2.43 ADVANCED THERMODYNAMICS

Spring Term 2024 LECTURE 18

Room 3-442 Friday, April 12, 11:00am - 1:00pm

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Systems with chemical reactions

energy and entropy balances notation and stoichiometry properties of reaction properties of formation

Energy and entropy balances for closed and open systems with chemical reactions

$$
\begin{bmatrix}\nV, n_{1R}, \ldots, n_{iR}, \ldots, n_{rR} \\
T_{R}, p_{R}, \mu_{R} \\
U_{R} = \sum_{i} n_{iR} u_{i}(T_{R}, p_{R}, \mathbf{y}_{R}) \\
S_{R} = \sum_{i} n_{iR} s_{i}(T_{R}, p_{R}, \mathbf{y}_{R}) \\
\text{initial state "R"} \\
U_{P} = n_{iR} + \nu_{i} \epsilon\n\end{bmatrix}\n\begin{bmatrix}\nV, n_{1P}, \ldots, n_{iP}, \ldots, n_{rP} \\
T_{P}, p_{P}, \mu_{P} \\
S_{P} = \sum_{i} n_{iP} u_{i}(T_{P}, p_{P}, \mathbf{y}_{P}) \\
S_{P} = \sum_{i} n_{iP} s_{i}(T_{P}, p_{P}, \mathbf{y}_{P}) \\
\text{initial state "R"} \\
U_{P} - U_{R} = \sum_{k} Q_{k}^{+} - W^{+} \\
S_{P} - S_{R} = \sum_{k} \frac{Q_{k}^{+}}{T_{k}} + S_{\text{irr}} \\
\text{final state "P"} \\
T_{R}, p_{R}, \mu_{R} \\
T_{R}, p_{R}, \mu_{R}\n\end{bmatrix}\n\begin{bmatrix}\n\text{steady state reactor} \\
S_{P} - S_{R} = \sum_{k} \frac{Q_{k}^{+}}{T_{k}} + S_{\text{irr}} \\
\text{final state "P"} \\
T_{P}, p_{P}, \mathbf{y}_{P}\n\end{bmatrix}
$$
\n
$$
\begin{bmatrix}\n\dot{n}_{1R}, \ldots, \dot{n}_{iR}, \ldots, \dot{n}_{iR} \\
\hline\nT_{R}, p_{R}, \mu_{R} \\
\hline\n\dot{R} = \sum_{i} \dot{n}_{iR} h_{i}(T_{R}, p_{R}, \mathbf{y}_{R}) \\
\hline\n\dot{S}_{R} = \sum_{i} \dot{n}_{iR} h_{i}(T_{R}, p_{R}, \mathbf{y}_{R}) \\
\hline\n\dot{S}_{R} = \sum_{i} \dot{n}_{iR} s_{i}(T_{R}, p_{R}, \mathbf{y}_{R})\n\end{bmatrix}\n\begin{bmatrix}\n\dot{n}_{1P} & \dot{n}_{1P}, \ldots, n_{iP}, \ldots, n_{iP} \\
\hline\n\dot{S}_{P} = \sum_{i} n_{iP} u_{i}(T_{P},
$$

Notation and stoichiometry: proportionality relations, reaction coordinates

chemical reaction mechanism $2H_2 + O_2 = 2H_2O$ or $-2H_2 - O_2 + 2H_2O = 0$ composition: n_{H_2} , n_{O_2} , $n_{\text{H}_2\text{O}}$ changes in composition: Δn_{H_2} , Δn_{O_2} , $\Delta n_{\text{H}_2\text{O}}$ proportionality relations $\frac{\Delta n_{\text{H}_2}}{2} = \frac{\Delta n_{\text{O}_2}}{1} = \frac{\Delta n_{\text{H}_2\text{O}}}{2} = \Delta \epsilon$ $\epsilon =$ reaction coordinate

