Fall 2004

Problem Set No. 3

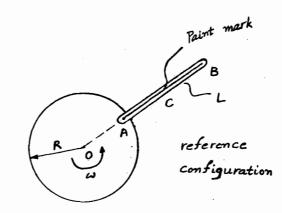
Problem 1

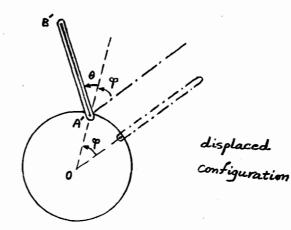
(a)

(6)

To find the angular velocity of the rod, compare orientation of AB to A'B':

 $\omega_{rod} = (\dot{\varphi} + \dot{\theta}) \underline{e}_z = (\omega + \dot{\theta}) \underline{e}_z$





Vc = VA + Wrod × AC

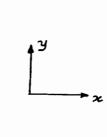
$$\underline{U}_A = \underline{\omega} \times \underline{OA} = \omega R \underline{e}_{\mu}$$

$$\therefore \quad \underline{v}_{c} = \omega R = \Psi + (\omega + \theta) \frac{L}{2} = \theta$$

$$e_{\psi} = -Sin\psi e_{x} + Cos\psi e_{y}$$

$$\underline{e}_{\theta} = -\cos\left[\theta - \left(\frac{\pi}{2} - \psi\right)\right] \underline{e}_{\chi}$$
$$-\sin\left[\theta - \left(\frac{\pi}{2} - \psi\right)\right] \underline{e}_{\chi}$$
$$\underline{e}_{\theta} = -\sin\left(\theta + \psi\right) \underline{e}_{\chi} + \cos\left(\theta + \psi\right) \underline{e}_{\chi}$$

If $\psi(t=0) = \psi_{0} \implies \psi(t) = \psi_{0} + \omega t$ So, $\psi_{c} = \left[-\omega_{R} \sin\psi - (\omega+\dot{\theta})\frac{L}{2} \sin(\theta+\psi)\right] e_{x} + \left[\omega_{R} \cos\psi + (\omega+\dot{\theta})\frac{L}{2} \cos(\theta+\psi)\right] e_{y}$



$$\mathcal{V}_{c} = \mathcal{V}_{e_{\mathcal{R}}} = \mathcal{V}_{cos\theta} = \mathcal{X} + \mathcal{V}_{sin\theta} = \mathcal{Y}_{cos\theta}$$

No slip between cylinder and

the bar:

So,

$$\mathcal{U}_{A}|_{bar} = \frac{\mathcal{U}_{A}}{\mathcal{Cylinder}}$$

$$\mathcal{V}_A|_{bar} = \mathcal{V}_{bar} \times \mathcal{B}_A \implies \mathcal{V}_A|_{bar} = -AB \theta e_{j}$$

$$\frac{\nabla A}{cylinder} = \frac{\nabla c}{c} + \frac{\omega}{cylinder} \times \frac{cA}{cA}$$
$$= \left(\nabla \cos \theta = x + \nabla \sin \theta = y \right) + R \frac{\omega}{cylinder} = x$$

$$\therefore \qquad AB \dot{\theta} \stackrel{e}{=}_{y} = \left(\mathcal{V} Cos \theta \stackrel{e}{=}_{\chi} + \mathcal{V} Sin \theta \stackrel{e}{=}_{y} \right) - R \omega_{cyl.} \stackrel{e}{=}_{\chi}$$

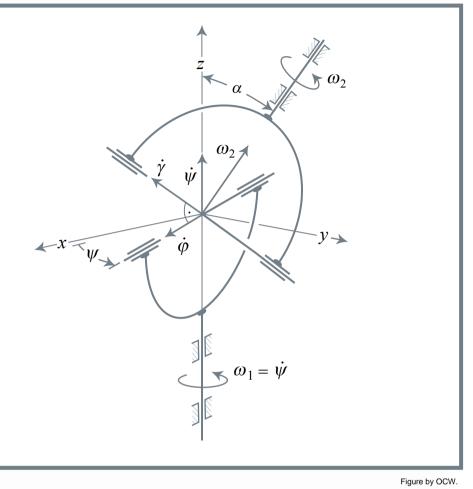
$$\Rightarrow (\nu Cos\theta - R \omega_{cyl.}) \stackrel{e}{=} \chi + (\nu sim\theta - AB \stackrel{o}{\theta}) \stackrel{e}{=} \chi = 0$$

$$\frac{CA}{AB} = \tan \frac{\theta}{2} \implies AB = R \cot \frac{\theta}{2}$$

(b)
$$\underline{\nu} \in f_{cylinder} = \underline{\nu}_{c} + \underline{\omega}_{cylinder} \times \underline{CE} = \left(\nu + \frac{\nu Cos\theta}{R} \right) \underline{e}_{z} = \nu \left(1 + \cos\theta \right) \underline{e}_{z}$$

