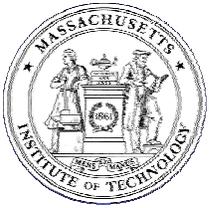


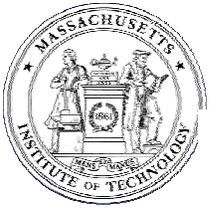
Computational Ocean Acoustics

- Ray Tracing
- Wavenumber Integration
- Normal Modes
- Parabolic Equation



Normal Modes

- Perturbation Approaches
 - Attenuation (5.8)
 - Group Velocity (5.8.1)
- Modes for Range-Dependent Envir.
 - Coupled Modes (5.9)
 - One-way Coupled Modes
 - Adiabatic Modes



Modal Group Velocity

$$u_n(\omega) = \frac{d\omega}{dk_{rn}}$$

$$u_n \simeq \frac{(\omega + \Delta\omega) - \omega}{k_{rn}(\omega + \Delta\omega) - k_{rn}(\omega)}$$

Perturbation Formulation

$$K^2(z) = \frac{(\omega + \Delta\omega)^2}{c^2(z)} \simeq \frac{\omega^2}{c^2(z)} + \frac{2 \Delta\omega \omega}{c^2(z)}$$

$$K^2 = K_0^2 + \epsilon K_1^2$$

$$K_0^2 = \omega^2 / c^2,$$

$$K_1^2 = 2\omega / c^2$$

$$\epsilon = \Delta\omega$$

$$k_{r1}^2 = \int_0^D \frac{2\omega}{c^2(z)} \frac{\Psi_0^2(z)}{\rho(z)} dz$$

Finite Difference Perturbation

$$k_r^2(\omega + \Delta\omega) \simeq k_{r0}^2(\omega) + \Delta\omega k_{r1}^2$$

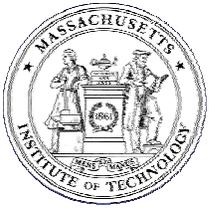
$$\frac{k_r^2(\omega + \Delta\omega) - k_{r0}^2(\omega)}{\Delta\omega} \simeq k_{r1}^2$$

$$\frac{k_r^2(\omega + \Delta\omega) - k_{r0}^2(\omega)}{\Delta\omega} \rightarrow_{\Delta\omega \rightarrow 0} \frac{dk_r^2}{d\omega}$$

$$\frac{d(k_r^2)}{d\omega} = 2k_r \frac{dk_r}{d\omega} = k_{r1}^2$$

Modal Group Slowness

$$\frac{dk_r}{d\omega} = \frac{k_{r1}^2}{2k_r} = \frac{\omega}{k_r} \int_0^D \frac{\Psi_0^2(z)}{\rho(z) c^2(z)} dz$$

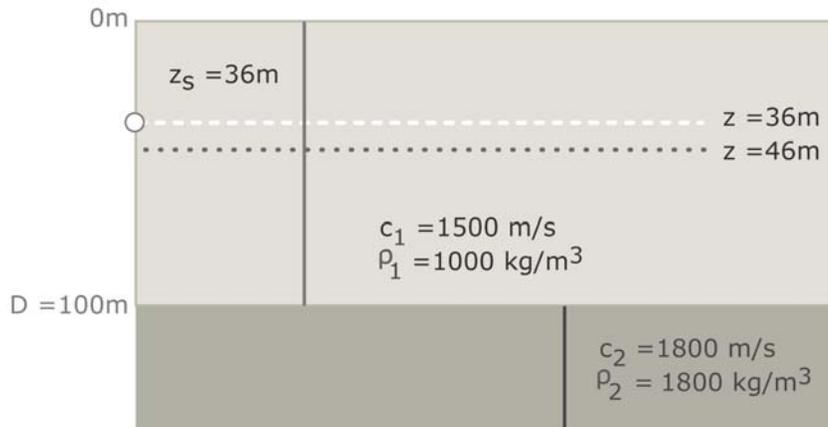


Modal Group Speed - Penetrable Bottom

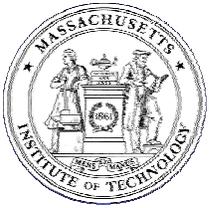
$$u_n = \frac{d\omega}{dk_{rn}} = \frac{k_{rn}}{\omega} \left[\int_0^\infty \frac{\Psi_0^2(z)}{\rho(z) c^2(z)} dz \right]^{-1} \frac{1}{v_n}$$

v_n
 ↑
 Modal Phase Velocity

Pekeris Waveguide



[See Fig 2.28b in Jensen, Kuperman, Porter and Schmidt. *Computational Ocean Acoustics*. New York: Springer-Verlag, 2000.]



