Solutions to Problems for Infinite Spatial Domain Prolems

and the Fourier Transform

18.303 Linear Partial Differential Equations

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1 Problem 1

(i) Show that

$$u(x,t) = u_0 \operatorname{erf}\left(\frac{x}{2\sqrt{t}}\right), \qquad t > 0, \qquad x \in \mathbb{R},$$
 (1)

where u_0 is a constant, is a solution of the heat equation

$$u_t = u_{xx}, (2)$$

and satisfies the initial condition

$$u(x,0) = f(x) = \begin{cases} u_0, & \text{if } x > 0, \\ 0, & \text{if } x = 0, \\ -u_0, & \text{if } x < 0 \end{cases}$$
 (3)

in the sense that

$$\lim_{t \to 0^{+}} u\left(x, t\right) = f\left(x\right).$$

Solution: We show by direct substitution that u(x,t) is the solution to the heat equation and initial condition. First, note that the error function is defined by

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-s^2} ds \tag{4}$$

Thus by the chain rule,

$$u_{t} = u_{0} \frac{2}{\sqrt{\pi}} \exp\left(-\frac{x^{2}}{4t}\right) \frac{x}{2t^{3/2}} \left(-\frac{1}{2}\right) = -\frac{u_{0}x}{2\sqrt{\pi}t^{3/2}} \exp\left(-\frac{x^{2}}{4t}\right)$$

$$u_{xx} = \frac{\partial^{2}}{\partial x^{2}} \left(u_{0} \operatorname{erf}\left(\frac{x}{2\sqrt{t}}\right)\right) = \frac{\partial}{\partial x} \left(u_{0} \frac{2}{\sqrt{\pi}} \frac{1}{2\sqrt{t}} \exp\left(-\frac{x^{2}}{4t}\right)\right)$$

$$= u_{0} \frac{2}{\sqrt{\pi}} \exp\left(-\frac{x^{2}}{4t}\right) \frac{1}{2\sqrt{t}} \left(-\frac{2x}{4t}\right)$$

$$= -\frac{u_{0}x}{2\sqrt{\pi}t^{3/2}} \exp\left(-\frac{x^{2}}{4t}\right)$$

$$= u_{t}$$

Thus Eq. (1) for u(x,t) satisfies the Heat Equation (2).

We now show that (1) satisfies the IC (3). Note that

$$x > 0: \lim_{t \to 0^{+}} u(x, t) = u_{0} \operatorname{erf} \left(\lim_{t \to 0^{+}} \frac{x}{2\sqrt{t}} \right) = u_{0} \operatorname{erf} (\infty) = u_{0}$$

$$x = 0: \lim_{t \to 0^{+}} u(0, t) = \lim_{t \to 0^{+}} 0 = 0$$

$$x < 0: \lim_{t \to 0^{+}} u(x, t) = u_{0} \operatorname{erf} \left(\lim_{t \to 0^{+}} \frac{x}{2\sqrt{t}} \right) = u_{0} \operatorname{erf} (-\infty) = -u_{0}$$

Thus, $\lim_{t\to 0^+} u(x,t) = f(x)$.

(ii) Give a physical interpretation of the solution. Sketch the curves $u\left(x,t\right)=const$ in the xt-plane.

Solution: u(x,t) is the temperature in an infinite rod with thermal diffusivity 1 and initial temperature u_0 for x > 0 and $-u_0$ for x < 0. You could also think of this as two infinite rods put together, with the temperature at x = 0 held at 0. The level curves are given by

$$u(x,t) = u_0 \operatorname{erf}\left(\frac{x}{2\sqrt{t}}\right) = const$$

and hence

$$\frac{x}{2\sqrt{t}} = const \implies t = cx^2$$

These are parabolas pointing upward, and are plotted in Figure 1.

(iii) Derive the solution (i) from the general solution we derived in class in terms of the heat kernel K(s, x, t), using the initial temperature u(x, 0) = f(x). NOTE: All that is required is a change of variable in the integral, and then writing the integral in terms of the error function erf. Also, f(x) does not decay as $x \to \infty$, but it turns out this requirement can be relaxed as long as the integrals exist.

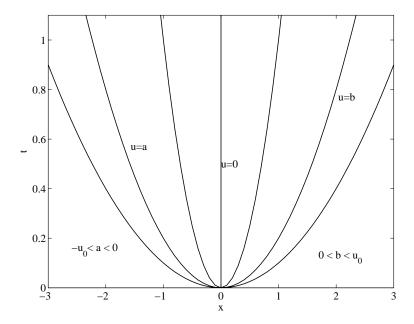


Figure 1: Level curves of u(x,t) for question 1(ii).

