

## Problem Set No. 8

### Problem 1

123 <sup>system</sup> coordinate is fixed to the inner gimbal.

$$\# \text{ DOF} = 6 - 3 = 3$$

generalized coordinates:

$$q_1 = \theta, \quad q_2 = \varphi, \quad q_3 = \psi$$

If we need torque  $M_\psi$  about spin axis and torque  $M_\varphi$  about the vertical axis to maintain the motion, the generalized forces can be found as follows:

$$\begin{aligned} \delta W &= M_\psi \delta\psi + M_\varphi (\delta\psi \cos\theta) + M_\varphi \delta\varphi + M_\psi (\delta\varphi \cos\theta) \\ &= (M_\psi + M_\varphi \cos\theta) \delta\psi + (M_\varphi + M_\psi \cos\theta) \delta\varphi \end{aligned}$$

$$\Rightarrow Q_\theta = 0$$

$$Q_\varphi = M_\varphi + M_\psi \cos\theta$$

$$Q_\psi = M_\psi + M_\varphi \cos\theta$$

$$\mathcal{L} = T - V$$

$$V = 0 \quad (\text{center of mass does not move})$$

$$T = \frac{1}{2} \underline{\omega}^T \underline{I}_c \underline{\omega}$$

$$\underline{\omega} = \dot{\varphi} \underline{e}_z + \dot{\theta} \underline{e}_2 + \dot{\psi} \underline{e}_1 = (\dot{\psi} + \dot{\varphi} \cos\theta) \underline{e}_1 + \dot{\theta} \underline{e}_2 + \dot{\varphi} \sin\theta \underline{e}_3$$

$$T = \frac{1}{2} I_1 (\dot{\psi} + \dot{\varphi} \cos\theta)^2 + \frac{1}{2} I_2 \dot{\theta}^2 + \frac{1}{2} I_2 \dot{\varphi}^2 \sin^2\theta = L \quad (V=0)$$

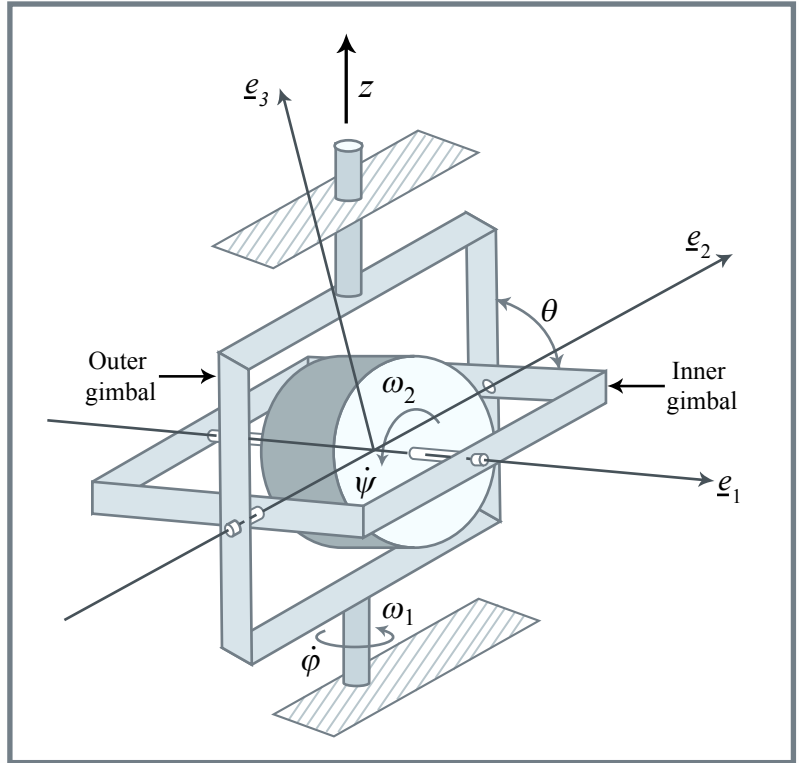
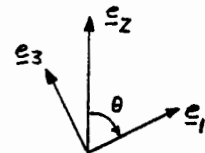
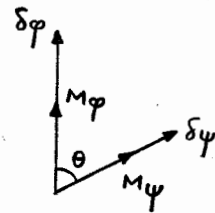


