2.43 ADVANCED THERMODYNAMICS

Spring Term 2024 LECTURE 11

Room 3-442 Tuesday, March 12, 2:30pm - 4:30pm

Instructor: Gian Paolo Beretta <u>beretta@mit.edu</u> Room 3-351d **Review of basic concepts:**

Exergies and first and second law efficiencies

in energy conversion technologies

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Review of basic concepts: Exergy in heating and cooling bulk flows



Review of basic concepts: Minimum exergy for low temperature heating a liquid or ideal-gas flow

Ratio between exergy and energy content for a liquid flow:

For example, for the heating of pressurized water from 60°C to 120°C (in an environment at 300 K)

$$\frac{\Delta \dot{E} x_{12}}{\dot{Q}} = 1 - \frac{T_a \ln(T_2/T_1)}{T_2 - T_1} = 1 - \frac{300 \ln(393/333)}{60} = 0.172 \qquad \frac{1}{0.172} = 5.8$$

For the heating of water from 20°C to 60°C (in an environment at 300 K)

$$\frac{\Delta \dot{E} x_{12}}{\dot{Q}} = 1 - \frac{T_a \ln(T_2/T_1)}{T_2 - T_1} = 1 - \frac{300 \ln(333/293)}{40} = 0.040 \qquad \begin{array}{c} \text{COP}_{\text{max}} = \\ \frac{1}{0.04} = 25 \end{array}$$

Review of basic concepts: Power-plant philosophy of best available flame-based fuel-to-power conversion technology



Overall thermodynamic efficiency

64%

$$\eta_{\rm II-overall} = \frac{\dot{W}_{\rm electric}}{\dot{E}x_{\rm fuel}} = \eta_{\rm II-combustion} \times \eta_{\rm II-thermal-cycle} \approx 64\%$$

Review of basic concepts: one way to get around the inherent irreversibility of flames is by oxidating the fuel in fuel cells



Integrating SOFC in Thermal Cycles (without CO2 sequestration) may yield: Fuel-Cell Rankine Cycle, 72% Fuel-Cell Combined Brayton-Rankine Cycle, 75% Fuel-Cell Regenerative Brayton Cycle, 76%

Separate production of electricity (simple thermoelectric generation)



Note that in this particular case, $\eta_{II} \approx \eta_{I}$ is a consequence of the fact that

$$\dot{E}x_{\text{fuel}} \approx \dot{Q}_{\text{fuel}} \approx \dot{m} \text{ LHV}$$

Separate production of heat at a fixed temperature



- The value of $\eta_{\rm I}$ for a boiler for an industrial production of process steam is around 95%. But remember that this does not mean that there is only a 5% margin of improvement. In fact..
- The value of η_{II} for an industrial boiler that produces steam at 120°C is equal to 0.95*(1-283/393) = 27% (for T_o = 283 K), so there is a large margin for improvement!



- The value of η_{II} for an 85% efficient domestic boiler which heats water from 20 to 60°C is around 4% (for T_o = 298 K) or 11% (for T_o = 273 K).
- For a 95% efficient industrial boiler that heats pressurized water from 20 to 120°C, η_{II} is around 12% (for T_o = 298 K) or 19% (for T_o = 273 K).
- There are ample margins for improvement, obtained for example with cogeneration.

Combined production of heat and mechanical energy (cogeneration)



Combined production of heat and mechanical energy (cogeneration)



$$\eta_{\rm I} = \frac{\dot{W} + \dot{Q}_B}{\dot{Q}_{\rm fuel}}$$

$$\eta_{\text{II}} = \frac{\dot{W} + \dot{Q}_B \left[1 - T_o \frac{s_2 - s_1}{h_2 - h_1} \right]}{\dot{E} x_{\text{fuel}}}$$

In the case of heating of water or air:

$$\eta_{\text{II}} = \frac{\dot{W} + \dot{Q}_B \left[1 - \frac{T_o \ln(T_2/T_1)}{T_2 - T_1} \right]}{\dot{E} x_{\text{fuel}}}$$

Separate production of chilled water (simple industrial refrigeration)



The value of COP for industrial refrigerators used for air conditioning, which cool water from 12 to 7°C, can be around 5, so (for $T_o = 303$ K) η_{II} is around 36%.



cooling

Review of basic concepts: Typical (efficient and inefficient) uses of fossil fuels Try to fill this table for your own country!