velocity of the cylinder at the point where it contacts the ground. $\frac{\omega_2}{\omega_2} \text{ is alway in yz plane}$ and \dot{q} in zy plane: $\frac{\omega_2}{\omega_2} = \frac{\omega_2}{\sum_{\substack{\substack{0 \\ \text{Sin w} \\ \text{Cos w}}}}$ $\dot{q} = \dot{q} \left\{ \begin{array}{c} \cos \psi \\ \sin \psi \\ 0 \end{array} \right\}$ $\dot{\omega}_1 = \omega_1 \left\{ \begin{array}{c} 0 \\ 0 \\ 1 \end{array} \right\}$ \dot{z} is perpendicular to both $\dot{\omega}_2$ and \dot{q} so the direction

of is can be found:



After Problem 4-14 in Crandall, S. H., et al. Dynamics of Mechanical and Electromechanical Systems. Malabar, FL: Krieger, 1982.

$$\frac{e_{\dot{\varphi}}}{e_{\dot{\varphi}}} = \frac{e_{\dot{\varphi}} \times e_{\omega_2}}{\left| e_{\dot{\varphi}} \times e_{\omega_2} \right|} = \begin{bmatrix} e_{\chi} & e_{\chi} & e_{\chi} \\ \cos \psi & \sin \psi & 0 \\ 0 & \sin \psi & \cos \psi \end{bmatrix} = \begin{bmatrix} e_{\dot{\varphi}} \times e_{\omega_2} \\ e_{\dot{\varphi}} \times e_{\omega_2} \end{bmatrix}$$

$$\frac{e_{\dot{y}}}{e_{\dot{y}}} = \begin{cases} 5in\psi\cos\alpha \\ -\cos\psi\cos\alpha \\ -\cos\psi\cos\alpha \\ \cos\psi\sin\alpha \end{cases} \frac{1}{\sqrt{\cos^{2}\alpha + 5in^{2}\alpha\cos^{2}\psi}} \implies \dot{y} = \frac{\dot{y}}{\sqrt{\cos^{2}\alpha + 5in^{2}\alpha\cos\psi}} \begin{cases} 5in\psi\cos\alpha \\ -\cos\psi\cos\alpha \\ -\cos\psi\cos\alpha \\ \cos\psi\sin\alpha \end{cases} \\ \begin{cases} -\cos\psi\cos\alpha \\ \cos\psi\sin\alpha \\ \cos\psi\sin\alpha \end{cases} \end{cases}$$

The angular velocity of the cross can be determined in two ways:

 $\omega_{cross} = \omega_1 + \dot{\varphi}$

$$\frac{\omega}{c_{ross}} = \frac{\omega}{2} + \frac{\delta}{2}$$

Since the cross is a rigid body, it has a unique angular velocity:

$$\therefore \quad \underline{\omega}_1 + \dot{\underline{\phi}} = \underline{\omega}_2 + \dot{\underline{\delta}} \qquad (*)$$

Problem 3

$$\begin{array}{cccc} (*) \implies \begin{cases} 0\\0\\\omega_1 \\ \end{cases} + \begin{cases} \dot{\varphi} \cos \psi\\\dot{\varphi} \sin \psi\\0 \\ \end{cases} = \begin{cases} 0\\\omega_2 \sin \alpha\\\omega_2 \cos \alpha \\ \end{cases} + \frac{\dot{\delta}}{\sqrt{\frac{2}{\cos \alpha} + \frac{2}{\sin \alpha} \cos \psi}} \begin{cases} \sin \psi \cos \alpha\\-\cos \psi \cos \alpha\\\cos \psi \\ \end{cases} \\ \begin{array}{c} \cos \psi \sin \alpha\\\cos \psi \\ \end{array} \end{cases}$$

Let
$$\dot{\delta}_1 = \frac{\dot{\delta}}{\sqrt{\cos^2 \alpha + 5m^2 \cos^2 \psi}}$$

So we get three equations in three unknowns w, &, and p:

$$\begin{cases} \dot{\varphi} \cos \psi = \dot{\lambda}_{1} \sin \psi \cos \alpha \\ \dot{\varphi} \sin \psi = \omega_{2} \sin \alpha - \dot{\lambda}_{1} \cos \psi \cos \alpha \\ \omega_{1} = \omega_{2} \cos \alpha + \dot{\lambda}_{1} \cos \psi \sin \alpha \end{cases}$$

