## 22 APPENDIX 1: MATH FACTS

### 22.1 Vectors

### 22.1.1 Definition

A vector has a dual definition: It is a segment of a a line with direction, or it consists of its projection on a reference system $0 x y z$, usually orthogonal and right handed. The first form is independent of any reference system, whereas the second (in terms of its components) depends directly on the coordinate system. Here we use the second notation, i.e., $\underline{x}$ is meant as a column vector, whose components are found as projections of an (invariant) directed segment on a specific reference system.
We use the overhead arrow to denote a column vector, i.e., a linear segment with a direction. For example, in three-space, we write a vector in terms of its components with respect to a reference system as

$$
\vec{a}=\left\{\begin{array}{l}
2 \\
1 \\
7
\end{array}\right\} .
$$

The elements of a vector have a graphical interpretation, which is particularly easy to see in two or three dimensions.


1. Vector addition:

$$
\begin{gathered}
\vec{a}+\vec{b}=\vec{c} \\
\left\{\begin{array}{l}
2 \\
1 \\
7
\end{array}\right\}+\left\{\begin{array}{l}
3 \\
3 \\
2
\end{array}\right\}=\left\{\begin{array}{l}
5 \\
4 \\
9
\end{array}\right\} .
\end{gathered}
$$

Graphically, addition is stringing the vectors together head to tail.
2. Scalar multiplication:

$$
-2 \times\left\{\begin{array}{l}
2 \\
1 \\
7
\end{array}\right\}=\left\{\begin{array}{r}
-4 \\
-2 \\
-14
\end{array}\right\}
$$

### 22.1.2 Vector Magnitude

The total length of a vector of dimension $m$, its Euclidean norm, is given by

$$
\|\vec{x}\|=\sqrt{\sum_{i=1}^{m} x_{i}^{2}}
$$

This scalar is commonly used to normalize a vector to length one.

### 22.1.3 Vector Dot or Inner Product

The dot product of two vectors is a scalar equal to the sum of the products of the corresponding components:

$$
\vec{x} \cdot \vec{y}=\vec{x}^{T} \vec{y}=\sum_{i=1}^{m} x_{i} y_{i} .
$$

The dot product also satisfies

$$
\vec{x} \cdot \vec{y}=\|\vec{x}\|\|\vec{y}\| \cos \theta
$$

where $\theta$ is the angle between the vectors.

### 22.1.4 Vector Cross Product

The cross product of two three-dimensional vectors $\vec{x}$ and $\vec{y}$ is another vector $\vec{z}, \vec{x} \times \vec{y}=\vec{z}$, whose

1. direction is normal to the plane formed by the other two vectors,
2. direction is given by the right-hand rule, rotating from $\vec{x}$ to $\vec{y}$,
3. magnitude is the area of the parallelogram formed by the two vectors - the cross product of two parallel vectors is zero - and
4. (signed) magnitude is equal to $\|\vec{x}\|\|\vec{y}\| \sin \theta$, where $\theta$ is the angle between the two vectors, measured from $\vec{x}$ to $\vec{y}$.

In terms of their components,

$$
\vec{x} \times \vec{y}=\left|\begin{array}{ccc}
\hat{i} & \hat{j} & \hat{k} \\
x_{1} & x_{2} & x_{3} \\
y_{1} & y_{2} & y_{3}
\end{array}\right|=\left\{\begin{array}{l}
\left(x_{2} y_{3}-x_{3} y_{2}\right) \hat{i} \\
\left(x_{3} y_{1}-x_{1} y_{3}\right) \hat{j} \\
\left(x_{1} y_{2}-x_{2} y_{1}\right) \hat{k}
\end{array}\right\} .
$$

### 22.2 Matrices

### 22.2.1 Definition

A matrix, or array, is equivalent to a set of column vectors of the same dimension, arranged side by side, say

$$
A=\left[\begin{array}{ll}
\vec{a} & \vec{b}
\end{array}\right]=\left[\begin{array}{ll}
2 & 3 \\
1 & 3 \\
7 & 2
\end{array}\right]
$$

This matrix has three rows $(m=3)$ and two columns $(n=2)$; a vector is a special case of a matrix with one column. Matrices, like vectors, permit addition and scalar multiplication. We usually use an upper-case symbol to denote a matrix.