Modal Group Speed - Penetrable Bottom

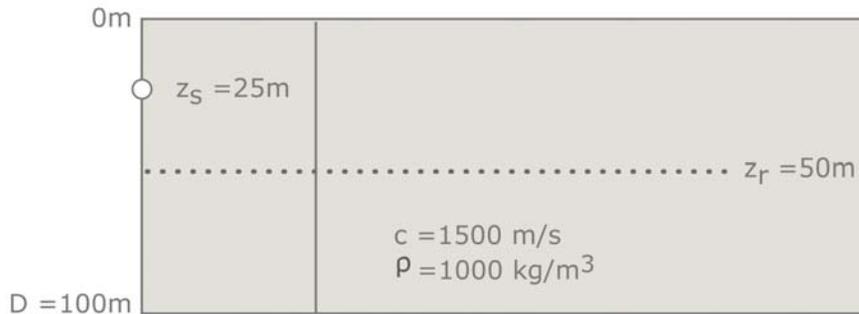
$$u_n = \frac{d\omega}{dk_{rn}} = \frac{k_{rn}}{\omega} \left[\int_0^\infty \frac{\Psi_0^2(z)}{\rho(z) c^2(z)} dz \right]^{-1} \frac{1}{v_n}$$

Isovelocity, Ideal Waveguide

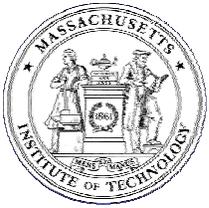
Modal Phase Velocity

$$u_n = \frac{k_{rn} c^2}{\omega} = \frac{c^2}{v_n}$$

Ideal Waveguide

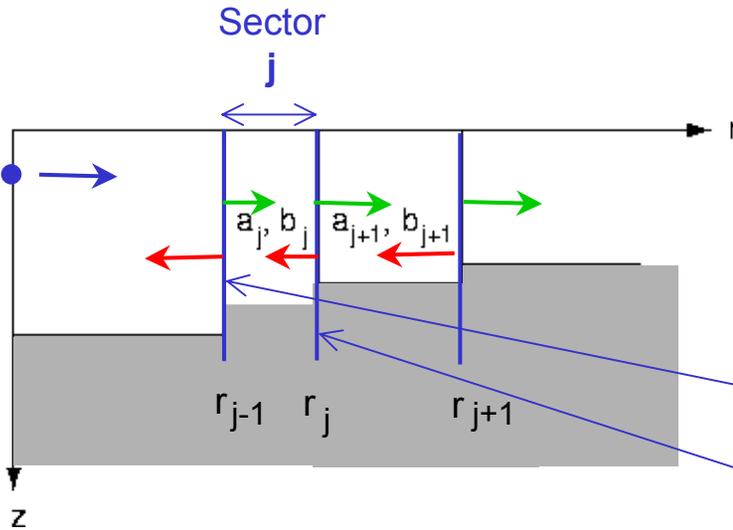


[See Jensen, Fig 2.22]



Normal Modes for Range-Dependent Environments

Coupled Modes



$$p^j(r, z) = \sum_{m=1}^M [a_m^j \widehat{H}1_m^j(r) + b_m^j \widehat{H}2_m^j(r)] \Psi_m^j(z),$$

Normalized Hankel Functions

$$\widehat{H}1_m^j(r) = \frac{H_0^{(1)}(k_{rm}^j r)}{H_0^{(1)}(k_{rm}^j r_{j-1})}, \quad \text{Forward}$$

$$\widehat{H}2_m^j(r) = \frac{H_0^{(2)}(k_{rm}^j r)}{H_0^{(2)}(k_{rm}^j r_j)}, \quad \text{Backward}$$