Method:	SP = Separate production		CH combin and p	P = ed heat oower	HYB = hybrid fuel and renewable		
Second-law efficiency of the conversion of fuel energy to:	Ave	BAT	Ave	BAT	Ave	BAT	
Electrical power	45%?	65%?	55%?	70%?	48%?	?	
Residential heat (20-60°C @T _a =0°C)	11%?		14%?				
Residential heat (20-60°C @T _a =25°C)	4%?		5%?				
Industrial heat (60-120°C @T _a =0°C)	23%?		32%?				
Industrial heat (60-120°C @T _a =25°C)	17%?		23%?				
 (other energy products)							

Review of basic concepts:

The learning curve of fuel-to-power conversion technologies



 η_{II}

Thermodynamic efficiency of the best-available mature technology for primary energy conversion to work or electricity

Methods for the ALLOCATION

of energy consumption and CO₂ production in combined heat and power (CHP) production

and of heat and/or power production in hybrid multi-resource facilities

Allocation problem in Heat&Power Cogeneration: α_W , α_Q , partial efficiencies, and PES



Allocation problem in CHP: Incremental Electricity-Centered Allocation



Allocation problem in CHP: Separate Production Reference Allocation



Allocation problem in CHP: Choice of reference values for η_W^{ref} and COP_Q^{ref}

SPR method:

Fixed values set by some local Authority



Exergy method:

Fixed values set by Thermodynamics

Effectively takes as references the REVERSIBLE heat engines!

 $\eta_{\rm W}^{\rm ref} \approx 1$

$$\operatorname{COP}_{Q}^{\operatorname{ref}} = 1 / \left(1 - \frac{T_{\operatorname{env}}}{T_{Q}} \right)$$

Allocation problem in CHP: "fair" reference values in a given local area

STALPR Method*: Self-Tuned-Average-Local-Productions-Reference

Adopt reference efficiencies



that are representative of the actual average efficiencies

 η_{W}^{loc} COP_Q^{loc}

of the energy production portfolio (typically the local area where the cogenerator itself is located) with which the resulting efficiencies of the cogenerator





* G.P. Beretta, P. Iora, and A.F. Ghoniem, Energy, Vol. 44, pp. 1107-1120 (2012)

Allocation Example in CHP: a comparison between allocation methods





0/]	IEC	Ex	ergy	STALPR		
70	η^{chp}_{W}	COP _Q ^{chp}	η^{chp}_{W}	$\operatorname{COP}_Q^{\operatorname{chp}}$	η^{chp}_W	COP _Q ^{chp}	
CC	52	354	50	234	69.2	92.0	
BPST	52	101	29	137	54.4	95.8	

Allocation problem in Hybrid Facilities: $\beta_{\rm F}$, $\beta_{\rm Q}$, partial efficiencies, and PES



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Allocation problem in Hybrid Facilities: Incremental Fossil-Centered Allocation



Allocation problem in Hybrid Facilities: Separate Production Reference Allocation



Allocation problem in Hybrid Facilities: Choice of reference values η_F^{ref}

SPR method: Fixed values set by some local Authority Hybrid facility **Reference SP facilities** $\eta_F^{ref}F$ **Fossil-only** F $\beta_F W$ $\frac{\eta_{\rm F}^{\rm ref} F}{\eta_{\rm F}^{\rm ref} F + \eta_{\rm F}^{\rm ref} R}$ power production $\eta_{\rm F}^{\rm ref} R$ $\eta_{\rm F}^{\rm ref} F + r'$ **SPR**

Exergy method:

R



Effectively takes as references the REVERSIBLE heat engines!

$$\eta_{\rm F}^{\rm ref} \approx 1$$

$$\eta_{\rm R}^{\rm ref} \approx 0.93$$

R

Solar-only

power

production

 $\eta_R^{\text{ref}}R$

 $a \eta_R^{ref}$

Allocation problem in Hybrid Facilities: "fair" reference values in a given local area