Figure by OCV.



$$\underline{e}_2 = \cos\theta \underline{e}_1 + \sin\theta \underline{e}_3$$

## Problem 1

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = Q_{\theta} \quad \rightarrow \quad I_2 \ddot{\theta} + I_1 (\dot{\psi} + \dot{\phi} \cos \theta) \dot{\phi} \sin \theta - I_2 \dot{\phi}^2 \sin \theta \cos \theta = 0$$

Note that  $\dot{\phi} = \omega_1$  and  $\dot{\psi} = \omega_2$  are known so the system has one degree of freedom  $\theta$  and the equation of motion would be:

$$\underline{I_2 \ddot{\theta} + (I_1 - I_2) \omega_1^2 \sin \theta \cos \theta + I_1 \omega_1 \omega_2 \sin \theta = 0}$$

To find the generalized forces  $Q_{\phi}$  and  $Q_{\psi}$ :

$$Q_{\phi} = \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\phi}} \right) - \frac{\partial L}{\partial \phi} = \frac{d}{dt} \left[ I_1 (\dot{\psi} + \dot{\phi} \cos \theta) \cos \theta + I_2 \dot{\phi} \sin^2 \theta \right]$$

$$Q_{\phi} = I_1 \dot{\omega}_1 \cos^2 \theta - I_1 \dot{\theta} \omega_2 \sin \theta + I_2 \dot{\omega}_1 \sin^2 \theta + (I_2 - I_1) \dot{\theta} \omega_1 \sin 2\theta$$

$$Q_{\psi} = \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\psi}} \right) - \frac{\partial L}{\partial \psi} = I_1 (\ddot{\psi} + \ddot{\phi} \cos \theta) - I_1 \dot{\phi} \dot{\theta} \sin \theta =$$

$$Q_{\psi} = I_1 \dot{\omega}_1 \cos \theta - I_1 \dot{\omega}_1 \dot{\theta} \sin \theta$$

$M_{\phi}$  and  $M_{\psi}$  which, respectively, correspond to  $M_z$  and  $M_x$  (found in the Quiz problem)

can be determined using the following equations:

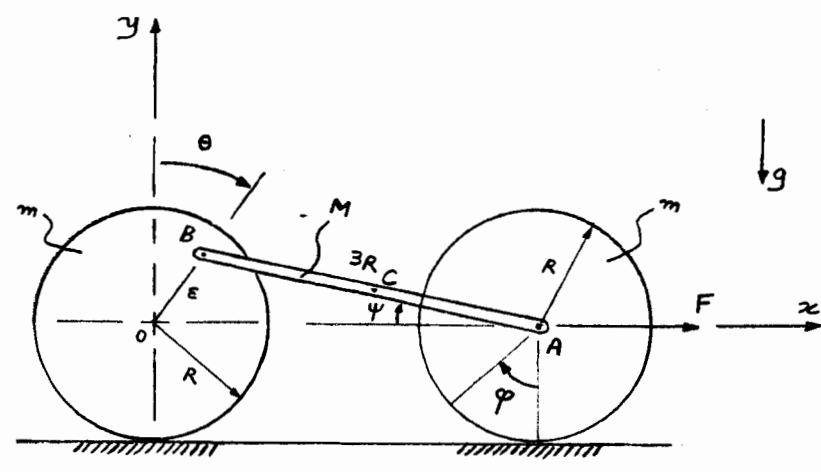
$$\begin{cases} M_{\phi} = \frac{Q_{\phi} - \cos \theta Q_{\psi}}{\sin^2 \theta} \\ M_{\psi} = \frac{Q_{\psi} - \cos \theta Q_{\phi}}{\sin^2 \theta} \end{cases}$$