Solution: In class, we derived the solution of the heat equation with general initial condition f(x),

$$u(x,t) = \int_{-\infty}^{\infty} \frac{f(s)}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{(x-s)^2}{4\kappa t}\right) ds$$

Substituting f(s) from (3) gives

$$u\left(x,t\right) = -\frac{u_0}{\sqrt{4\pi\kappa t}} \int_{-\infty}^{0} \exp\left(-\frac{\left(x-s\right)^2}{4\kappa t}\right) ds + \frac{u_0}{\sqrt{4\pi\kappa t}} \int_{0}^{\infty} \exp\left(-\frac{\left(x-s\right)^2}{4\kappa t}\right) ds$$

Making the change of variable

$$\alpha = \frac{s - x}{\sqrt{4\kappa t}}$$

in the integrals gives

$$u(x,t) = -\frac{u_0}{\pi^{1/2}} \int_{-\infty}^{-x/\sqrt{4\kappa t}} e^{-\alpha^2} d\alpha + \frac{u_0}{\pi^{1/2}} \int_{-x/\sqrt{4\kappa t}}^{\infty} e^{-\alpha^2} ds$$

$$= -\frac{u_0}{\pi^{1/2}} \left(\int_{-\infty}^{0} e^{-\alpha^2} d\alpha + \int_{0}^{-x/\sqrt{4\kappa t}} e^{-\alpha^2} d\alpha \right)$$

$$+ \frac{u_0}{\pi^{1/2}} \left(\int_{-x/\sqrt{4\kappa t}}^{0} e^{-\alpha^2} ds + \int_{0}^{\infty} e^{-\alpha^2} ds \right)$$

$$= \frac{2u_0}{\pi^{1/2}} \int_{-x/\sqrt{4\kappa t}}^{0} e^{-\alpha^2} ds$$

$$= u_0 \frac{2}{\pi^{1/2}} \int_{0}^{x/\sqrt{4\kappa t}} e^{-\alpha^2} ds$$
(5)

since

$$\int_{-\infty}^{0} e^{-\alpha^2} d\alpha = \int_{0}^{\infty} e^{-\alpha^2} ds.$$

Substituting the definition of the error function (4) into (5) and setting $\kappa = 1$ gives

$$u(x,t) = u_0 \operatorname{erf}\left(\frac{x}{2\sqrt{t}}\right)$$

as required.

2 Problem 2

(i) Find the temperature u(x,t) of a semi-infinite rod $(x \ge 0)$, whose end (x = 0) is kept at a temperature of zero, and with an initial hot-spot, u(x,0) = f(x), where

$$f(x) = \begin{cases} u_0, & \text{if } x \in (x_0, x_1) \\ 0, & \text{if } x \in [0, x_0) \cup (x_1, \infty) \end{cases}$$
 (6)

with x_0 , x_1 constants, $0 \le x_0 < x_1$. Sketch the temperature profiles t = const (i.e., $u(x, t_0)$ in the ux-plane for various fixed times t_0), x = const (i.e., $u(x_0, t)$ in the ut-plane for various fixed x_0) and the level curves u(x, t) = const in the xt-plane. See note below.

- (ii) Repeat (i) with the end of the rod (x = 0) insulated. See note below.
- (iii) Referring to (ii), show that the temperature of the insulated end is a maximum at time

$$t = \frac{x_1^2 - x_0^2}{4\kappa \left(\log x_1 - \log x_0\right)}$$

where κ is the thermal diffusivity.

NOTE: in both (i) and (ii), just use the general solution we derived in class with the heat kernel, by suitably extending f(x) to the whole real line (i.e. odd extension or even extension - see class notes). The integrals in the solution can then be expressed as the sum of four terms involving error functions erf.