STALPR Method*: Self-Tuned-Average-Local-Productions-Reference

Adopt reference efficiencies



that are representative of the actual average efficiencies



of the energy production portfolio (typically the local area where the hybrid plant itself is located) with which the resulting efficiencies of the hybrid plant





* G.P. Beretta, P. Iora, and A.F. Ghoniem, Energy, Vol. 60, pp. 344-360 (2013)

Allocation Example in Hybrid Facilities: a comparison between allocation methods





%	FC		Exergy		SPR (38%,15.3%)			STALRP				
	η_{F}^{hyb}	η_R^{hyb}	β_R^{hyb}	η_{F}^{hyb}	η_R^{hyb}	β_R^{hyb}	$\eta_{\rm F}^{\rm hyb}$	η_R^{hyb}	β_R^{hyb}	$\eta_{\rm F}^{\rm hyb}$	η_R^{hyb}	β_R^{hyb}
	38	204	25.2	42	35.5	17.4	46.2	18.6	9.11	46.1	18.7	9.19

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Part II: chemical potentials and multicomponent equilibria Contents:

- Modeling ideal and nonideal gas mixtures and solutions
- Mutual equilibrium across semi-permeable membranes
- Minimum work of separation, maximum work of mixing
- Osmotic pressure
- Liquid-vapor and liquid-liquid phase equilibria
- Metastable states and spinodal decomposition
- Modeling chemically reacting mixtures
- Chemical equilibrium
- Chemical kinetics standard model
- Electrochemical potentials



Stable-equilibrium properties of MIXTURES (within the simple-system model)

Partial properties Gibbs-Dalton ideal mixtures Mixing and separation Osmotic pressure and blue energy Stratification

Summary of SES relations: (small systems) (Mixtures) Specific properties in general depend on *n* ...

$$\begin{split} \overline{Eu} &= E - TS + pV - \boldsymbol{\mu} \cdot \boldsymbol{n} = F + pV - \boldsymbol{\mu} \cdot \boldsymbol{n} = G - \boldsymbol{\mu} \cdot \boldsymbol{n} = H - TS - \boldsymbol{\mu} \cdot \boldsymbol{n} \\ dEu &= -S \, dT + V \, dp - \boldsymbol{n} \cdot d\boldsymbol{\mu} \quad \Rightarrow \quad Eu = Eu(T, p, \boldsymbol{\mu}) \\ eu &= \frac{Eu}{n} = e - Ts + pv - \boldsymbol{\mu} \cdot \boldsymbol{y} = g - \boldsymbol{\mu} \cdot \boldsymbol{y} = f + pv - \boldsymbol{\mu} \cdot \boldsymbol{y} = h - Ts - \boldsymbol{\mu} \cdot \boldsymbol{y} \\ \frac{deu = \frac{1}{n} dEu - Eu \frac{dn}{n^2} = -s \, dT + v \, dp - \boldsymbol{y} \cdot d\boldsymbol{\mu} - eu \frac{dn}{n} \quad \Rightarrow \quad eu = eu(T, p, \boldsymbol{\mu}, n) = \frac{n_0}{n} eu(T, p, \boldsymbol{\mu}, n_0) \\ \overline{dE} = T \, dS - p \, dV + \boldsymbol{\mu} \cdot d\boldsymbol{n} \quad \Rightarrow \quad E = E(S, V, n\boldsymbol{y}) \quad \Rightarrow \quad \mu_i = \left(\frac{\partial E}{\partial n_i}\right)_{S,V,\boldsymbol{n}_i^i} = \mu_i(S, V, \boldsymbol{n}) \\ \frac{de = \frac{1}{n} dE - E \frac{dn}{n^2} = T \, ds - p \, dv + \boldsymbol{\mu} \cdot d\boldsymbol{y} - eu \frac{dn}{n} \quad \Rightarrow \quad e = e(s, v, \boldsymbol{y}, n) \\ \overline{dF} = -S \, dT - p \, dV + \boldsymbol{\mu} \cdot d\boldsymbol{n} \quad \Rightarrow \quad F = E - TS \\ = F(T, V, n\boldsymbol{y}) \quad \Rightarrow \quad \mu_i = \left(\frac{\partial F}{\partial n_i}\right)_{T,V,\boldsymbol{n}_i^\prime} = \mu_i(T, V, \boldsymbol{n}) \\ \frac{df = \frac{1}{n} dF - F \frac{dn}{n^2} = -s \, dT + p \, dv + \boldsymbol{\mu} \cdot d\boldsymbol{y} - eu \frac{dn}{n} \quad \Rightarrow \quad f = f(T, v, \boldsymbol{y}, n) \\ \overline{dG} = -S \, dT + V \, dp + \boldsymbol{\mu} \cdot d\boldsymbol{n} \quad \Rightarrow \quad G = E - TS + pV \\ = G(T, p, n\boldsymbol{y}) \quad \Rightarrow \quad \mu_i = \left(\frac{\partial G}{\partial n_i}\right)_{T,p,\boldsymbol{n}_i^\prime} = \mu_i(T, p, \boldsymbol{n}) \\ \frac{dg = \frac{1}{n} dG - G \frac{dn}{n^2} = -s \, dT + v \, dp + \boldsymbol{\mu} \cdot d\boldsymbol{y} - eu \frac{dn}{n} \quad \Rightarrow \quad g = g(T, p, \boldsymbol{y}, n) \\ \overline{dH = T \, dS + V \, dp + \boldsymbol{\mu} \cdot d\boldsymbol{n}} \quad \Rightarrow \quad H = E + pV \\ = H(S, p, n\boldsymbol{y}) \quad \Rightarrow \quad \mu_i = \left(\frac{\partial H}{\partial n_i}\right)_{S,p,\boldsymbol{n}_i^\prime} = \mu_i(S, p, \boldsymbol{n}) \\ \frac{dh = \frac{1}{n} dH - H \frac{dn}{n^2} = T \, ds + v \, dp + \boldsymbol{\mu} \cdot d\boldsymbol{y} - eu \frac{dn}{n} \quad \Rightarrow \quad h = h(s, p, \boldsymbol{y}, n) \end{aligned}$$