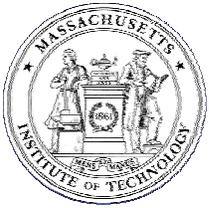
Asymptotic

Forward

$$\widehat{H}1_m^j(r) \simeq H1_m^j(r) = \sqrt{\frac{r_{j-1}}{r}} e^{ik_{rm}^j(r-r_{j-1})},$$

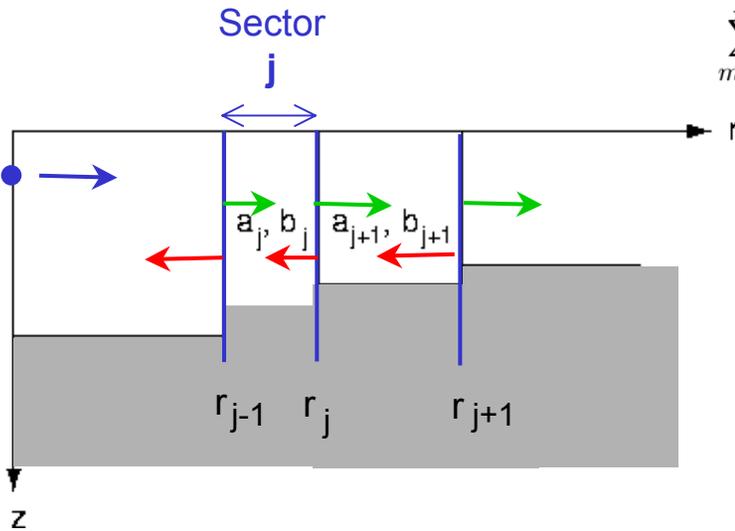
Backward

$$\widehat{H}2_m^j(r) \simeq H2_m^j(r) = \sqrt{\frac{r_j}{r}} e^{ik_{rm}^j(r_j-r)}.$$



Continuity of Pressure

*j*th interface



$$\sum_{m=1}^M (a_m^{j+1} + b_m^{j+1} H 2_m^{j+1}(r_j)) \Psi_m^{j+1}(z) = \sum_{m=1}^M [a_m^j H 1_m^j(r_j) + b_m^j] \Psi_m^j(z).$$

Coupling Operator

$$\int (\cdot) \frac{\Psi_l^{j+1}(z)}{\rho_{j+1}(z)} dz,$$

Orthogonality

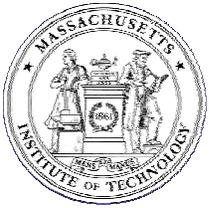
$$\int \frac{\Psi_m^{j+1}(z) \Psi_l^{j+1}(z)}{\rho_{j+1}(z)} dz = \delta_{lm},$$

$$a_l^{j+1} + b_l^{j+1} H 2_l^{j+1}(r_j) = \sum_{m=1}^M [a_m^j H 1_m^j(r_j) + b_m^j] \tilde{c}_{lm}, \quad l = 1, \dots, M,$$

$$\tilde{c}_{lm} = \int \frac{\Psi_l^{j+1}(z) \Psi_m^j(z)}{\rho_{j+1}(z)} dz.$$

Matrix Notation

$$\mathbf{a}^{j+1} + \mathbf{H}_2^{j+1} \mathbf{b}^{j+1} = \tilde{\mathbf{C}}^j (\mathbf{H}_1^j \mathbf{a}^j + \mathbf{b}^j),$$



Continuity of Radial Particle Velocity

$$\frac{1}{\rho_j} \frac{\partial p^j(r, z)}{\partial r} \simeq \frac{1}{\rho_j} \sum_{m=1}^M k_{rm}^j [a_m^j H1_m^j(r) - b_m^j H2_m^j(r)] \Psi_m^j(z).$$