### 22.2.2 Multiplying a Vector by a Matrix

If $A_{i j}$ denotes the element of matrix $A$ in the $i$ 'th row and the $j$ 'th column, then the multiplication $\vec{c}=A \vec{v}$ is constructed as:

$$
c_{i}=A_{i 1} v_{1}+A_{i 2} v_{2}+\cdots+A_{i n} v_{n}=\sum_{j=1}^{n} A_{i j} v_{j}
$$

where $n$ is the number of columns in $A . \vec{c}$ will have as many rows as $A$ has rows $(m)$. Note that this multiplication is defined only if $\vec{v}$ has as many rows as $A$ has columns; they have
consistent inner dimension $n$. The product $\vec{v} A$ would be well-posed only if $A$ had one row, and the proper number of columns. There is another important interpretation of this vector multiplication: Let the subscript : indicate all rows, so that each $A_{: j}$ is the $j$ 'th column vector. Then

$$
\vec{c}=A \vec{v}=A_{: 1} v_{1}+A_{: 2} v_{2}+\cdots+A_{: n} v_{n}
$$

We are multiplying column vectors of $A$ by the scalar elements of $\vec{v}$.

### 22.2.3 Multiplying a Matrix by a Matrix

The multiplication $C=A B$ is equivalent to a side-by-side arrangement of column vectors $C_{: j}=A B_{: j}$, so that

$$
C=A B=\left[\begin{array}{llll}
A B_{: 1} & A B_{: 2} & \cdots & A B_{: k}
\end{array}\right],
$$

where $k$ is the number of columns in matrix $B$. The same inner dimension condition applies as noted above: the number of columns in $A$ must equal the number of rows in $B$. Matrix multiplication is:

1. Associative. $(A B) C=A(B C)$.
2. Distributive. $A(B+C)=A B+A C,(B+C) A=B A+C A$.
3. NOT Commutative. $A B \neq B A$, except in special cases.

### 22.2.4 Common Matrices

Identity. The identity matrix is usually denoted $I$, and comprises a square matrix with ones on the diagonal, and zeros elsewhere, e.g.,

$$
I_{3 \times 3}=\left[\begin{array}{ccc}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]
$$

The identity always satisfies $A I_{n \times n}=I_{m \times m} A=A$.

Diagonal Matrices. A diagonal matrix is square, and has all zeros off the diagonal. For instance, the following is a diagonal matrix:

$$
A=\left[\begin{array}{ccc}
4 & 0 & 0 \\
0 & -2 & 0 \\
0 & 0 & 3
\end{array}\right]
$$

The product of a diagonal matrix with another diagonal matrix is diagonal, and in this case the operation is commutative.

### 22.2.5 Transpose

The transpose of a vector or matrix, indicated by a $T$ superscript results from simply swapping the row-column indices of each entry; it is equivalent to "flipping" the vector or matrix around the diagonal line. For example,

$$
\begin{aligned}
& \vec{a}=\left\{\begin{array}{l}
1 \\
2 \\
3
\end{array}\right\} \longrightarrow \vec{a}^{T}=\left\{\begin{array}{lll}
1 & 2 & 3
\end{array}\right\} \\
& A=\left[\begin{array}{ll}
1 & 2 \\
4 & 5 \\
8 & 9
\end{array}\right] \rightarrow A^{T}=\left[\begin{array}{lll}
1 & 4 & 8 \\
2 & 5 & 9
\end{array}\right] .
\end{aligned}
$$