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Summary of SES relations: (large *n* limit: simple-system approximation) (Mixtures) ... but that dependence vanishes for large *n*

$$\begin{split} \overline{Eu} &= 0 \quad \Rightarrow \quad \overline{E} - TS + pV - \mu \cdot n = 0 \quad \text{Euler relation} \\ &\Rightarrow \quad \mu \cdot n = G = H - TS = F + pV = E - TS + pV \\ dEu &= -S \, dT + V \, dp - n \cdot d\mu = 0 \quad \text{Gibbs-Duhem relation} \\ eu &= 0 = e - Ts + pv - \mu \cdot y \quad \Rightarrow \quad \mu \cdot y = g = h - Ts = f + pv = e - Ts + pv \\ dE &= T \, dS - p \, dV + \mu \cdot dn \quad \Rightarrow \quad E = E(S, V, ny) \quad \Rightarrow \quad \mu_i = \left(\frac{\partial E}{\partial n_i}\right)_{S,V,n'_i} = \mu_i(S, V, n) \\ \frac{de}{dF} &= -S \, dT - p \, dV + \mu \cdot dn \quad \Rightarrow \quad F = E - TS \\ = F(T, V, ny) \quad \Rightarrow \quad \mu_i = \left(\frac{\partial F}{\partial n_i}\right)_{T,V,n'_i} = \mu_i(T, V, n) \\ \frac{df}{dF} &= \frac{1}{n} dF - F \frac{dn}{n^2} = -s \, dT + p \, dv + \mu \cdot dy \quad \Rightarrow \quad f = f(T, v, y) \\ \frac{df}{dG} &= -S \, dT + V \, dp + \mu \cdot dn \quad \Rightarrow \quad G = E - TS + pV \\ &= G(T, p, ny) \quad \Rightarrow \quad \mu_i = \left(\frac{\partial G}{\partial n_i}\right)_{T,p,n'_i} = \mu_i(T, p, n) \\ \frac{dg}{dF} &= \frac{1}{n} dG - G \frac{dn}{n^2} = -s \, dT + v \, dp + \mu \cdot dy \quad \Rightarrow \quad g = g(T, p, y) \\ \frac{dH = T \, dS + V \, dp + \mu \cdot dn}{dH = H (S, p, ny)} \quad \Rightarrow \quad \mu_i = \left(\frac{\partial H}{\partial n_i}\right)_{S,p,n'_i} = \mu_i(S, p, n) \\ \frac{dh = \frac{1}{n} dH - H \frac{dn}{n^2} = T \, ds + v \, dp + \mu \cdot dy \quad \Rightarrow \quad h = h(s, p, y) \end{split}$$