Solution (i): In class, we derived that solution to the heat equation on a semi-infinite rod $(x \ge 0)$ whose end is kept at zero is

$$u(x,t) = \int_{-\infty}^{\infty} \frac{\tilde{f}(s)}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{(x-s)^2}{4\kappa t}\right) ds$$
 (7)

where $\hat{f}(s)$ is the odd extension of f(s),

$$\tilde{f}(s) = \begin{cases} f(s), & s > 0 \\ 0, & s = 0 \\ -f(-s), & s < 0 \end{cases}$$
 (8)

Substituting (8) into (7) gives

$$u\left(x,t\right) = -\int_{-\infty}^{0} \frac{f\left(-s\right)}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{\left(x-s\right)^{2}}{4\kappa t}\right) ds + \int_{0}^{\infty} \frac{f\left(s\right)}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{\left(x-s\right)^{2}}{4\kappa t}\right) ds$$

Substituting for f(s) from (6) gives

$$u\left(x,t\right) = -\frac{u_0}{\sqrt{4\pi\kappa t}} \int_{-x_1}^{-x_0} \exp\left(-\frac{\left(x-s\right)^2}{4\kappa t}\right) ds + \frac{u_0}{\sqrt{4\pi\kappa t}} \int_{x_0}^{x_1} \exp\left(-\frac{\left(x-s\right)^2}{4\kappa t}\right) ds$$

Making the change of variable

$$\alpha = \frac{s - x}{\sqrt{4\kappa t}}$$

gives

$$\begin{array}{lll} u\left(x,t\right) & = & -\frac{u_{0}}{\sqrt{\pi}} \int_{\frac{-x_{1}-x}{\sqrt{4\kappa t}}}^{\frac{-x_{0}-x}{\sqrt{4\kappa t}}} e^{-\alpha^{2}} d\alpha + \frac{u_{0}}{\sqrt{\pi}} \int_{\frac{x_{0}-x}{\sqrt{4\kappa t}}}^{\frac{x_{1}-x}{\sqrt{4\kappa t}}} e^{-\alpha^{2}} d\alpha \\ & = & \frac{u_{0}}{\sqrt{\pi}} \left(-\int_{\frac{-x_{1}-x}{\sqrt{4\kappa t}}}^{0} e^{-\alpha^{2}} d\alpha - \int_{0}^{\frac{-x_{0}-x}{\sqrt{4\kappa t}}} e^{-\alpha^{2}} d\alpha + \int_{\frac{x_{0}-x}{\sqrt{4\kappa t}}}^{0} e^{-\alpha^{2}} d\alpha + \int_{0}^{\frac{x_{1}-x}{\sqrt{4\kappa t}}} e^{-\alpha^$$

Substituting the definition of the error function (4) gives

$$u\left(x,t\right) = \frac{u_0}{2} \left(-\operatorname{erf}\left(\frac{x_1 + x}{\sqrt{4\kappa t}}\right) + \operatorname{erf}\left(\frac{x_0 + x}{\sqrt{4\kappa t}}\right) - \operatorname{erf}\left(\frac{x_0 - x}{\sqrt{4\kappa t}}\right) + \operatorname{erf}\left(\frac{x_1 - x}{\sqrt{4\kappa t}}\right) \right)$$

The temperature profiles and the level curves are plotted in Figures 2 to 5.

Solution (ii): In class, we derived that solution to the heat equation on a semi-infinite rod $(x \ge 0)$ whose end is insulated,

$$u(x,t) = \int_{-\infty}^{\infty} \frac{\tilde{f}(s)}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{(x-s)^2}{4\kappa t}\right) ds$$
 (9)

where $\tilde{f}(s)$ is the odd extension of f(s),

$$\tilde{f}(s) = \begin{cases} f(s), & s > 0\\ f(-s), & s < 0 \end{cases}$$
(10)

Substituting (10) into (9) gives

$$u\left(x,t\right) = \int_{-\infty}^{0} \frac{f\left(-s\right)}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{\left(x-s\right)^{2}}{4\kappa t}\right) ds + \int_{0}^{\infty} \frac{f\left(s\right)}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{\left(x-s\right)^{2}}{4\kappa t}\right) ds$$

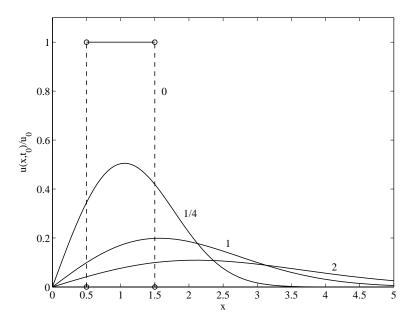


Figure 2: $u(x, t_0)$ curves for question 2(i). Numbers adjacent to curves indicate time t_0 .