A very useful property of the transpose is

$$
(A B)^{T}=B^{T} A^{T}
$$

### 22.2.6 Determinant

The determinant of a square matrix $A$ is a scalar equal to the volume of the parallelepiped enclosed by the constituent vectors. The two-dimensional case is particularly easy to remember, and illustrates the principle of volume:

$$
\begin{aligned}
\operatorname{det}(A) & =A_{11} A_{22}-A_{21} A_{12} \\
\operatorname{det}\left(\left[\begin{array}{rr}
1 & -1 \\
1 & 1
\end{array}\right]\right) & =1+1=2 .
\end{aligned}
$$



In higher dimensions, the determinant is more complicated to compute. The general formula allows one to pick a row $k$, perhaps the one containing the most zeros, and apply

$$
\operatorname{det}(A)=\sum_{j=1}^{j=n} A_{k j}(-1)^{k+j} \Delta_{k j}
$$

where $\Delta_{k j}$ is the determinant of the sub-matrix formed by neglecting the $k$ 'th row and the $j$ 'th column. The formula is symmetric, in the sense that one could also target the $k$ 'th column:

$$
\operatorname{det}(A)=\sum_{j=1}^{j=n} A_{j k}(-1)^{k+j} \Delta_{j k}
$$

If the determinant of a matrix is zero, then the matrix is said to be singular - there is no volume, and this results from the fact that the constituent vectors do not span the matrix dimension. For instance, in two dimensions, a singular matrix has the vectors colinear; in three dimensions, a singular matrix has all its vectors lying in a (two-dimensional) plane. Note also that $\operatorname{det}(A)=\operatorname{det}\left(A^{T}\right)$. If $\operatorname{det}(A) \neq 0$, then the matrix is said to be nonsingular.

### 22.2.7 Inverse

The inverse of a square matrix $A$, denoted $A^{-1}$, satisfies $A A^{-1}=A^{-1} A=I$. Its computation requires the determinant above, and the following definition of the $n \times n$ adjoint matrix:

$$
\operatorname{adj}(A)=\left[\begin{array}{ccc}
(-1)^{1+1} \Delta_{11} & \cdots & (-1)^{1+n} \Delta_{1 n} \\
\cdots & \cdots & \cdots \\
(-1)^{n+1} \Delta_{n 1} & \cdots & (-1)^{n+n} \Delta_{n n} .
\end{array}\right]^{T}
$$

Once this computation is made, the inverse follows from

$$
A^{-1}=\frac{\operatorname{adj}(A)}{\operatorname{det}(A)}
$$

If $A$ is singular, i.e., $\operatorname{det}(A)=0$, then the inverse does not exist. The inverse finds common application in solving systems of linear equations such as

$$
A \vec{x}=\vec{b} \longrightarrow \vec{x}=A^{-1} \vec{b} .
$$

### 22.2.8 Trace

The trace of a square matrix is the sum of the diagonals:

$$
\operatorname{tr}(A)=\sum_{i=1}^{n} A_{i i} .
$$

### 22.2.9 Eigenvalues and Eigenvectors

A typical eigenvalue problem is stated as

$$
A \vec{x}=\lambda \vec{x}
$$

where $A$ is an $n \times n$ matrix, $\vec{x}$ is a column vector with $n$ elements, and $\lambda$ is a scalar. We ask for what nonzero vectors $\vec{x}$ (right eigenvectors), and scalars $\lambda$ (eigenvalues) will the equation be satisfied. Since the above is equivalent to $(A-\lambda I) \vec{x}=\overrightarrow{0}$, it is clear that $\operatorname{det}(A-\lambda I)=0$. This observation leads to the solutions for $\lambda$; here is an example for the two-dimensional case:

$$
\begin{aligned}
A & =\left[\begin{array}{ll}
4 & -5 \\
2 & -3
\end{array}\right] \longrightarrow \\
A-\lambda I & =\left[\begin{array}{cc}
4-\lambda & -5 \\
2 & -3-\lambda
\end{array}\right] \longrightarrow \\
\operatorname{det}(A-\lambda I) & =(4-\lambda)(-3-\lambda)+10 \\
& =\lambda^{2}-\lambda-2 \\
& =(\lambda+1)(\lambda-2)
\end{aligned}
$$