Summary of SES relations: (small systems) (Pure substances)

Specific properties in general depend on *n*

$$\begin{split} Eu &= E - TS + pV - \mu n = F + pV - \mu n = G - \mu n = H - TS - \mu n \\ dEu &= -S \, dT + V \, dp - n \, d\mu \quad \Rightarrow \quad Eu = Eu(T, p, \mu) \\ eu &= \frac{Eu}{n} = e - Ts + pv - \mu = g - \mu = f + pv - \mu = h - Ts - \mu \\ deu &= \frac{1}{n} dEu - Eu \frac{dn}{n^2} = -s \, dT + v \, dp - d\mu - eu \frac{dn}{n} \quad \Rightarrow \quad eu = eu(T, p, \mu, n) = \frac{n_0}{n} eu(T, p, \mu, n_0) \\ dE &= T \, dS - p \, dV + \mu \, dn \quad \Rightarrow \quad E = E(S, V, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial E}{\partial n}\right)_{S,V} = \mu(S, V, n) \\ dE &= -S \, dT - p \, dV + \mu \, dn \quad \Rightarrow \quad F = E - TS \\ dF &= -S \, dT - p \, dV + \mu \, dn \quad \Rightarrow \quad F = E - TS \\ dF &= -S \, dT - p \, dV + \mu \, dn \quad \Rightarrow \quad F = E - TS \\ eF &= F(T, V, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial F}{\partial n}\right)_{T,V} = \mu(T, V, n) \\ dG &= -S \, dT + V \, dp + \mu \, dn \quad \Rightarrow \quad G = E - TS + pV \\ = G(T, p, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial G}{\partial n}\right)_{T,p} = \mu(T, p, n) \\ dg &= \frac{1}{n} dG - G \frac{dn}{n^2} = -s \, dT + v \, dp - eu \frac{dn}{n} \quad \Rightarrow \quad g = g(T, p, n) \\ dH &= T \, dS + V \, dp + \mu \, dn \quad \Rightarrow \quad H = E + pV \\ = H(S, p, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial H}{\partial n}\right)_{S,p} = \mu(S, p, n) \\ dh &= \frac{1}{n} dH - H \frac{dn}{n^2} = T \, ds + v \, dp - eu \frac{dn}{n} \quad \Rightarrow \quad h = h(s, p, n) \end{aligned}$$

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Summary of SES relations: (large *n* limit: simple-system approximation) (Pure substances) ... but that dependence vanishes for large *n*

$$\begin{array}{l} \overline{Eu} = 0 \quad \Rightarrow \quad \overline{E} - TS + pV - \mu n = 0 \quad \mbox{Euler relation} \\ \Rightarrow \quad \mu n = G = H - TS = F + pV = E - TS + pV \\ dEu = -S \, dT + V \, dp - n \, d\mu = 0 \quad \Rightarrow \quad d\mu = -s \, dT + v \, dp \quad \mbox{Gibbs-Duhem relation} \\ eu = 0 = e - Ts + pv - \mu \quad \Rightarrow \quad \mu = g = h - Ts = f + pv = e - Ts + pv = \mu(T, p) \\ dE = T \, dS - p \, dV + \mu \, dn \quad \Rightarrow \quad E = E(S, V, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial E}{\partial n}\right)_{S,V} = \mu(S, V, n) \\ \hline de = \frac{1}{n} dE - E \frac{dn}{n^2} = T \, ds - p \, dv \quad \Rightarrow \quad e = e(s, v) \\ dF = -S \, dT - p \, dV + \mu \, dn \quad \Rightarrow \quad F = E - TS \\ = F(T, V, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial F}{\partial n}\right)_{T,V} = \mu(T, V, n) \\ \hline df = \frac{1}{n} dF - F \frac{dn}{n^2} = -s \, dT + p \, dv \quad \Rightarrow \quad f = f(T, v) \\ \hline dG = -S \, dT + V \, dp + \mu \, dn \quad \Rightarrow \quad G = E - TS + pV \\ = G(T, p, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial G}{\partial n}\right)_{T,p} = \mu(T, p, n) \\ \hline dg = \frac{1}{n} dG - G \frac{dn}{n^2} = -s \, dT + v \, dp \quad \Rightarrow \quad g = g(T, p) \quad \Rightarrow \quad \mu = \mu(T, p) \\ \hline dH = T \, dS + V \, dp + \mu \, dn \quad \Rightarrow \quad H = E + pV \\ = H(S, p, n) \quad \Rightarrow \quad \mu = \left(\frac{\partial H}{\partial n}\right)_{S,p} = \mu(S, p, n) \\ \hline dh = \frac{1}{n} dH - H \frac{dn}{n^2} = T \, ds + v \, dp \quad \Rightarrow \quad h = h(s, p) \end{array}$$