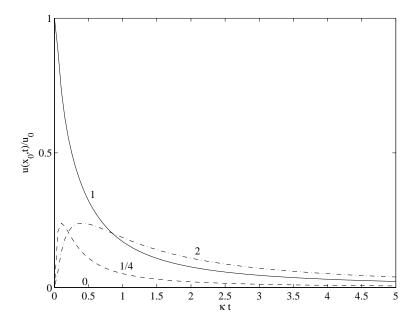


Figure 3: $u(x_0,t)$ curves for question 2(i). Numbers adjacent to curves indicate position x_0 .

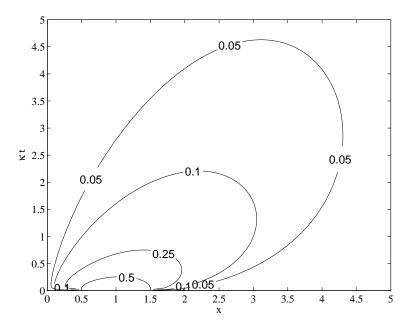


Figure 4: Level curves of u(x,t) for question 2(i). Numbers adjacent to curves indicate value of u on level curve.

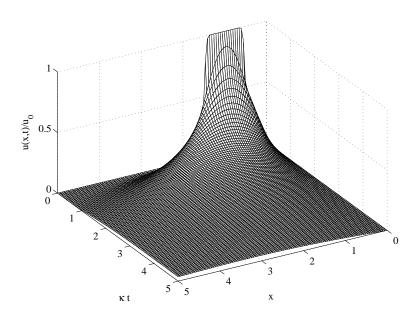


Figure 5: 3D plot of solution u(x,t) to question 2(i).

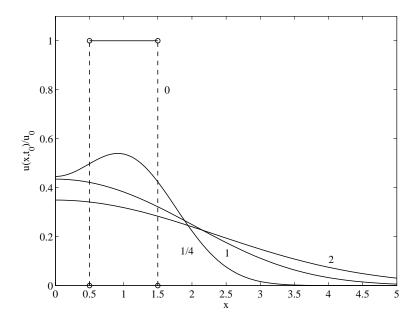


Figure 6: $u(x, t_0)$ curves for question 2(ii). Numbers adjacent to curves indicate time t_0 .

Substituting for f(s) from (6) gives

$$u\left(x,t\right) = \frac{u_0}{\sqrt{4\pi\kappa t}} \int_{-x_1}^{-x_0} \exp\left(-\frac{\left(x-s\right)^2}{4\kappa t}\right) ds + \frac{u_0}{\sqrt{4\pi\kappa t}} \int_{x_0}^{x_1} \exp\left(-\frac{\left(x-s\right)^2}{4\kappa t}\right) ds$$

Making the change of variable

$$\alpha = \frac{s - x}{\sqrt{4\kappa t}}$$

gives

$$u(x,t) = \frac{u_0}{\sqrt{\pi}} \int_{\frac{-x_1 - x}{\sqrt{4\kappa t}}}^{\frac{-x_0 - x}{\sqrt{4\kappa t}}} e^{-\alpha^2} d\alpha + \frac{u_0}{\sqrt{\pi}} \int_{\frac{x_0 - x}{\sqrt{4\kappa t}}}^{\frac{x_1 - x}{\sqrt{4\kappa t}}} e^{-\alpha^2} d\alpha$$

$$= \frac{u_0}{\sqrt{\pi}} \left(\int_{\frac{-x_1 - x}{\sqrt{4\kappa t}}}^{0} e^{-\alpha^2} d\alpha + \int_{0}^{\frac{-x_0 - x}{\sqrt{4\kappa t}}} e^{-\alpha^2} d\alpha + \int_{0}^{0} e^{-\alpha^2} d\alpha + \int_{0}^{\frac{x_1 - x}{\sqrt{4\kappa t}}} e^{-\alpha^2} d\alpha + \int_{0}^{\infty} e^{-\alpha^2} d\alpha + \int_{0}^{\frac{x_1 - x}{\sqrt{4\kappa t}}} e^{-\alpha^2} d\alpha + \int_{0}^{\infty} e^{-\alpha^2}$$

Substituting the definition of the error function (4) gives

$$u\left(x,t\right) = \frac{u_0}{2} \left(\operatorname{erf}\left(\frac{x_1 + x}{\sqrt{4\kappa t}}\right) - \operatorname{erf}\left(\frac{x_0 + x}{\sqrt{4\kappa t}}\right) - \operatorname{erf}\left(\frac{x_0 - x}{\sqrt{4\kappa t}}\right) + \operatorname{erf}\left(\frac{x_1 - x}{\sqrt{4\kappa t}}\right) \right)$$

The temperature profiles and the level curves are plotted in Figures 6 to 9.