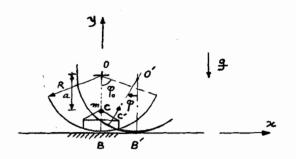
$$\Rightarrow \qquad \dot{s}_{1} = \omega_{1} \frac{Simar Cos \psi}{Cos^{2} \kappa + Simar Cos^{2} \psi} \qquad \dot{s} = \omega_{1} \frac{Simar Cos \psi}{\sqrt{Cos^{2} \kappa + Simar Cos^{2} \psi}}$$
$$\dot{\phi} = \omega_{1} \frac{Simar Cos \kappa Sim \psi}{Cos^{2} \kappa + Simar Cos^{2} \psi}$$

$$\omega_2 = \omega_1 \frac{\cos \alpha}{\cos \alpha + \sin^2 \alpha \cos^2 \psi}$$

equation of motion :

$$(R^2 + a^2 - 2aR \cos \varphi) \ddot{\varphi} + aR \sin \varphi \dot{\varphi}^2 + ag \sin \varphi = 0$$

- $\varphi \cdot \langle \varphi \langle \varphi \rangle$
 $(range of validity of this governing eq.)$



system is conservative. ___ T+V = Const.

$$r_{c} = (R\varphi - a \sin \varphi) \underline{e}_{x} + (R - a \cos \varphi) \underline{e}_{y}$$
$$\underline{v}_{c} = \dot{r}_{c} = (R - a \cos \varphi) \dot{\varphi} \underline{e}_{x} + a \sin \varphi \dot{\varphi} \underline{e}_{y}$$

$$T = \frac{1}{2}m\dot{\phi}^{2}\left\{\left(R-a\cos\varphi\right)^{2} + a^{2}\sin^{2}\varphi\right\}$$

$$T + V = \text{ const.} = E, \qquad \dot{\phi} = \frac{1}{2}\sqrt{2\frac{E_{*}-mg(R-a\cos\varphi)}{m(R^{2}+a^{2}-2aR\cos\varphi)}},$$

$$a \text{ real solution for } \dot{\phi} \text{ exists if } E_{*} \xrightarrow{} mg(R-a\cos\varphi).$$

For $E_{\bullet}^{(n)} \ge mg(R-a)$, if $E_{\bullet}^{(n)} = mg(R-a \operatorname{Gs} \varphi^{*(n)})$ then one can examine the behavior of

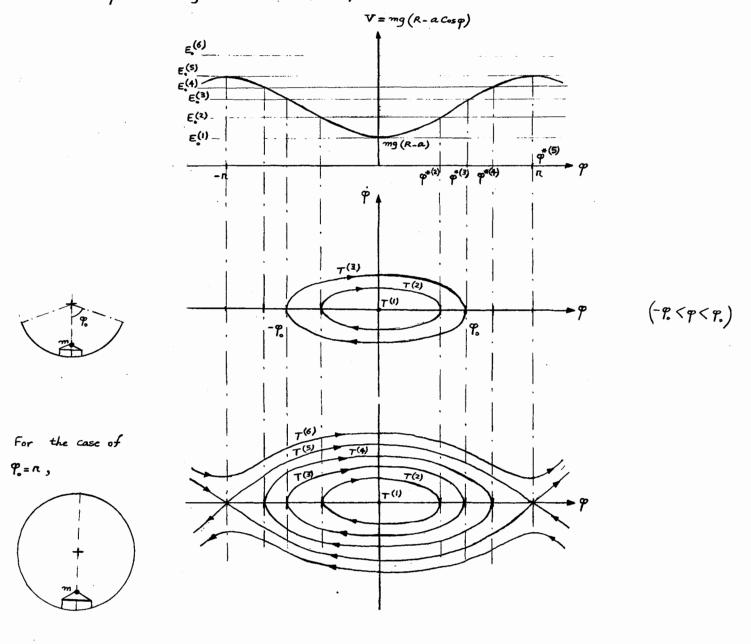
if it the neighborhood of $\varphi^{*(n)}$, $\varphi^{*(n)}_{+} \varphi$ (φ small):

$$\dot{\varphi} = \pm \sqrt{2 \frac{E_{\circ}^{(n)} - mg(R - a Cos(\varphi^{*(n)} + \varphi))}{m(R^{2} + a^{2} - 2aR Cos(\varphi^{*(n)} + \varphi))}} = \pm \sqrt{2 \frac{E_{\circ}^{(n)} - mg(R - a Cos \varphi^{*(n)}) - mga Sin \varphi^{*(n)} - mga Sin \varphi^{*($$

$$\therefore \quad \dot{\varphi} \sim \sqrt{|\varphi|} \qquad (\varphi \quad close \quad to \quad \varphi^{*(n)})$$

if $\sin \varphi^{*(n)} = 0$ $\left(\varphi^{*(n)} = \pm n\right)$, $\dot{\varphi} \sim \pm \varphi$

5



Qualitative plots of trajectories in the phase plane :

Note that $\varphi = 0$ is a stable equilibrium point and $\varphi = \pm \pi$ (in the case of $\varphi = \pi$) is an unstable equilibrium point.