# Problem 2

$$\omega|_{\text{left cyl.}} = -\dot{\theta} e_z$$

$$\omega|_{\text{right cyl.}} = -\dot{\phi} e_z$$

$$\omega|_{\text{rod}} = -\dot{\psi} e_z$$

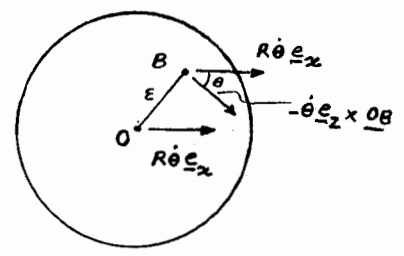


$$v_A = R\dot{\phi} e_x$$

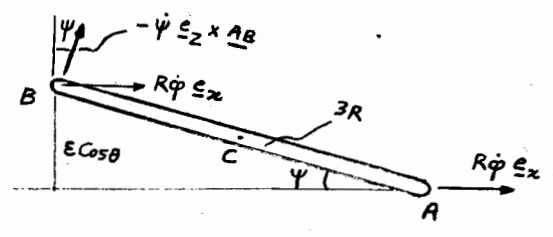
$$v_B = v_O + (-\dot{\theta} e_z) \times \underline{OB}$$

$$= R\dot{\theta} e_x + \dot{\theta} E (\cos\theta e_x - \sin\theta e_y)$$

$$= (R + E\cos\theta)\dot{\theta} e_x - E\sin\theta\dot{\theta} e_y$$



(point B on the left cylinder)



$$v_B|_{\text{rod}} = v_A + \omega|_{\text{rod}} \times \underline{AB}$$

$$= R\dot{\phi} e_x + 3R\dot{\psi} (\sin\psi e_x + \cos\psi e_y)$$

$$= (R\dot{\phi} + 3R\dot{\psi} \sin\psi) e_x + 3R\dot{\psi} \cos\psi e_y$$

$$v_B|_{\text{left cyl.}} = v_B|_{\text{rod}} \Rightarrow$$

$$\begin{cases} (R + E\cos\theta)\dot{\theta} = R\dot{\phi} + 3R\dot{\psi} \sin\psi \\ -E\sin\theta\dot{\theta} = 3R\dot{\psi} \cos\psi \end{cases}$$

$$\Rightarrow \begin{cases} \dot{\psi} = \frac{-E\sin\theta\dot{\theta}}{\sqrt{9R^2 - E^2\cos^2\theta}} \\ \dot{\phi} = \frac{(R + E\cos\theta)\dot{\theta}}{R} + \frac{E^2\sin\theta\cos\theta\dot{\theta}}{R\sqrt{9R^2 - E^2\cos^2\theta}} \end{cases}$$

(\*)

## Problem 2

$$\underline{v}_C = \underline{v}_A + \underline{\omega}|_{rod} \times \underline{AC} = R\dot{\varphi} \underline{e}_x + \dot{\psi} \frac{3R}{2} (\sin\psi \underline{e}_x + \cos\psi \underline{e}_y)$$

$$= \left( R\dot{\varphi} + \frac{3R}{2} \dot{\psi} \sin\psi \right) \underline{e}_x + \frac{3R}{2} \dot{\psi} \cos\psi \underline{e}_y$$

$$\underline{v}_C = \left[ (R + \epsilon \cos\theta) \dot{\theta} + \frac{\epsilon^2 \sin\theta \cos\theta \dot{\theta}}{2\sqrt{9R^2 - \epsilon^2 \cos^2\theta}} \right] \underline{e}_x - \frac{\epsilon \sin\theta \dot{\theta}}{2} \underline{e}_y = G(\theta) \dot{\theta} \underline{e}_x - \frac{\epsilon \sin\theta \dot{\theta}}{2} \underline{e}_y$$

$$\# \text{ DOF} = 3 \times 3 - 2 - 2 - 2 - 2 = 1$$

$$q_1 = \theta$$

F is the only active non-potential force.

$$(*) \rightarrow \delta\varphi = \left( \frac{R + \epsilon \cos\theta}{R} + \frac{\epsilon^2 \sin\theta \cos\theta}{R\sqrt{9R^2 - \epsilon^2 \cos^2\theta}} \right) \delta\theta$$

$$\text{To find } Q_\theta, \quad \delta W_\theta = F(R\delta\varphi) = F \left( R + \epsilon \cos\theta + \frac{\epsilon^2 \sin\theta \cos\theta}{\sqrt{9R^2 - \epsilon^2 \cos^2\theta}} \right) \delta\theta$$

$$\Rightarrow Q_\theta = F \left( R + \epsilon \cos\theta + \frac{\epsilon^2 \sin\theta \cos\theta}{\sqrt{9R^2 - \epsilon^2 \cos^2\theta}} \right)$$