Experimental measurement of SES properties of mixtures:

Partial pressures to measure chemical potentials



Recall the definition of partial pressure of constituent i in the mixture

 $\mu_i(T, p, \boldsymbol{n}) = \mu_{ii}(T, p_{ii}, n_{ii}) \implies p_{ii} = p_{ii}(T, p, \boldsymbol{n}, n_{ii}) \stackrel{n \text{ large}}{\longrightarrow} p_{ii} = p_{ii}(T, p, \boldsymbol{y})$ For large *n*, within the SS model, recall that for the pure substance $\mu = g$, hence, $\mu_i = \mu_i(T, p, \boldsymbol{y}) = \mu_{ii}(T, p_{ii}) = g_{ii}(T, p_{ii})$ **Experimental measurement of SES properties of mixtures:**

Partial properties from the chemical potentials

$$\begin{aligned} G &= E - TS + pV = G(T, p, \boldsymbol{n}) & \mathrm{d}G = -S \,\mathrm{d}T + V \,\mathrm{d}p + \boldsymbol{\mu} \cdot \mathrm{d}\boldsymbol{n} \\ \left(\frac{\partial G}{\partial T}\right)_{T, p, \boldsymbol{n}} &= -S(T, p, \boldsymbol{n}) \quad \left(\frac{\partial G}{\partial p}\right)_{T, p, \boldsymbol{n}} = V(T, p, \boldsymbol{n}) \quad \mu_i = \left(\frac{\partial G}{\partial n_i}\right)_{T, p, \boldsymbol{n}'_i} = \mu_i(T, p, n\boldsymbol{y}) \\ \mathrm{d}\mu_i &= \underbrace{\left(\frac{\partial \mu_i}{\partial T}\right)_{p, \boldsymbol{n}}}_{-s_i} \mathrm{d}T + \underbrace{\left(\frac{\partial \mu_i}{\partial p}\right)_{T, \boldsymbol{n}}}_{V_i} \mathrm{d}p + \sum_{j=1}^r \underbrace{\left(\frac{\partial \mu_i}{\partial n_j}\right)_{T, p, \boldsymbol{n}'_j}}_{\mu_{i,j} = \mu_{i,j}} \mathrm{d}n_j = -s_i \,\mathrm{d}T + \underbrace{v_i \,\mathrm{d}p + \sum_{j=1}^r \mu_{i,j} \,\mathrm{d}n_j}_{\mathrm{d}\mu_i|_T} \end{aligned}$$