Coupling Operator

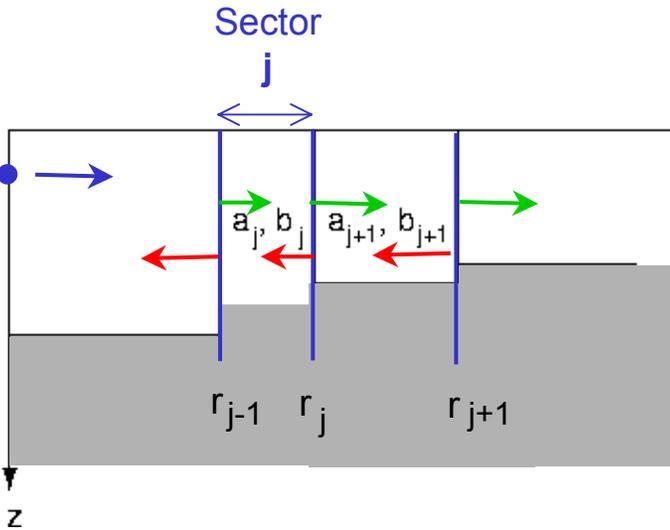
$$\int (\cdot) \Psi_l^{j+1}(z) dz,$$

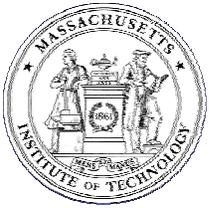
$$a_l^{j+1} - b_l^{j+1} H2_l^{j+1}(r_j) = \sum_{m=1}^M [a_m^j H1_m^j(r_j) - b_m^j] \hat{c}_{lm}, \quad l = 1, \dots, M,$$

$$\hat{c}_{lm} = \frac{k_{rm}^j}{k_{rl}^{j+1}} \int \frac{\Psi_l^{j+1}(z) \Psi_m^j(z)}{\rho_j(z)} dz.$$

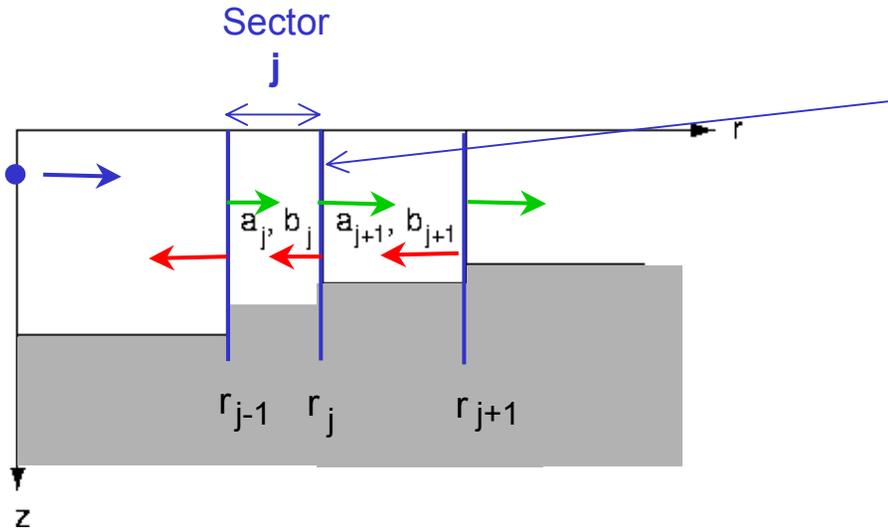
Matrix Notation

$$\mathbf{a}^{j+1} - \mathbf{H}_2^{j+1} \mathbf{b}^{j+1} = \widehat{\mathbf{C}}^j (\mathbf{H}_1^j \mathbf{a}^j - \mathbf{b}^j).$$





Combined Coupling Equations



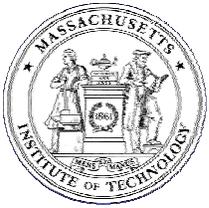
$$\begin{bmatrix} \mathbf{a}^{j+1} \\ \mathbf{b}^{j+1} \end{bmatrix} = \begin{bmatrix} \mathbf{R}_1^j & \mathbf{R}_2^j \\ \mathbf{R}_3^j & \mathbf{R}_4^j \end{bmatrix} \begin{bmatrix} \mathbf{a}^j \\ \mathbf{b}^j \end{bmatrix},$$

$$\mathbf{R}_1^j = \frac{1}{2} (\widetilde{\mathbf{C}}^j + \widehat{\mathbf{C}}^j) \mathbf{H}_1^j,$$

$$\mathbf{R}_2^j = \frac{1}{2} (\widetilde{\mathbf{C}}^j - \widehat{\mathbf{C}}^j),$$