Construct Lagrangian:  $\mathcal{L} = T - V$

$$T = \frac{1}{2} m \underline{v}_O \cdot \underline{v}_O + \frac{1}{2} \left( \frac{1}{2} m R^2 \right) \dot{\theta}^2 + \frac{1}{2} m \underline{v}_A \cdot \underline{v}_A + \frac{1}{2} \left( \frac{1}{2} m R^2 \right) \dot{\varphi}^2 \\ + \frac{1}{2} M \underline{v}_C \cdot \underline{v}_C + \frac{1}{2} \left( \frac{1}{12} M (3R)^2 \right) \dot{\psi}^2$$

$$T = \frac{1}{2} m (R\dot{\theta})^2 + \frac{1}{4} m R^2 \dot{\theta}^2 + \frac{1}{2} m (R\dot{\varphi})^2 + \frac{1}{4} m R^2 \dot{\varphi}^2 + \frac{1}{2} M \left[ G(\theta)^2 + \frac{\epsilon^2 \sin^2\theta}{4} \right] \dot{\theta}^2 \\ + \frac{3}{8} M R^2 \frac{\epsilon^2 \sin^2\theta}{9R^2 - \epsilon^2 \cos^2\theta} \dot{\theta}^2$$

$$T = \frac{3}{4} m R^2 \dot{\theta}^2 + \frac{3}{4} m R^2 \dot{\varphi}^2 + \frac{1}{2} M \left[ G(\theta)^2 + \frac{\epsilon^2 \sin^2\theta}{4} \right] \dot{\theta}^2 + \frac{3}{8} M R^2 \frac{\epsilon^2 \sin^2\theta}{9R^2 - \epsilon^2 \cos^2\theta} \dot{\theta}^2$$

$$V = M g y_C = M g \left( \frac{3R}{2} \sin\psi \right) = M g \frac{\epsilon \cos\theta}{2} \quad (y_O = 0, y_A = 0)$$

## Problem 2

$$\begin{aligned} \mathcal{L} &= \frac{3}{4} m R^2 \dot{\theta}^2 + \frac{3}{4} m \left[ R + \epsilon \cos \theta + \frac{\epsilon^2 \sin \theta \cos \theta}{\sqrt{9R^2 - \epsilon^2 \cos^2 \theta}} \right]^2 \dot{\theta}^2 \\ &+ \frac{1}{2} M \left\{ \left[ R + \epsilon \cos \theta + \frac{\epsilon^2 \sin \theta \cos \theta}{2\sqrt{9R^2 - \epsilon^2 \cos^2 \theta}} \right]^2 + \frac{\epsilon^2 \sin^2 \theta}{4} \right\} \dot{\theta}^2 + \frac{3}{8} M R^2 \frac{\epsilon^2 \sin^2 \theta}{9R^2 - \epsilon^2 \cos^2 \theta} \dot{\theta}^2 \\ - Mg \frac{\epsilon}{2} \cos \theta &= \underline{H(\theta) \dot{\theta}^2 - Mg \frac{\epsilon}{2} \cos \theta} \end{aligned}$$

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\theta}} \right) - \frac{\partial \mathcal{L}}{\partial \theta} = Q_\theta$$

$$\frac{d}{dt} [2H(\theta) \dot{\theta}] - \left[ \frac{dH}{d\theta} \dot{\theta}^2 + Mg \frac{\epsilon}{2} \sin \theta \right] = Q_\theta$$

$$2H(\theta) \ddot{\theta} + \frac{dH}{d\theta} \dot{\theta}^2 - Mg \frac{\epsilon}{2} \sin \theta = Q_\theta$$

governing equation of motion

where

$$\begin{aligned} H(\theta) &= \frac{3}{4} m R^2 + \frac{3}{4} m \left[ R + \epsilon \cos \theta + \frac{\epsilon^2 \sin \theta \cos \theta}{\sqrt{9R^2 - \epsilon^2 \cos^2 \theta}} \right]^2 \\ &+ \frac{1}{2} M \left\{ \left[ R + \epsilon \cos \theta + \frac{\epsilon^2 \sin \theta \cos \theta}{2\sqrt{9R^2 - \epsilon^2 \cos^2 \theta}} \right]^2 + \frac{\epsilon^2 \sin^2 \theta}{4} \right\} + \frac{3}{8} M R^2 \frac{\epsilon^2 \sin^2 \theta}{9R^2 - \epsilon^2 \cos^2 \theta} \end{aligned}$$

and

$$Q_\theta = F \left( R + \epsilon \cos \theta + \epsilon^2 \frac{\sin \theta \cos \theta}{\sqrt{9R^2 - \epsilon^2 \cos^2 \theta}} \right)$$