Thus, $A$ has two eigenvalues, $\lambda_{1}=-1$ and $\lambda_{2}=2$. Each is associated with a right eigenvector $\vec{x}$. In this example,

$$
\begin{aligned}
\left(A-\lambda_{1} I\right) \vec{x}_{1} & =\overrightarrow{0} \longrightarrow \\
{\left[\begin{array}{cc}
5 & -5 \\
2 & -2
\end{array}\right] \vec{x}_{1} } & =\overrightarrow{0} \longrightarrow \\
\vec{x}_{1} & =\{\sqrt{2} / 2, \quad \sqrt{2} / 2\}^{T} \\
\left(A-\lambda_{2} I\right) \vec{x}_{2} & =\overrightarrow{0} \longrightarrow \\
{\left[\begin{array}{rr}
2 & -5 \\
2 & -5
\end{array}\right] \vec{x}_{2} } & =\overrightarrow{0} \longrightarrow \\
\vec{x}_{2} & = \begin{cases}5 \sqrt{29} / 29, & 2 \sqrt{29} / 29\}^{T}\end{cases}
\end{aligned}
$$

Eigenvectors are defined only within an arbitrary constant, i.e., if $\vec{x}$ is an eigenvector then $c \vec{x}$ is also an eigenvector for any $c \neq 0$. They are often normalized to have unity magnitude, and positive first element (as above). The condition that $\operatorname{rank}\left(A-\lambda_{i} I\right)=\operatorname{rank}(A)-1$ indicates that there is only one eigenvector for the eigenvalue $\lambda_{i}$; more precisely, a unique direction for the eigenvector, since the magnitude can be arbitrary. If the left-hand side rank is less than this, then there are multiple eigenvectors that go with $\lambda_{i}$.
The above discussion relates only the right eigenvectors, generated from the equation $A \vec{x}=$ $\lambda \vec{x}$. Left eigenvectors, defined as $\vec{y}^{T} A=\lambda \vec{y}^{T}$, are also useful for many problems, and can be defined simply as the right eigenvectors of $A^{T}$. $A$ and $A^{T}$ share the same eigenvalues $\lambda$, since they share the same determinant. Example:

$$
\left(A^{T}-\lambda_{1} I\right) \vec{y}_{1}=\overrightarrow{0} \longrightarrow
$$

$$
\begin{aligned}
{\left[\begin{array}{rr}
5 & 2 \\
-5 & -2
\end{array}\right] \vec{y}_{1} } & =\overrightarrow{0} \longrightarrow \\
\vec{y}_{1} & =\{2 \sqrt{29} / 29,-5 \sqrt{29} / 29\}^{T} \\
\left(A^{T}-\lambda_{2} I\right) \vec{y}_{2} & =\overrightarrow{0} \longrightarrow \\
{\left[\begin{array}{rr}
2 & 2 \\
-5 & -5
\end{array}\right] \vec{y}_{2} } & =\overrightarrow{0} \longrightarrow \\
\vec{y}_{2} & =\{\sqrt{2} / 2,-\sqrt{2} / 2\}^{T} .
\end{aligned}
$$

