## 19 Stream functions and conformal maps

There is a useful device for thinking about two dimensional flows, called the stream function of the flow. The stream function $\psi(x, y)$ is defined as follows

$$
\begin{equation*}
\boldsymbol{u}=(u, v)=\left(\frac{\partial \psi}{\partial y},-\frac{\partial \psi}{\partial x}\right) . \tag{469}
\end{equation*}
$$

The velocity field described by $\psi$ automatically satisfies the incompressibility condition, and it should be noted that

$$
\begin{equation*}
\boldsymbol{u} \cdot \nabla \psi=u \frac{\partial \psi}{\partial x}+v \frac{\partial \psi}{\partial y}=0 . \tag{470}
\end{equation*}
$$

Thus $\psi$ is constant along streamlines of the flow. Besides it's physical convenience, another great thing about the stream function is the following. By definition

$$
\begin{align*}
u & =\frac{\partial \psi}{\partial y}=\frac{\partial \phi}{\partial x},  \tag{471a}\\
v & =-\frac{\partial \psi}{\partial x}=\frac{\partial \phi}{\partial y}, \tag{471b}
\end{align*}
$$

where $\phi$ is the velocity potential for an irrotational flow. Thus, both $\phi$ and $\psi$ obey the well known Cauchy-Riemann equations of complex analysis.

### 19.1 The Cauchy-Riemann equations

In complex analysis you work with the complex variable $z=x+i y$. Thus, if you have some complex function $f(z)$ what is $d f / d z$ ? Well, $f(z)$ can be separated into a real part $u(x, y)$ and an imaginary part $v(x, y)$, where $u$ and $v$ are real functions, i.e.:

$$
\begin{equation*}
f(z)=f(x+i y)=u(x, y)+i v(x, y) . \tag{472}
\end{equation*}
$$

For example, if $f(z)=z^{2}$ then $u=x^{2}-y^{2}$ and $v=2 x y$. What then is $d f / d z$ ? Since we are now in two-dimensions we can approach a particular point $z$ from the $x$-direction or the $y$-direction (or any other direction, for that matter). On one hand we could define

$$
\begin{equation*}
\frac{d f}{d z}=\frac{\partial f}{\partial x}=\frac{\partial u}{\partial x}+i \frac{\partial v}{\partial x} \tag{473a}
\end{equation*}
$$

Or, alternatively

$$
\begin{equation*}
\frac{d f}{d z}=\frac{\partial f}{\partial(i y)}=-i \frac{\partial f}{\partial y}=-i \frac{\partial u}{\partial y}+\frac{\partial v}{\partial y} . \tag{473b}
\end{equation*}
$$

For the definition of the derivative to make sense requires $\partial u / \partial x=\partial v / \partial y$ and $-\partial u / \partial y=$ $\partial v / \partial x$, the Cauchy-Riemann equations. If this is true then $f(z)$ is said to be analytic and we can simply differentiate with respect to $z$ in the usual manner. For our simple example $f(z)=z^{2}$ we have that $d f / d z=2 z$ (confirm for yourself that $z^{2}$ is analytic as there are many functions that are not, e.g., $|z|$ is not an analytic function.)

### 19.2 Conformal mapping

We can now use the power of complex analysis to think about two dimensional potential flow problems. Since $\phi$ and $\psi$ obey the Cauchy-Riemann equations, this implies that $w=\phi+i \psi$ is an analytic function of the complex variable $z=x+i y$. We call $w$ the complex potential. Another important property of 2D incompressible flow is that both $\phi$ and $\psi$ satisfy Laplace's equation. For example, using the Cauchy-Riemann equations we see that

$$
\begin{equation*}
\frac{\partial \psi}{\partial x^{2}}+\frac{\partial \psi}{\partial y^{2}}=-\frac{\partial^{2} \psi}{\partial x \partial y}+\frac{\partial^{2} \psi}{\partial y \partial x}=0 \tag{474}
\end{equation*}
$$

The same proof can be used for $\phi$. We can therefore consider any analytic function (e.g., $\left.\sin z, z^{4}, \ldots\right)$, calculate the real and imaginary parts and both of them satisfy Laplace's equation.

The velocity components $u$ and $v$ are directly related to $d w / d z$, which is conveniently calculated as follows:

$$
\begin{equation*}
\frac{d w}{d z}=\frac{\partial \phi}{\partial x}+i \frac{\partial \psi}{\partial x}=u-i v . \tag{475}
\end{equation*}
$$

As a simple example consider uniform flow at an angle $\alpha$ to the $x$-axis. The corresponding complex potential is $w=u_{0} z e^{-i \alpha}$. In this case $d w / d z=u_{0} e^{-i \alpha}$. Using the above relation, this tells us that $u=u_{0} \cos \alpha$ and $v=u_{0} \sin \alpha$.