Problem 3

# DOF = 2

generalized coordinates:

$$q_1 = s, \quad q_2 = \varphi$$

generalized forces:

$$Q_s = 0, \quad Q_\varphi = M$$

$$(\delta W = M \delta \varphi)$$

$$L = T - V$$

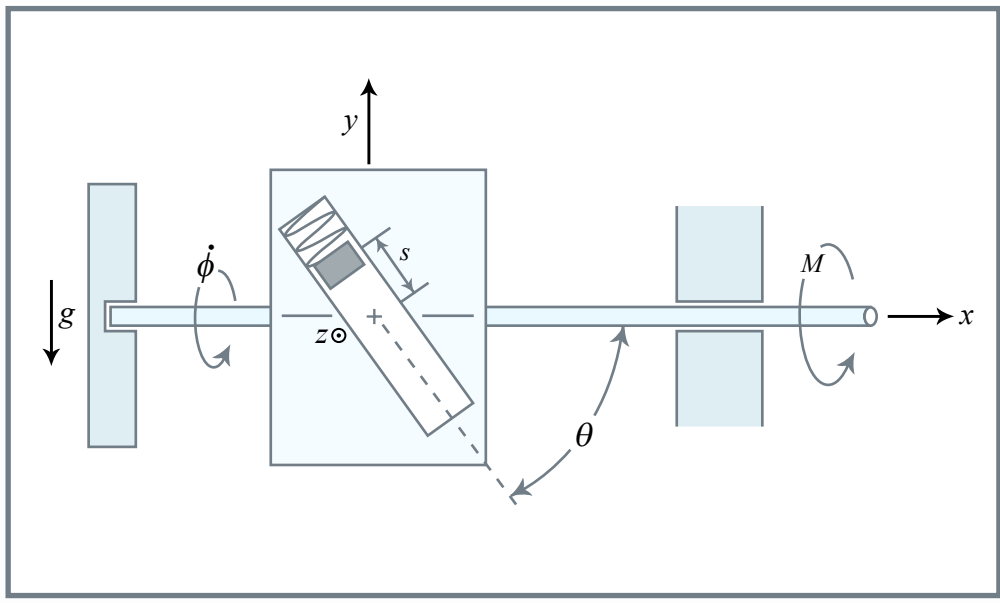


Figure by OCW. After problem 6.41 in Ginsberg, J. H. *Advanced Engineering Dynamics*. 2nd ed. New York: Cambridge University Press, 1998.

$$V_{slider} = mgs \sin \theta \cos \varphi$$

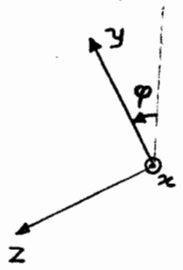
$$V_{spring} = \frac{1}{2} ks^2 \quad (k \text{ Spring stiffness})$$

$$V_{housing} = 0$$

$$T_{housing} = \frac{1}{2} I \dot{\varphi}^2$$

$$T_{slider} = \frac{1}{2} m v_s^2$$

xyz coordinate system rotates with  $\dot{\varphi}$  about x:



$$r_s = -s \cos \theta e_x + s \sin \theta e_y$$

$$v_s = \dot{r}_s + \omega \times r_s = \underbrace{-\dot{s} \cos \theta e_x + \dot{s} \sin \theta e_y}_{\dot{\varphi} e_x \times r_s} + \underbrace{\dot{\varphi} e_x \times r_s}_{\dot{\varphi} s \sin \theta e_z}$$

$$T_{slider} = \frac{1}{2} m (\dot{s}^2 + \dot{\varphi}^2 s^2 \sin^2 \theta)$$

$$\therefore L = \frac{1}{2} I \dot{\varphi}^2 + \frac{1}{2} m (\dot{s}^2 + \dot{\varphi}^2 s^2 \sin^2 \theta) - mgs \sin \theta \cos \varphi - \frac{1}{2} ks^2$$