### 22.2.10 Modal Decomposition

For simplicity, we consider matrices that have unique eigenvectors for each eigenvalue. The right and left eigenvectors corresponding to a particular eigenvalue $\lambda$ can be defined to have unity dot product, that is $\vec{x}_{i}^{T} \vec{y}_{i}=1$, with the normalization noted above. The dot products of a left eigenvector with the right eigenvectors corresponding to different eigenvalues are zero. Thus, if the set of right and left eigenvectors, $V$ and $W$, respectively, is

$$
\begin{aligned}
V & =\left[\vec{x}_{1} \cdots \vec{x}_{n}\right], \text { and } \\
W & =\left[\vec{y}_{1} \cdots \vec{y}_{n}\right]
\end{aligned}
$$

then we have

$$
\begin{aligned}
W^{T} V & =I, \text { or } \\
W^{T} & =V^{-1}
\end{aligned}
$$

Next, construct a diagonal matrix containing the eigenvalues:

$$
\Lambda=\left[\begin{array}{ccc}
\lambda_{1} & & 0 \\
& \cdot & \\
0 & & \lambda_{n}
\end{array}\right]
$$

it follows that

$$
\begin{aligned}
A V & =V \Lambda \longrightarrow \\
A & =V \Lambda W^{T} \\
& =\sum_{i=1}^{n} \lambda_{i} \vec{v}_{i} \vec{w}_{i}^{T} .
\end{aligned}
$$

Hence $A$ can be written as a sum of modal components. ${ }^{2}$

[^0]
### 22.2.11 Singular Value

Let $G$ be an $m \times n$ real or complex matrix. The singular value decomposition (SVD) computes three matrices satisfying

$$
G=U \Sigma V^{*}
$$

where $U$ is $m \times m, \Sigma$ is $m \times n$, and $V$ is $n \times n$. The star notation indicates a complex-conjugate transpose (the Hermitian of the matrix). The matrix $\Sigma$ has the form

$$
\Sigma=\left[\begin{array}{cc}
\Sigma_{1} & 0 \\
0 & 0
\end{array}\right], \text { where } \Sigma_{1}=\left[\begin{array}{ccc}
\sigma_{1} & 0 & 0 \\
0 & . & 0 \\
0 & 0 & \sigma_{p}
\end{array}\right]
$$

and $p=\min (m, n)$. Each nonzero entry on the diagonal of matrix $\Sigma_{1}$ is a real, positive singular value, ordered such that $\sigma_{1}>\sigma_{2}>\cdots \sigma_{p}$. Each $\sigma_{i}^{2}$ is an eigenvalue of $G^{H} G$ (or $G G^{H}$ ). The notation is common that $\sigma_{1}=\bar{\sigma}$, the maximum singular value, and $\sigma_{p}=\underline{\sigma}$, the minimum singular value. The auxiliary matrices $U$ and $V$ are unitary, i.e., they satisfy $X^{*}=X^{-1}$. They are defined explicitly: $U$ is the matrix of right eigenvectors of $G G^{H}$, and $V$ is the matrix of right eigenvectors of $G^{H} G$. Like eigenvalues, the singular values of $G$ are related to projections. $\sigma_{i}$ represents the Euclidean size of the matrix $G$ along the $i$ 'th singular vector:

$$
\begin{aligned}
\bar{\sigma} & =\max _{\|x\|=1}\|G x\| \\
\underline{\sigma} & =\min _{\|x\|=1}\|G x\| .
\end{aligned}
$$

Other properties of the singular value include:

- $\bar{\sigma}(A B) \leq \bar{\sigma}(A) \bar{\sigma}(B)$.
- $\bar{\sigma}(A)=\sqrt{\lambda_{\max }\left(A^{*} A\right)}$.
- $\underline{\sigma}(A)=\sqrt{\lambda_{\min }\left(A^{*} A\right)}$.
- $\underline{\sigma}(A)=1 / \bar{\sigma}\left(A^{-1}\right)$.
- $\bar{\sigma}(A)=1 / \underline{\sigma}\left(A^{-1}\right)$.