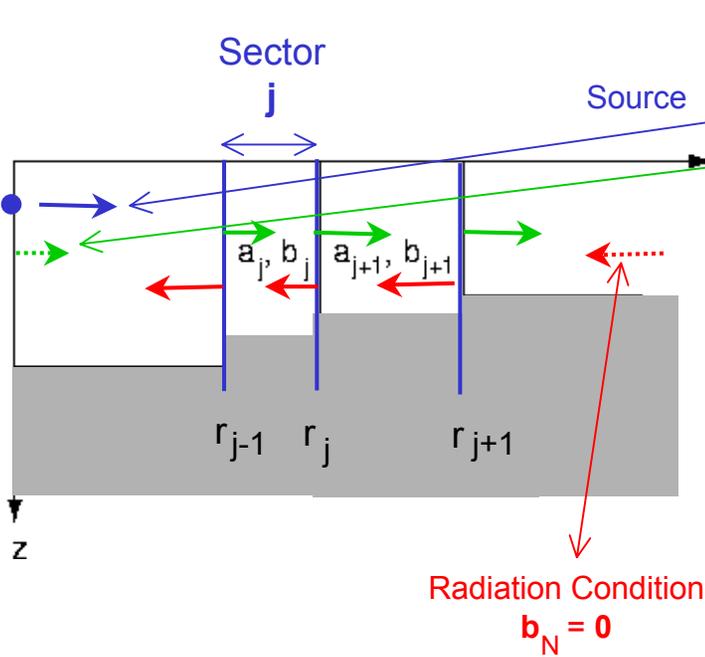
$$\mathbf{R}_3^j = \frac{1}{2} (\widetilde{\mathbf{C}}^j - \widehat{\mathbf{C}}^j) (\mathbf{H}_2^{j+1})^{-1} \mathbf{H}_1^j,$$

$$\mathbf{R}_4^j = \frac{1}{2} (\widetilde{\mathbf{C}}^j + \widehat{\mathbf{C}}^j) (\mathbf{H}_2^{j+1})^{-1}.$$



Initial Condition at Origin

'Rigid' Condition



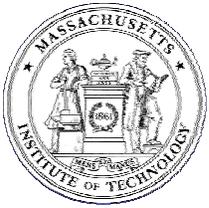
$$a_m^1 = \frac{i}{4\rho(z_s)} \Psi_m(z_s) H_0^{(1)}(k_{rm}^1 r_1) + b_m^1 \frac{H_0^{(1)}(k_{rm}^1 r_1)}{H_0^{(2)}(k_{rm}^1 r_1)}, \quad m = 1, \dots, M.$$

Block-Diagonal Matrix Problem

$$\begin{bmatrix} \mathbf{I} & -\mathbf{D} & 0 \\ \mathbf{R}_1^1 & \mathbf{R}_2^1 & \mathbf{I} & 0 \\ \mathbf{R}_3^1 & \mathbf{R}_4^1 & 0 & \mathbf{I} \\ \dots & \dots & \dots & \dots \\ \mathbf{R}_1^{N-2} & \mathbf{R}_2^{N-2} & \mathbf{I} & 0 \\ \mathbf{R}_3^{N-2} & \mathbf{R}_4^{N-2} & 0 & \mathbf{I} \\ & & \mathbf{R}_1^{N-1} & \mathbf{R}_2^{N-1} & \mathbf{I} \\ & & \mathbf{R}_3^{N-1} & \mathbf{R}_4^{N-1} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{a}^1 \\ \mathbf{b}^1 \\ \mathbf{a}^2 \\ \mathbf{b}^2 \\ \vdots \\ \mathbf{a}^{N-1} \\ \mathbf{b}^{N-1} \\ \mathbf{a}^N \end{bmatrix} = \begin{bmatrix} \mathbf{s} \\ 0 \\ 0 \\ 0 \\ \vdots \\ 0 \\ 0 \\ 0 \end{bmatrix},$$