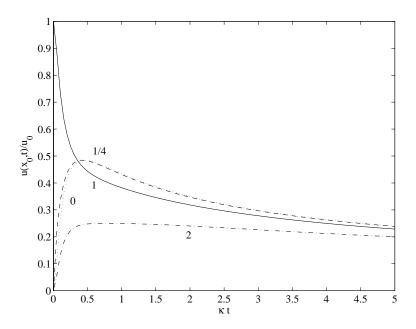


Figure 7: $u(x_0, t)$ curves for question 2(ii). Numbers adjacent to curves indicate position x_0 .

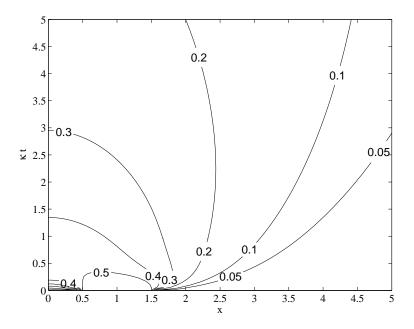


Figure 8: Level curves of u(x,t) for question 2(ii). Numbers adjacent to curves indicate value of u on level curve.

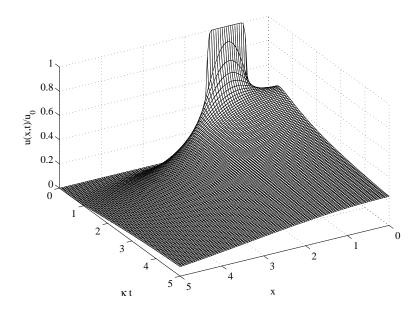


Figure 9: 3D plot of solution u(x, t) to question 2(ii).

Solution (iii): The temperature at the insulated end of the solution in problem (ii) is

$$u(0,t) = u_0 \left(\operatorname{erf} \left(\frac{x_1}{\sqrt{4\kappa t}} \right) - \operatorname{erf} \left(\frac{x_0}{\sqrt{4\kappa t}} \right) \right)$$

The temperature is maximum when $u_t(0,t) = 0$,

$$0 = u_t(0, t) = -\frac{u_0}{4\sqrt{\kappa}t^{3/2}} \left(x_1 \exp\left(-\frac{x_1^2}{4\kappa t}\right) - x_0 \exp\left(-\frac{x_0^2}{4\kappa t}\right) \right)$$

Rearranging gives

$$x_1 = x_0 \exp\left(\frac{x_1^2 - x_0^2}{4\kappa t}\right)$$

Taking the log of both sides and solving for t yields

$$t = \frac{x_1^2 - x_0^2}{4\kappa \log(x_1/x_0)} = \frac{x_1^2 - x_0^2}{4\kappa (\log(x_1) - \log(x_0))}$$

3 Problem 3

Show that

$$u(x,y) = \frac{2u_0}{\pi} \arctan\left(\frac{x}{y}\right) \tag{11}$$

where u_0 is constant, is a solution of Laplace's equation

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0,$$

and satisfies the boundary condition

$$\lim_{y \to 0^+} u\left(x, y\right) = f\left(x\right)$$

Give a physical interpretation of the solution (i.e. how does this relate to what Heat Problem?). Sketch the isothermal curves (level curves) u(x,y) = const in the xy-plane. Note that in polar coordinates,

$$\theta = \arctan\left(\frac{x}{y}\right)$$

where θ is the angle measured from the y-axis ($\theta = 0$ is the y-axis) and increasing clockwise.