### 22.3 Laplace Transform

### 22.3.1 Definition

The Laplace transform projects time-domain signals into a complex frequency-domain equivalent. The signal $y(t)$ has transform $Y(s)$ defined as follows:

$$
Y(s)=L(y(t))=\int_{0}^{\infty} y(\tau) e^{-s \tau} d \tau
$$

where $s$ is a complex variable, properly constrained within a region so that the integral converges. $Y(s)$ is a complex function as a result. Note that the Laplace transform is linear, and so it is distributive: $L(x(t)+y(t))=L(x(t))+L(y(t))$. The following table gives a list of some useful transform pairs and other properties, for reference.

The last two properties are of special importance: for control system design, the differentiation of a signal is equivalent to multiplication of its Laplace transform by $s$; integration of a signal is equivalent to division by $s$. The other terms that arise will cancel if $y(0)=0$, or if $y(0)$ is finite.

### 22.3.2 Convergence

We note first that the value of $s$ affects the convergence of the integral. For instance, if $y(t)=e^{t}$, then the integral converges only for $\operatorname{Re}(s)>1$, since the integrand is $e^{1-s}$ in this case. Although the integral converges within a well-defined region in the complex plane, the function $Y(s)$ is defined for all $s$ through analytic continuation. This result from complex analysis holds that if two complex functions are equal on some arc (or line) in the complex plane, then they are equivalent everywhere. It should be noted however, that the Laplace transform is defined only within the region of convergence.

### 22.3.3 Convolution Theorem

One of the main points of the Laplace transform is the ease of dealing with dynamic systems. As with the Fourier transform, the convolution of two signals in the time domain corresponds with the multiplication of signals in the frequency domain. Consider a system whose impulse response is $g(t)$, being driven by an input signal $x(t)$; the output is $y(t)=g(t) * x(t)$. The Convolution Theorem is

$$
y(t)=\int_{0}^{t} g(t-\tau) x(\tau) d \tau \Longleftrightarrow Y(s)=G(s) X(s)
$$

Here's the proof given by Siebert:

$$
\begin{aligned}
Y(s) & =\int_{0}^{\infty} y(t) e^{-s t} d t \\
& =\int_{0}^{\infty}\left[\int_{0}^{t} g(t-\tau) x(\tau) d \tau\right] e^{-s t} d t \\
& =\int_{0}^{\infty}\left[\int_{0}^{\infty} g(t-\tau) h(t-\tau) x(\tau) d \tau\right] e^{-s t} d t \\
& =\int_{0}^{\infty} x(\tau)\left[\int_{0}^{\infty} g(t-\tau) h(t-\tau) e^{-s t} d t\right] d \tau
\end{aligned}
$$

$$
\begin{aligned}
& y(t) \longleftrightarrow Y(s) \\
& \delta(t) \longleftrightarrow \frac{1}{1} \\
& \text { (Impulse) } 1(t) \longleftrightarrow \frac{1}{s} \\
& \text { (Unit Step) } \quad \longleftrightarrow \frac{1}{s^{2}} \\
& \text { (Unit Ramp) } t \longleftrightarrow \frac{1}{s+\alpha} \\
& e^{-\alpha t} \longleftrightarrow \frac{\omega}{s^{2}+\omega^{2}} \\
& \sin \omega t \longleftrightarrow \frac{s}{s^{2}+\omega^{2}} \\
& \cos \omega t \longleftrightarrow \frac{\omega}{(s+\alpha)^{2}+\omega^{2}} \\
& e^{-\alpha t} \sin \omega t \longleftrightarrow \frac{s+\alpha}{(s+\alpha)^{2}+\omega^{2}} \\
& e^{-\alpha t} \cos \omega t \longleftrightarrow \frac{1}{b-a}\left(e^{-a t}-e^{-b t}\right)
\end{aligned} \longleftrightarrow \longleftrightarrow \frac{1}{(s+a)(s+b)}
$$