$$\left. \begin{aligned} \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\varphi}} \right) - \frac{\partial L}{\partial \varphi} &= M \quad \rightarrow \quad (I + ms^2 \sin^2 \theta) \ddot{\varphi} + 2ms\dot{\varphi} \sin^2 \theta - mgs \sin \theta \sin \varphi = M \\ \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{s}} \right) - \frac{\partial L}{\partial s} &= 0 \quad \rightarrow \quad m\ddot{s} - ms\dot{\varphi}^2 \sin^2 \theta + mgs \sin \theta \cos \varphi + ks = 0 \end{aligned} \right\} \text{equations of motion}$$

# Problem 4

# DOF = 3

$q_1 = \theta, \quad q_2 = \psi, \quad q_3 = s$

(Spring is unstretched when  $s=0$  and  $\psi=0$ )

All active forces are potential.  $\Rightarrow Q_j = 0 \quad (j=1,2,3)$

$L = T - V$

$T = \frac{1}{2} m v_C^2 + \frac{1}{2} \omega^T I_C \omega$

$xyz$  rotates with  $\psi$  about vertical axis so that the bar is always in  $xy$  plane.  $XYZ$  is fixed to the bar.

Introducing additional generalized coordinates and using Lagrange multipliers, one can find the constraint forces as well.

The constraint force at point  $O$  has three components. The vertical component is simply  $-k_e s$  since the collar is massless.

To find the horizontal components of the constraint force:

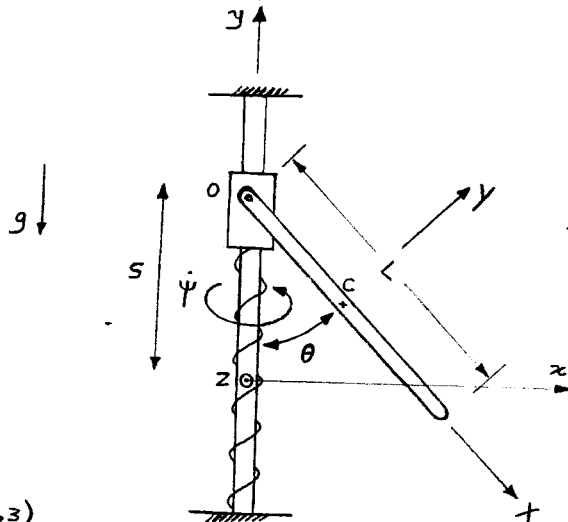
select  $q_4 = x, \quad q_5 = z$  ( $x$  and  $z$  are the coordinates of the upper end of the rod,

$\rightarrow n = 5$

Constraints:  $\begin{cases} x = 0 & \text{or} & q_4 = 0 \\ z = 0 & \text{or} & q_5 = 0 \end{cases} \rightarrow m = 2 \rightarrow \text{two Lagrange multipliers } \lambda_1, \lambda_2$

$q_4 = 0 \rightarrow a_{11} = a_{12} = a_{13} = a_{15} = 0, \quad a_{14} = 1$

$q_5 = 0 \rightarrow a_{21} = a_{22} = a_{23} = a_{24} = 0, \quad a_{25} = 1$



### Problem 4

$$\underline{r}_C = \left(\frac{L}{2} \sin \theta + z\right) \underline{e}_x + \left(s - \frac{L}{2} \cos \theta\right) \underline{e}_y + z \underline{e}_z$$

$$\underline{v}_C = \dot{\underline{r}}_C = \dot{\underline{r}}_C + \underbrace{\underline{\omega}_{xyz}}_{\dot{\psi} \underline{e}_y} \times \underline{r}_C = \left(\dot{x} + \dot{\psi} z + \frac{L}{2} \cos \theta \dot{\theta}\right) \underline{e}_x + \left(\dot{s} + \frac{L}{2} \sin \theta \dot{\theta}\right) \underline{e}_y + \left(\dot{z} - \dot{\psi} x - \dot{\psi} \frac{L}{2} \sin \theta\right) \underline{e}_z$$

