$p = m_0 \gamma u \approx m_0 u.$$

• High-speed limit $|\beta| \approx 1 \ (\gamma \gg 1, E \gg E_0)$:

$$E \approx cp$$

This becomes exact (E = cp) for particles moving with speed of light, like photons and gravitons.

• $-\mathbf{P}^t \eta \mathbf{P} = (E/c)^2 - p^2$ is invariant also for *system* of particles, since

$$\mathbf{P}_{\mathrm{tot}}' \equiv \sum_i \mathbf{P}'_i = \sum_i \mathbf{\Lambda} \mathbf{P}_i = \mathbf{\Lambda} \left(\sum_i \mathbf{P}_i \right) = \mathbf{\Lambda} \mathbf{P}_{\mathrm{tot}}.$$

• We derived $\mathbf{p} = m_0 \gamma u$ only for 1-dimensional collision. But any collision is 1-dimensional in the frame where the total momentum is zero!

Acceleration & force (optional!)

- The acceleration 4-vector ${\bf A}$ and the Force 4-vector ${\mathbb F}$ are less useful than their 4-vector cousins \mathbf{X} , \mathbf{U} , \mathbf{P} and \mathbf{K} . We'll use \mathbb{F} mainly for deriving the force transformation law, which will in turn give us the transformation law for electromagnetic fields. We'll use upper case ${\bf A}$ for the acceleration 4-vector to avoid confusion with the the acceleration 3-vector \mathbf{a} , and the annoying symbol \mathbb{F} for the force 4-vector to avoid confusion with the the force 3-vector \mathbf{F} .
- Acceleration 4-vector:

$$\mathbf{A} \equiv \frac{d\mathbf{U}}{d\tau} = \gamma_u \frac{d\mathbf{U}}{dt} = \gamma_u \frac{d}{dt} \gamma_u \begin{pmatrix} u_x \\ u_y \\ u_z \\ c \end{pmatrix} = \gamma_u \frac{d}{dt} \gamma_u \begin{pmatrix} \mathbf{u} \\ c \end{pmatrix}$$
$$= \gamma_u^2 \begin{pmatrix} \dot{\mathbf{u}} \\ 0 \end{pmatrix} + \gamma_u \dot{\gamma}_u \begin{pmatrix} \mathbf{u} \\ c \end{pmatrix} = \gamma_u \begin{pmatrix} \mathbf{a} + \dot{\gamma}_u \mathbf{u} \\ \dot{\gamma}_u c \end{pmatrix}$$
$$= \gamma_u^2 \begin{pmatrix} \mathbf{a} \\ 0 \end{pmatrix} + \gamma_u^4 \frac{\mathbf{u} \cdot \mathbf{a}}{c^2} \begin{pmatrix} \mathbf{u} \\ c \end{pmatrix},$$

,

where in the last step, we have used the fact that

$$\dot{\gamma}_u = \gamma_u^3 \frac{\mathbf{u} \cdot \mathbf{a}}{c^2}.$$

• Force 4-vector:

$$\mathbb{F} \equiv \frac{d}{d\tau} \mathbf{P} = \gamma_u \frac{d}{dt} \mathbf{P} = \gamma_u \frac{d}{dt} m_0 \mathbf{U} = m_0 \frac{d}{d\tau} \mathbf{U},$$

so by definition, we have

$$\mathbb{F}=m_0\mathbf{A}.$$

(Note that this does *not* apply the Newtonian result $\mathbf{F} = ma!$)

• Interpretation of Force 4-vector:

$$\mathbb{F} = \gamma_u \frac{d}{dt} \mathbf{P} = \gamma_u \begin{pmatrix} \dot{\mathbf{p}} \\ \dot{E}/c \end{pmatrix} = \gamma_u \begin{pmatrix} \mathbf{F} \\ P/c \end{pmatrix},$$

where $\mathbf{F} = \dot{p}$ is the familiar force 3-vector and $P = \dot{E}$ is the power, the energy change per unit time (in Watts).

• Work-energy theorem:

$$dE = \mathbf{F} \cdot d\mathbf{r} = \mathbf{F} \cdot \frac{d\mathbf{r}}{dt} dt = \mathbf{F} \cdot \mathbf{u} dt,$$

so the power satisfies

$$P = \dot{E} = \mathbf{u} \cdot \mathbf{F}.$$

• Force 3-vector explicitly: Dividing the above equation $\mathbb{F} = m_0 \mathbf{A}$ by γ_u gives

$$\frac{\mathbf{F}}{m_0\gamma_u} = \mathbf{a} + \gamma_u^2 \frac{\mathbf{u} \cdot \mathbf{a}}{c^2} \mathbf{u}.$$

• Special case where **u** and **a** are parallel, *e.g.*, for linear motion:

$$\frac{\mathbf{F}}{m_0\gamma_u} = \mathbf{a} + \gamma_u^2 \frac{u^2 \mathbf{a}}{c^2} = \left(1 + \gamma_u^2 \beta^2\right) \mathbf{a} = \gamma_u^2 \mathbf{a}.$$