where we define

 $-\left(\frac{\partial \mu_i}{\partial T}\right)_{n,\boldsymbol{n}} = -\left(\frac{\partial^2 G}{\partial T \partial n_i}\right)_{n,\boldsymbol{n}'} = \left(\frac{\partial S}{\partial n_i}\right)_{T,n,\boldsymbol{n}'} = s_i(T,p,n\boldsymbol{y}) \quad \text{Partial entropy of constituent } i$ $\left(\frac{\partial \mu_i}{\partial p}\right)_{T,\boldsymbol{n}} = \left(\frac{\partial^2 G}{\partial p \partial n_i}\right)_{T,\boldsymbol{n}'} = \left(\frac{\partial V}{\partial n_i}\right)_{T,p,\boldsymbol{n}'} = v_i(T,p,n\boldsymbol{y}) \quad \text{Partial volume of constituent } i$ $\left(\frac{\partial \mu_i}{\partial n_i}\right)_{T,p,\boldsymbol{n}'_i} = \left(\frac{\partial^2 G}{\partial n_i \partial n_i}\right)_{T,p,\boldsymbol{n}'_i} = \left(\frac{\partial \mu_j}{\partial n_i}\right)_{T,p,\boldsymbol{n}'_i} = \mu_{i,j}(T,p,n\boldsymbol{y}) = \mu_{j,i}(T,p,n\boldsymbol{y})$ notice that, from $dE = T dS - p dV + \boldsymbol{\mu} \cdot d\boldsymbol{n}$ and $dH = T dS + V dp + \boldsymbol{\mu} \cdot d\boldsymbol{n}$ we also have $\left(\frac{\partial E}{\partial n_i}\right)_{T,n,\mathbf{r}'} = e_i(T, p, n\mathbf{y}) = T s_i - p v_i + \mu_i$ Partial energy of constituent *i* $\left(\frac{\partial H}{\partial n_i}\right)_{T,p,\boldsymbol{n}'_i} = h_i(T,p,n\boldsymbol{y}) = T\,s_i + \mu_i = \left(\frac{\partial(\mu_i/T)}{\partial(1/T)}\right)_{p,\boldsymbol{n}} \quad \text{Partial enthalpy of constituent } i$ All partial properties can be evaluated once we know the chemical potentials as functions of T, p, y and n, i.e., $\mu_i = \mu_i(T, p, ny)$. Of course, the dependence on n vanishes for large n and $\mu_i = \mu_i(T, p, \boldsymbol{y})$. © 2024 Gian Paolo Beretta @MIT 2.43 Advanced Thermodynamics

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Experimental measurement of SES properties of mixtures:

Mixture properties from the partial properties

 $Eu = E - TS + pV - \boldsymbol{\mu} \cdot \boldsymbol{n} = Eu(T, p, \boldsymbol{\mu})$ $dEu = -S dT + V dp - \mathbf{n} \cdot d\mathbf{\mu} \qquad \Rightarrow \quad \left(\frac{\partial Eu}{\partial T}\right)_{T, \mu} = -S \quad \left(\frac{\partial Eu}{\partial n}\right)_{T, \mu} = V \quad \left(\frac{\partial Eu}{\partial \mu}\right)_{T, \mu, \mu'} = -n_i$ Substituting $\boldsymbol{\mu} = \boldsymbol{\mu}(T, p, \boldsymbol{n})$ we may view Eu as a function of T, p, \boldsymbol{n} $Eu = Eu(T, p, \boldsymbol{\mu}(T, p, \boldsymbol{n})) = Eu(T, p, \boldsymbol{n})$ using the chain rule and the definitions of partial entropies and volumes, its partial derivatives are $\left(\frac{\partial Eu}{\partial T}\right)_{p,\boldsymbol{n}} = \left(\frac{\partial Eu}{\partial T}\right)_{p,\boldsymbol{\mu}} + \sum_{i=1}^{\prime} \left(\frac{\partial Eu}{\partial \mu_i}\right)_{T,p,\boldsymbol{\mu}'_i} \left(\frac{\partial \mu_i}{\partial T}\right)_{p,\boldsymbol{n}} = -S + \sum_{i=1}^{\prime} n_i s_i$ $\left(\frac{\partial Eu}{\partial p}\right)_{T,\boldsymbol{n}} = \left(\frac{\partial Eu}{\partial p}\right)_{T,\boldsymbol{\mu}} + \sum_{i=1}^{\prime} \left(\frac{\partial Eu}{\partial \mu_i}\right)_{T,p,\boldsymbol{\mu}_i} \left(\frac{\partial \mu_i}{\partial p}\right)_{T,\boldsymbol{n}} = V - \sum_{i=1}^{\prime} n_i v_i$ $\left(\frac{\partial Eu}{\partial n_i}\right)_{T,n,n'} = \sum_{i=1}^{r} \left(\frac{\partial Eu}{\partial \mu_i}\right)_{T,n,n'} \left(\frac{\partial \mu_i}{\partial n_i}\right)_{T,n'} = -\sum_{i=1}^{r} n_i \mu_{i,j}$ Therefore, in general (i.e., for small and large systems) we have $S = \sum_{i=1}^{r} n_i s_i - \left(\frac{\partial Eu}{\partial T}\right)_{p,n} \qquad \stackrel{n \text{ large}}{\longrightarrow} \qquad S = \sum_{i=1}^{r} n_i s_i$ $V = \sum_{i=1}^{r} n_i v_i + \left(\frac{\partial Eu}{\partial p}\right)_{T, \mathbf{n}} \qquad \stackrel{n \text{ large}}{\longrightarrow} \qquad V = \sum_{i=1}^{r} n_i v_i$ $\sum_{i=1}^{r} n_i \,\mu_{i,j} = -\left(\frac{\partial E u}{\partial n_i}\right)_{T,p,\boldsymbol{n}'_i} \qquad \stackrel{n \text{ large}}{\longrightarrow} \qquad \sum_{i=1}^{r} n_i \,\mu_{i,j} = 0 \quad \text{Duhem-Margules relation}$