(Pure Delay) $y(t-\tau) 1(t-\tau) \longleftrightarrow Y(s) e^{-s \tau}$

$$
\text { (Time Derivative) } \frac{d y(t)}{d t} \longleftrightarrow s Y(s)-y(0)
$$

$$
\text { (Time Integral) } \quad \int_{0}^{t} y(\tau) d \tau \longleftrightarrow \frac{Y(s)}{s}+\frac{\int_{0-}^{0+} y(t) d t}{s}
$$

$$
\begin{aligned}
& =\int_{0}^{\infty} x(\tau) G(s) e^{-s \tau} d \tau \\
& =G(s) X(s)
\end{aligned}
$$

where $h(t)$ is the unit step function. When $g(t)$ is the impulse response of a dynamic system, then $y(t)$ represents the output of this system when it is driven by the external signal $x(t)$.

### 22.3.4 Solution of Differential Equations by Laplace Transform

The Convolution Theorem allows one to solve (linear time-invariant) differential equations in the following way:

1. Transform the system impulse response $g(t)$ into $G(s)$, and the input signal $x(t)$ into $X(s)$, using the transform pairs.
2. Perform the multiplication in the Laplace domain to find $Y(s)$.
3. Ignoring the effects of pure time delays, break $Y(s)$ into partial fractions with no powers of $s$ greater than 2 in the denominator.
4. Generate the time-domain response from the simple transform pairs. Apply time delay as necessary.

Specific examples of this procedure are given in a later section on transfer functions.

### 22.4 Background for the Mapping Theorem

The mapping theorem uses concepts from complex analysis, specifically Cauchy's Theorem and the Residue Theorem. References for this section include Ogata and Hildebrand.
First, consider a function of the complex variable $s=u+i v: f(s)$. We say $f(s)$ is analytic in a region $S$ if it has finite derivative and takes only one value at each point $s$ in $S$. Therefore discontinuous or multi-valued functions, e.g., $\sqrt{s}$, are not analytic functions. Polynomials in $s$ are analytic, as are many functions that can be written as a Taylor or Laurent series. An especially important class for control system design is the rational function, a polynomial in $s$ divided by another polynomial in $s$. Rational functions are consequently zero for certain values of $s$, the roots of the numerator, and undefined for other values of $s$, the roots of the numerator, also called the poles.
The integral of interest here is

$$
\int f(s) d s
$$

taken on a path in the $s$-plane. A closed path in the complex $s$-plane leads to a closed path in the $f(s)$ plane, but more than one point in the $s$ plane can map to a single $f(s)$-plane point, so the number of complete loops may not be the same.
The usual rules of integration apply in complex analysis, so that, insofar as the antiderivative of $f(s)$, denoted $F(s)$ exists, and $f(s)$ is analytic on the path, we have

$$
\int_{s_{1}}^{s_{2}} f(s) d s=F\left(s_{2}\right)-F\left(s_{1}\right)
$$

It appears that this integral is zero for a closed path, since $s_{1}=s_{2}$. Indeed, Cauchy's theorem states that it is, provided that $f(s)$ is analytic on the path, and everywhere within the region enclosed. This latter condition results from the following observation. Consider the function $f(s)=\alpha s^{n}$; on a circular path of radius $r$, we have $s=r e^{i \theta}$, and thus

$$
\begin{aligned}
\int f(s) d s & =i \alpha r^{n+1} \int_{0}^{2 \pi} e^{i(n+1) \theta} d \theta \\
& =i \alpha r^{n+1} \int_{0}^{2 \pi}[\cos (n+1) \theta+i \sin (n+1) \theta] d \theta
\end{aligned}
$$