• Special case where **u** and **a** are perpendicular, eg, for circular motion:

$$\frac{\mathbf{F}}{m_0 \gamma_u} = \mathbf{a}$$

- Note that in relativity, **F** and **a** are generally *not* parallel, but that they are parallel for these two special cases.
- Acceleration 3-vector explicitly:

$$\mathbf{a} = \frac{\mathbf{F}}{m_0 \gamma_u} - \frac{\mathbf{u} \cdot \mathbf{F}}{m_0 \gamma_u c^2} \mathbf{u} = \frac{\mathbf{F}}{m} - \frac{P}{mc^2} \mathbf{u}.$$

The last term (the departure from $\mathbf{F} = m\mathbf{a}$) is seen to have the form of a friction term proportional to the power put into the particle. Derivation: the three steps below.

$$\dot{\gamma}_u = \frac{d}{dt} \frac{m_0 \gamma_u c^2}{m_0 c^2} = \frac{d}{dt} \frac{E}{m_0 c^2} = \frac{\dot{E}}{m_0 c^2} = \frac{\mathbf{u} \cdot \mathbf{F}}{m_0 c^2} = \frac{P}{m_0 c^2}.$$

Combining this with the other expression for $\dot{\gamma}_u$ above gives

$$\mathbf{u} \cdot \mathbf{a} = \frac{\mathbf{u} \cdot \mathbf{F}}{\gamma_u^3 m_0}.$$

The above equation for ${\bf F}$ now becomes

$$\frac{\mathbf{F}}{m_0\gamma_u} = \mathbf{a} + \frac{P}{m_0\gamma_u c^2}\mathbf{u} = \mathbf{a} + \gamma_u^2 \frac{\mathbf{u} \cdot \mathbf{a}}{c^2}\mathbf{u}.$$

Transformation of force

• Let's compute the transformation law for force by transforming to a frame S' moving with velocity v in the x-direction relative to S:

$$\mathbb{F}' = \gamma_{u'} \begin{pmatrix} F'_x \\ F'_y \\ F'_z \\ P'/c \end{pmatrix} = \mathbf{\Lambda} \mathbb{F} = \gamma_u \begin{pmatrix} \gamma & 0 & 0 & -\gamma\beta \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\gamma\beta & 0 & 0 & \gamma \end{pmatrix} \begin{pmatrix} F_x \\ F_z \\ P/c \end{pmatrix}$$

$$= \gamma_u \begin{pmatrix} \gamma[F_x - \beta P/c] \\ F_y \\ F_z \\ \gamma[P/c - \beta F_x] \end{pmatrix} = \frac{\gamma_{u'}}{\gamma(1 - \frac{u_x v}{c^2})} \begin{pmatrix} \gamma[F_x - \beta P/c] \\ F_y \\ F_z \\ \gamma[P/c - \beta F_x] \end{pmatrix}.$$

In the last step, we used the relation $\gamma_{u'} = \gamma_u \gamma [1 - u_x v/c^2]$ which we proved earlier when transforming the velocity 4-vector U — it followed from the fact that its normalization is Lorentz invariant, *i.e.*, $\mathbf{U}'^t \boldsymbol{\eta} \mathbf{U}' = \mathbf{U}^t \boldsymbol{\eta} \mathbf{U}$.

• The 4 components now give our desired force transformation equations:

$$\begin{array}{rcl} F_x' & = & \displaystyle \frac{F_x - \frac{v}{c^2}P}{1 - \frac{u_x v}{c^2}}, \\ F_y' & = & \displaystyle \frac{F_y}{\gamma \left(1 - \frac{u_x v}{c^2}\right)}, \\ F_z' & = & \displaystyle \frac{F_z}{\gamma \left(1 - \frac{u_x v}{c^2}\right)}, \\ P' & = & \displaystyle \frac{P - vF_x}{1 - \frac{u_x v}{c^2}}, \end{array}$$

where $P = \mathbf{u} \cdot \mathbf{F}$ as usual.

• If we take S to be the rest frame of the particle, then $\mathbf{u} = 0$, $P = \mathbf{u} \cdot \mathbf{F} = 0$ and this simplifies to $F'_x = F_x$, $F'_y = F_y/\gamma$, $F'_z = F_z/\gamma$, so in the frame S' where the particle is moving, the force is unaffected in the parallel direction and suppressed by γ in the transverse directions.

Transformation of acceleration

- We could derive expressions using an approach like for force, but the results are so messy that it's not particularly useful — it's better to deal with explicit problems as needed.
- Here's a useful special case that you get to derive on a problem set (probably PS7): For an arbitrary acceleration **a** in *S*, the acceleration **a'** in *S'* is related to **a** via

$$a_x = \frac{a'_x}{\gamma^3 (1 + v u'_x / c^2)^3} a_y = \frac{a'_y}{\gamma^2 (1 + v u'_x / c^2)^2},$$

with the **important caveat** that the expression for a_y is only valid for the case where either $u'_y = 0$ or $a'_x = 0$.