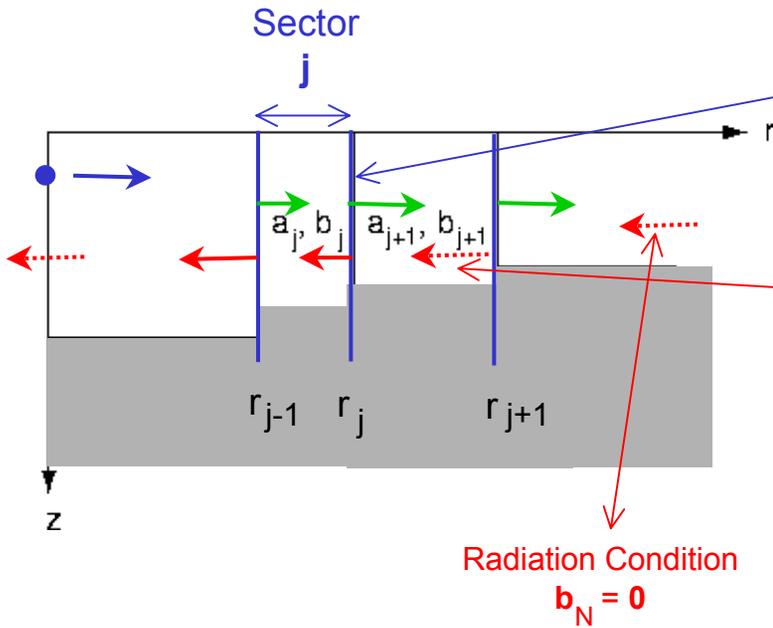
$$d_{ii} = \frac{H_0^{(1)}(k_{ri}^1 r_1)}{H_0^{(2)}(k_{ri}^1 r_1)}$$

$$s_m = \frac{i}{4\rho(z_s)} \Psi_m(z_s) H_0^{(1)}(k_{rm}^1 r_1).$$



One-Way Coupled Modes

Coupling Equations Interface j



$$\begin{bmatrix} \mathbf{a}^{j+1} \\ \mathbf{b}^{j+1} \end{bmatrix} = \begin{bmatrix} \mathbf{R}_1 & \mathbf{R}_2 \\ \mathbf{R}_3 & \mathbf{R}_4 \end{bmatrix} \begin{bmatrix} \mathbf{a}^j \\ \mathbf{b}^j \end{bmatrix} .$$

Ignore Backscatter from next interface:

$$\mathbf{b}^{j+1} = 0$$

Back-Scattered Amplitudes

$$\mathbf{b}^j = -\mathbf{R}_4^{-1} \mathbf{R}_3 \mathbf{a}^j .$$

Forward-Scattered Amplitudes

$$\mathbf{a}^{j+1} = (\mathbf{R}_1 - \mathbf{R}_2 \mathbf{R}_4^{-1} \mathbf{R}_3) \mathbf{a}^j ,$$

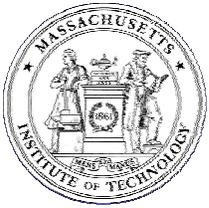
Approximate Single-Scatter Solution

$$\mathbf{a}^{j+1} = \mathbf{R}_1 \mathbf{a}^j .$$

$$\mathbf{R}_1^j = \frac{1}{2} (\widetilde{\mathbf{C}}^j + \widehat{\mathbf{C}}^j) \mathbf{H}_1^j ,$$

Mean of pressure and velocity coupling

Other: $p/(\rho c)^{1/2}$ continuous



Continuous Mode Coupling

Helmholtz Equation

$$\frac{\rho}{r} \frac{\partial}{\partial r} \left(\frac{r}{\rho} \frac{\partial p}{\partial r} \right) + \rho \frac{\partial}{\partial z} \left(\frac{1}{\rho} \frac{\partial p}{\partial z} \right) + \frac{\omega^2}{c^2(r, z)} p = -\frac{\delta(r) \delta(z - z_s)}{2\pi r}.$$

Range Factorization

$$p(r, z) = \sum_m \Phi_m(r) \Psi_m(r, z),$$

Local Modes $\Psi_m(r, z)$

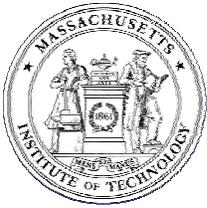
$$\rho(r, z) \frac{\partial}{\partial z} \left[\frac{1}{\rho(r, z)} \frac{\partial \Psi_m(r, z)}{\partial z} \right] + \left[\frac{\omega^2}{c^2(r, z)} - k_{rm}^2(r) \right] \Psi_m(r, z) = 0.$$