The second integral is clearly zero for all $n$, whereas the first is zero except in the case of $n=-1$, for which we obtain

$$
\int f(s) d s=i \alpha 2 \pi
$$

This result does not depend on $r$, and so applies to a vanishingly small circle around the point $s=0$. It can be shown also that the result holds for any closed path around the simple pole at $s=0$, which characterizes the function. The residue theorem is an extension to an arbitrary number of simple poles which are enclosed by a path:

$$
\int f(s) d s=i 2 \pi \sum \alpha_{i}
$$

The constants $\alpha_{i}$ are the residues associated with the poles, i.e.,

$$
f(s)=\frac{\alpha_{1}}{s-p_{1}}+\frac{\alpha_{2}}{s-p_{2}}+\ldots
$$

We show in another section how any strictly proper rational function (that is, the polynomial order in the numerator is less than that of the denominator) in $s$ can be put into this form. The connection between Cauchy's theorem and the residue theorem is indicated in the figure below. Here, the integral inside the closed path is zero because it excludes the three simple poles. Without the cuts however, the integral over the outermost curve would be equal to the summation of residues for the poles within (multiplied by $i 2 \pi$ ). Note that it is only the terms with $\left(s-a_{i}\right)^{-1}$, i.e., simple poles at $a_{i}$, that generate nonzero components.


Looking toward the mapping theorem now, consider the function

$$
f(s)=\frac{\left(s-z_{1}\right)^{m_{1}}\left(s-z_{2}\right)^{m_{2}} \cdots}{\left(s-p_{1}\right)^{n_{1}}\left(s-p_{2}\right)^{n_{2}} \cdots}
$$

Working through some algebra, it is straightforward that

$$
f^{\prime}(s) / f(s)=\frac{m_{1}}{s-z_{1}}+\frac{m_{2}}{s-z_{2}}+\ldots-\frac{n_{1}}{s-p_{1}}-\frac{n_{2}}{s-p_{2}} \ldots
$$

resulting in

$$
\int f^{\prime}(s) / f(s) d s=i 2 \pi(Z-P)
$$

where $Z=m_{1}+m_{2}+\ldots$ is the number of zeros in $f$, and $P=n_{1}+n_{2}+\ldots$ is the number of zeros. The mapping theorem comes from now looking in detail at the integral above:

$$
\begin{aligned}
f(s) & =|f(s)| e^{i \theta(s)} \\
f^{\prime}(s) & =\frac{d|f(s)|}{d s} e^{i \theta(s)}+i|f(s)| e^{i \theta(s)} \frac{d \theta(s)}{d s} \\
f^{\prime}(s) / f(s) & =\frac{1}{|f(s)|} \frac{d|f(s)|}{d s}+i \frac{d \theta(s)}{d s} \\
& =\frac{d \log |f(s)|}{d s}+i \frac{d \theta(s)}{d s},
\end{aligned}
$$

Considering the integral of $f^{\prime}(s) / f(s)$ over a closed contour in the $s$-plane, we take advantage of the fact that exact differentials $d \log |f(s)|$ and $d \theta(s)$ have been found. Both terms pertain to the $f(s)$ plane, not the $f^{\prime}(s) / f(s)$ plane. The first integral is zero,

$$
\int_{s_{1}}^{s_{2}} d \log |f|=0
$$

since $\log |f(s)|$ has to be the same at the start and end of the closed contour. Taking $\theta(s)$ counterclockwise through angle $n 2 \pi$ results in the second term

$$
\int_{0}^{n 2 \pi} i d \theta=i n 2 \pi
$$

As noted above, a single circuit in the s-plane may or may not map to a single circuit in the $f(s)$-plane, so $n$ depends directly on the function $f(s)$, and is not necessarily one. Assembling the results, we have

$$
i 2 \pi(Z-P)=i n 2 \pi \longrightarrow Z-P=n,
$$

which is the mapping theorem. In words, it states that the number of zeros minus the number of poles enclosed in a contour in the $s$-plane is equal to the number of counterclockwise encirclements of the origin, in the $f(s)$ plane.


[^0]:    ${ }^{2}$ By carrying out successive multiplications, it can be shown that $A^{k}$ has its eigenvalues at $\lambda_{i}^{k}$, and keeps the same eigenvectors as $A$.