We can also determine the complex potential for flow past a cylinder since we know that

$$
\begin{equation*}
\phi=u_{0}\left(r+\frac{R^{2}}{r}\right) \cos \theta \tag{476}
\end{equation*}
$$

and this is just the real part of the complex potential

$$
\begin{equation*}
w=u_{0}\left(z+\frac{R^{2}}{z}\right) . \tag{477}
\end{equation*}
$$

Check this by substituting in $z=r e^{i \theta}$. What is the corresponding stream function? Also $w(z)=-i \ln z$ is the complex potential for a point vortex since

$$
\begin{equation*}
\operatorname{Re}(w(z))=\operatorname{Re}\left(-i \ln \left(r e^{i \theta}\right)\right)=\theta \tag{478}
\end{equation*}
$$

and we know that $\phi=\theta$ is the real potential for a point vortex. Thus

$$
\begin{equation*}
w(z)=u_{0}\left(z+\frac{R^{2}}{z}\right)-\frac{i \Gamma}{2 \pi} \ln z \tag{479}
\end{equation*}
$$

is the complex potential for flow past a cylinder with circulation $\Gamma$.
So let's assume that the only problem we know how to solve is flow past a cylinder, when really we want to know how to solve for flow past an aerofoil. The idea is to now consider two complex planes $(x, y)$ and $(X, Y)$. In the first plane we have the complex variable $z=x+i y$ and in the latter we have $Z=X+i Y$. If we construct a mapping $Z=F(z)$ which is analytic, with an inverse $z=F^{-1}(Z)$, then $W(Z)=w\left(F^{-1}(Z)\right)$ is also analytic, and may be considered a complex potential in the new co-ordinate system. Because $W(Z)$ and $w(z)$ take the same value at corresponding points of the two planes it follows that $\Psi$ and $\psi$ are the same at corresponding points. Thus streamlines are mapped into streamlines. In particular a solid boundary in the $z$-plane, which is necessarily a streamline, gets mapped into a streamline in the $Z$-plane, which could accordingly be viewed as a rigid boundary. Thus all we have done is distort the streamlines and the boundary leaving us with the key question: Given flow past a circular cylinder in the $z$-plane can we choose a mapping so as to obtain in the $Z$-plane uniform flow past a more wing-like shape? (Note that we have brushed passed some technical details here, such as the requirement that $d F / d z \neq 0$ at any point, as this would cause a blow-up of the velocity).

### 19.3 Simple conformal maps

The simplest map is

$$
\begin{equation*}
Z=F(z)=z+b, \tag{480}
\end{equation*}
$$

which corresponds to a translation. Then there is

$$
\begin{equation*}
Z=F(z)=z e^{i \alpha}, \tag{481}
\end{equation*}
$$

which corresponds to a rotation through angle $\alpha$. In this case, the complex potential for uniform flow past a cylinder making angle $\alpha$ with the stream is

$$
\begin{equation*}
W(Z)=u_{0}\left(Z e^{-i \alpha}+\frac{R^{2}}{Z} e^{i \alpha}\right)-\frac{i \Gamma}{2 \pi} \ln Z . \tag{482}
\end{equation*}
$$

Note, that this expression could also include the term $\ln e^{i \alpha}=i \alpha$ which I have neglected. This is just a constant however and doesn't change the velocity.

Finally there is the non-trivial Joukowski transformation,

$$
\begin{equation*}
Z=F(z)=z+\frac{c^{2}}{z} . \tag{483}
\end{equation*}
$$

What does this do to the circle? Well, $z=a e^{i \theta}$ becomes

$$
\begin{equation*}
Z=a e^{i \theta}+\frac{c^{2}}{a} e^{-i \theta}=\left(a+\frac{c^{2}}{a}\right) \cos \theta+i\left(a-\frac{c^{2}}{a}\right) \sin \theta . \tag{484}
\end{equation*}
$$

Defining $X=\operatorname{Re}(Z), Y=\operatorname{Im}(Z)$, it is easily shown that

$$
\begin{equation*}
\left(\frac{X}{a+\frac{c^{2}}{a}}\right)^{2}+\left(\frac{Y}{a-\frac{c^{2}}{a}}\right)^{2}=1 \tag{485}
\end{equation*}
$$

which is the equation of an ellipse, provided $c<a$.

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