Solution: We can show u(x,y) is a solution of Laplace's equation directly,

$$\frac{\partial u}{\partial x} = \frac{2u_0}{\pi} \frac{1}{1 + (x/y)^2} \frac{1}{y}, \qquad \frac{\partial^2 u}{\partial x^2} = -\frac{2u_0}{\pi} \frac{2xy}{(y^2 + x^2)^2}
\frac{\partial u}{\partial y} = \frac{2u_0}{\pi} \frac{1}{1 + (x/y)^2} \frac{-x}{y^2}, \qquad \frac{\partial^2 u}{\partial y^2} = \frac{2u_0}{\pi} \frac{2xy}{(y^2 + x^2)^2}$$

Thus

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = -\frac{2u_0}{\pi} \frac{2xy}{(y^2 + x^2)^2} + \frac{2u_0}{\pi} \frac{2xy}{(y^2 + x^2)^2} = 0$$

and hence (11) is a solution of Laplace's equation. The limit $y \to 0^+$ of (11) is

$$x > 0: \lim_{y \to 0^{+}} u(x, y) = \frac{2u_{0}}{\pi} \lim_{y \to 0^{+}} \arctan\left(\frac{x}{y}\right) = \frac{2u_{0}}{\pi} \arctan\left(\infty\right) = \frac{2u_{0}}{\pi} \frac{\pi}{2} = u_{0}$$

$$x = 0: \lim_{y \to 0^{+}} u(0, y) = \frac{2u_{0}}{\pi} \lim_{y \to 0^{+}} 0 = 0$$

$$x < 0: \lim_{y \to 0^{+}} u(x, y) = \frac{2u_{0}}{\pi} \lim_{y \to 0^{+}} \arctan\left(\frac{x}{y}\right) = \frac{2u_{0}}{\pi} \arctan\left(-\infty\right) = -\frac{2u_{0}}{\pi} \frac{\pi}{2} = -u_{0}$$

Thus

$$\lim_{y \to 0^+} u\left(x, y\right) = f\left(x\right)$$

where f(x) is given in (3) in problem 1.

We can also show that (11) is a solution of Laplace's equation by substituting f(s) into the general solution we derived in class using the Fourier Transform,

$$u(x,y) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(s) \frac{2y}{(x-s)^2 + y^2} ds$$

$$= \frac{1}{2\pi} \left(\int_{-\infty}^{0} f(s) \frac{2y}{(x-s)^2 + y^2} ds + \int_{0}^{\infty} f(s) \frac{2y}{(x-s)^2 + y^2} ds \right)$$

$$= \frac{yu_0}{\pi} \left(-\int_{-\infty}^{0} \frac{ds}{(x-s)^2 + y^2} + \int_{0}^{\infty} \frac{ds}{(x-s)^2 + y^2} \right)$$

Note that

$$\int_{a}^{b} \frac{ds}{(s-x)^{2} + y^{2}} = \int_{a-x}^{b-x} \frac{d\alpha}{\alpha^{2} + y^{2}} = \left[\frac{1}{y}\arctan\left(\frac{\alpha}{y}\right)\right]_{\alpha=a-x}^{b-x}$$
$$= \frac{1}{y}\arctan\left(\frac{b-x}{y}\right) - \frac{1}{y}\arctan\left(\frac{a-x}{y}\right)$$

Thus

$$u(x,y) = \frac{yu_0}{\pi} \left(-\frac{1}{y} \arctan\left(\frac{-x}{y}\right) + \frac{1}{y} \arctan\left(\frac{-\infty - x}{y}\right) \right)$$

$$+ \frac{yu_0}{\pi} \left(\frac{1}{y} \arctan\left(\frac{\infty - x}{y}\right) - \frac{1}{y} \arctan\left(\frac{-x}{y}\right) \right)$$

$$= \frac{u_0}{\pi} \left(2 \arctan\left(\frac{x}{y}\right) - \frac{\pi}{2} + \frac{\pi}{2} \right)$$

$$= \frac{2u_0}{\pi} \arctan\left(\frac{x}{y}\right)$$

as required.

Physical interpretation: the solution u(x,y) is the steady-state temperature of the upper half plane with boundary condition u(x,0) = f(x). Since $\theta = \arctan(x/y)$ where θ is the angle measured from the y-axis ($\theta = 0$ is the y-axis) and increasing clockwise, we have

$$u\left(x,y\right) = u_0 \frac{\theta}{\pi/2}$$

The level curves u = const are thus lines through the origin, $\theta = const$, and are sketched in Figure 10.

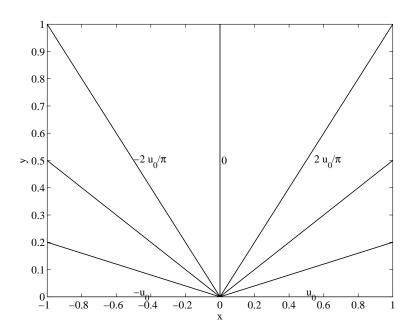


Figure 10: Level curves of u(x, y) for question 3.