Experimental measurement of SES properties of mixtures: Mixture properties from the partial properties

So, in general (i.e., for small and large systems) we proved

$$S = \sum_{i=1}^{r} n_{i} s_{i} - \left(\frac{\partial Eu}{\partial T}\right)_{p,n} \xrightarrow{n \text{ large}} S = \sum_{i=1}^{r} n_{i} s_{i}$$

$$V = \sum_{i=1}^{r} n_{i} v_{i} + \left(\frac{\partial Eu}{\partial p}\right)_{T,n} \xrightarrow{n \text{ large}} V = \sum_{i=1}^{r} n_{i} v_{i}$$

$$\sum_{i=1}^{r} n_{i} \mu_{i,j} = -\left(\frac{\partial Eu}{\partial n_{j}}\right)_{T,p,n'_{j}} \xrightarrow{n \text{ large}} \sum_{i=1}^{r} n_{i} \mu_{i,j} = 0$$
Duhem-Margules relation

Moreover, recalling the definitions of Eu and partial properties

$$Eu = E - TS + pV - \boldsymbol{\mu} \cdot \boldsymbol{n} = F + pV - \boldsymbol{\mu} \cdot \boldsymbol{n} = G - \boldsymbol{\mu} \cdot \boldsymbol{n} = H - TS - \boldsymbol{\mu} \cdot \boldsymbol{n}$$
$$\mu_i = e_i - Ts_i + pv_i = f_i + pv_i = g_i = h_i - Ts_i$$

we also find the following relations, which justify calling partial properties that way:

$$E = \sum_{i=1}^{r} n_{i} e_{i} + Eu - T\left(\frac{\partial Eu}{\partial T}\right)_{p,n} - p\left(\frac{\partial Eu}{\partial p}\right)_{T,n} \xrightarrow{n \text{ large}} E = \sum_{i=1}^{r} n_{i} e_{i}$$

$$F = \sum_{i=1}^{r} n_{i} f_{i} + Eu - p\left(\frac{\partial Eu}{\partial p}\right)_{T,n} \xrightarrow{n \text{ large}} F = \sum_{i=1}^{r} n_{i} f_{i}$$

$$G = \sum_{i=1}^{r} n_{i} g_{i} + Eu = \sum_{i=1}^{r} n_{i} \mu_{i} + Eu \xrightarrow{n \text{ large}} G = \sum_{i=1}^{r} n_{i} g_{i} = \sum_{i=1}^{r} n_{i} \mu_{i}$$

$$H = \sum_{i=1}^{r} n_{i} h_{i} + Eu - T\left(\frac{\partial Eu}{\partial T}\right)_{p,n} \xrightarrow{n \text{ large}} H = \sum_{i=1}^{r} n_{i} h_{i}$$

$$(2024 \text{ Gian Paolo Beretta} \quad (MIT 2.43 \text{ Advanced Thermodynamics})$$

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