$$\underline{\omega}_{rod} = \dot{\psi} \underline{e}_y + \dot{\theta} \underline{e}_z = \dot{\psi} (-\cos \theta \underline{e}_x + \sin \theta \underline{e}_y) + \dot{\theta} \underline{e}_z$$

$$\underline{I}_C = \frac{1}{12} mL^2 \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\begin{aligned} \rightarrow T &= \frac{1}{2} m \left[ \left(\dot{x} + \dot{\psi} z + \frac{L}{2} \cos \theta \dot{\theta}\right)^2 + \left(\dot{s} + \frac{L}{2} \sin \theta \dot{\theta}\right)^2 + \left(\dot{z} - \dot{\psi} x - \dot{\psi} \frac{L}{2} \sin \theta\right)^2 \right] \\ &+ \frac{1}{2} \left(\frac{1}{12} mL^2\right) (\dot{\theta}^2 + \dot{\psi}^2 \sin^2 \theta) \end{aligned}$$

$$V = mg y_C + \frac{1}{2} k_e s^2 + \frac{1}{2} k_t \psi^2 = mg \left(s - \frac{L}{2} \cos \theta\right) + \frac{1}{2} k_e s^2 + \frac{1}{2} k_t \psi^2$$

$$Q_j = 0 \quad (j = 1, \dots, 5)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = \lambda_1 a_{11} + \lambda_2 a_{21} = 0$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\psi}} - \frac{\partial L}{\partial \psi} = \lambda_1 a_{12} + \lambda_2 a_{22} = 0$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{s}} - \frac{\partial L}{\partial s} = \lambda_1 a_{13} + \lambda_2 a_{23} = 0$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = \lambda_1 a_{14} + \lambda_2 a_{24} = \lambda_1 \overset{K_4}{\quad}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{z}} - \frac{\partial L}{\partial z} = \lambda_1 a_{15} + \lambda_2 a_{25} = \lambda_2 \overset{K_5}{\quad}$$

$$\delta W_{\text{non-potential}} = \overset{K_4}{F_x} \delta x + \overset{K_5}{F_z} \delta z$$

$(F_x, F_z)$  horizontal components of the constraint force



## Problem 4

Applying the constraints  $x=0$  &  $z=0$ , the first three equations yield the governing equations:

$$\begin{cases} \frac{1}{3} mL\ddot{\theta} + \frac{1}{2} mL \sin\theta \ddot{s} - \frac{1}{3} mL^2 \sin\theta \cos\theta \dot{\psi}^2 + \frac{1}{2} mgL \sin\theta = 0 \\ \frac{1}{3} mL^2 \sin^2\theta \ddot{\psi} + \frac{2}{3} mL^2 \dot{\psi} \dot{\theta} \sin\theta \cos\theta + k_t \psi = 0 \\ m\ddot{s} + \frac{1}{2} mL \sin\theta \ddot{\theta} + \frac{1}{2} mL \dot{\theta}^2 \cos\theta + k_e s + mg = 0 \end{cases}$$

Constraint forces  $F_x$  and  $F_z$ :

$$F_x = K_4 = \lambda_1 = \left( \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} \right) \Bigg|_{\substack{x=\dot{x}=0 \\ z=\dot{z}=0}} = \underline{\underline{m \frac{L}{2} \ddot{\theta} \cos\theta - \frac{1}{2} mL \dot{\theta}^2 \sin\theta - \frac{1}{2} mL \dot{\psi}^2 \sin\theta}}$$

$$F_z = K_5 = \lambda_2 = \left( \frac{d}{dt} \frac{\partial L}{\partial \dot{z}} - \frac{\partial L}{\partial z} \right) \Bigg|_{\substack{x=\dot{x}=0 \\ z=\dot{z}=0}} = \underline{\underline{-m \frac{L}{2} \ddot{\psi} \sin\theta - mL \dot{\psi} \dot{\theta} \cos\theta}}$$

Note that the constraint torque about  $y$  axis is  $-k_t \psi$  and the constraint torque about  $x$  axis can be found by introducing a rotation about this axis.

Center of mass of the collar <sup>is stationary in</sup>  $x$  &  $z$  so the horizontal constraint force between the collar and the vertical rod is equal to  $(F_x, F_z)$ .