Substitution into Helmholtz Equation

$$\sum_m \frac{\rho}{r} \frac{\partial}{\partial r} \left(\frac{r}{\rho} \frac{\partial (\Phi_m \Psi_m)}{\partial r} \right) + \sum_m k_{rm}^2(r) \Phi_m \Psi_m = -\frac{\delta(r) \delta(z - z_s)}{2\pi r},$$

Rearranging Terms

$$\sum_m \left[\frac{\rho}{r} \frac{\partial}{\partial r} \left(\frac{r}{\rho} \frac{\partial \Phi_m}{\partial r} \right) \Psi_m + 2 \frac{\partial \Phi_m}{\partial r} \frac{\partial \Psi_m}{\partial r} + \frac{\rho}{r} \frac{\partial}{\partial r} \left(\frac{r}{\rho} \frac{\partial \Psi_m}{\partial r} \right) \Phi_m \right] + \sum_m k_{rm}^2(r) \Phi_m \Psi_m = -\frac{\delta(r) \delta(z - z_s)}{2\pi r}.$$



Orthogonality Operator

$$\rho = \rho(z) : \int (\cdot) \frac{\Psi_n(r, z)}{\rho(z)} dz ,$$

$$\frac{1}{r} \frac{d}{dr} \left(r \frac{d\Phi_n}{dr} \right) + \sum_m 2B_{mn} \frac{d\Phi_m}{dr} + \sum_m A_{mn} \Phi_m + k_{rn}^2(r) \Phi_n = -\frac{\delta(r) \Psi_n(z_s)}{2\pi r} ,$$

$$A_{mn} = \int \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \Psi_m}{\partial r} \right) \frac{\Psi_n}{\rho} dz ,$$

ODE for
Continuously
Coupled Modes

$$B_{mn} = \int \frac{\partial \Psi_m}{\partial r} \frac{\Psi_n}{\rho} dz .$$

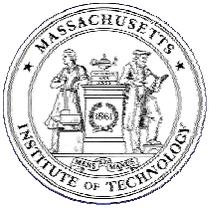
Solved e.g. by FD

$$B_{mn} = -B_{nm}$$

Orthogonality

$$\int \frac{\Psi_m(z) \Psi_n(z)}{\rho(z)} dz = \delta_{mn} ,$$

$$\int \frac{\partial \Psi_m(z)}{\partial r} \frac{\Psi_n(z)}{\rho(z)} dz + \int \frac{\Psi_m(z)}{\rho(z)} \frac{\partial \Psi_n(z)}{\partial r} dz = 0 .$$



Adiabatic Approximation

Decoupled Equations

$$\frac{1}{r} \frac{d}{dr} \left(r \frac{d\Phi_n}{dr} \right) + k_{rn}^2(r) \Phi_n = -\frac{\delta(r) \Psi_n(z_s)}{2\pi r},$$

WKB Approximation

$$\Phi_n(r) \simeq A \frac{e^{i \int_0^r k_{rn}(r') dr'}}{\sqrt{k_{rn}(r)}}.$$

Range-independent Source Condition

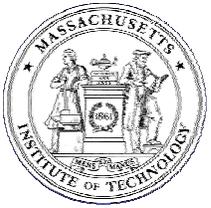
$$A = \frac{i}{\rho(z_s) \sqrt{8\pi r}} e^{-i\pi/4} \Psi_n(z_s).$$

Adiabatic Mode Approximation

$$p(r, z) \simeq \frac{i}{\rho(z_s) \sqrt{8\pi r}} e^{-i\pi/4} \sum_{m=1}^{\infty} \Psi_m(z_s) \Psi_m(r, z) \frac{e^{i \int_0^r k_{rm}(r') dr'}}{\sqrt{k_{rm}(r)}}.$$

Reciprocal Adiabatic Approximation (ad hoc)

$$p(r, z) \simeq \frac{i}{\rho(z_s) \sqrt{8\pi}} e^{-i\pi/4} \sum_{m=1}^{\infty} \Psi_m(z_s) \Psi_m(r, z) \frac{e^{i \int_0^r k_{rm}(r') dr'}}{\sqrt{\int_0^r k_{rm}(r') dr'}}.$$



Warm-Core Eddy Propagation

[Examples: See Jensen Figs 5.17, 5.18, 5.19a]