chemical symbols $A_1 = H_2, A_2 = O_2, A_3 = H_2O$ stoichiometric coefficients $\nu_1 = -2, \nu_2 = -1, \nu_3 = 2$ rewrite the mechanism as $\nu_1A_1 + \nu_2A_2 + \nu_3A_3 = 0$ $\sum_{i} \nu_i A_i = 0$ proportionality relations $\frac{\Delta n_1}{\nu_1} = \frac{\Delta n_2}{\nu_2} = \frac{\Delta n_3}{\nu_3} = \Delta \epsilon$ $\Delta n_i = \nu_i \, \Delta \epsilon \quad \forall i$ $n_{i\text{P}} = n_{i\text{R}} + \nu_i \epsilon$ where $\epsilon = \epsilon_{\rm P} - \epsilon_{\rm R}$ $\dot{n}_i = \nu_i \dot{\epsilon} \quad \forall i$ $\dot{n}_{i\text{p}} = \dot{n}_{i\text{B}} + \nu_i \dot{\epsilon}$

Properties of reaction

The balance equations require these differences: $U_P - U_R$, $S_P - S_R$, $\dot{H}_P - \dot{H}_R$, $\dot{S}_P - \dot{S}_R$. For example, $U_{\rm P}-U_{\rm R}=\sum_i n_{iP} u_i(T_{\rm P}, p_{\rm P}, y_{iP})-\sum_i n_{iR} u_i(T_{\rm R}, p_{\rm R}, y_{iR})$ = add and subtract $u_{ii}(T, p_o)$ $=\sum_{i}n_{i\text{P}}\left[u_{i}(T_{\text{P}},p_{\text{P}},y_{i\text{P}})-u_{ii}(T,p_{o})\right]-\sum_{i}n_{i\text{R}}\left[u_{i}(T_{\text{R}},p_{\text{R}},y_{i\text{R}})-u_{ii}(T,p_{o})\right]+\sum_{i}\left(n_{i\text{P}}-n_{i\text{R}}\right)u_{ii}(T,p_{o})$ $\Delta u_i^{\text{mix}}|_{T_P, p_P} + \Delta u_i^{\text{mix}}|_{T_R, p_R} + u_{ii}(T_P, p_P) - u_{ii}(T, p_o)$ $+ u_{ii}(T_R, p_R) - u_{ii}(T, p_o)$ $\epsilon \sum_i \nu_i u_{ii}(T, p_o) = \epsilon \Delta u^o(T)$ $H_P - H_R = \sum_i \dot{n}_{iP} h_i(T_P, p_P, y_{iP}) - \sum_i \dot{n}_{iR} h_i(T_R, p_R, y_{iR}) =$ add and subtract $h_{ii}(T, p_o)$ $= \sum_{i} \dot{n}_{iP} \left[h_i(T_P, p_P, y_{iP}) - h_{ii}(T, p_o) \right] - \sum_{i} \dot{n}_{iR} \left[h_i(T_R, p_R, y_{iR}) - h_{ii}(T, p_o) \right] + \sum_{i} (\dot{n}_{iP} - \dot{n}_{iR}) h_{ii}(T, p_o)$ $\Delta h_i^{\text{mix}}|_{T_P, p_P} + \Delta h_i^{\text{mix}}|_{T_R, p_R} + h_{ii}(T_P, p_P) - h_{ii}(T, p_o)$ $\qquad + h_{ii}(T_R, p_R) - h_{ii}(T, p_o)$ $\qquad \qquad + h_{ii}(T_R, p_R) - h_{ii}(T, p_o)$

So, we defined the properties of reaction

at *T* and standard pressure
\n
$$
p_o = 1 \text{ atm:}
$$
\n
$$
\Delta u^o(T) = \sum_i \nu_i u_{ii}(T, p_o)
$$
\n
$$
\Delta s^o(T) = \sum_i \nu_i s_{ii}(T, p_o)
$$
\n
$$
\Delta h^o(T) = \sum_i \nu_i h_{ii}(T, p_o)
$$
\n
$$
\Delta g^o(T) = \sum_i \nu_i g_{ii}(T, p_o)
$$
\n
$$
\Delta g^o(T) = \Delta h^o(T) - T \Delta s^o(T)
$$
\n
$$
\Delta v^o(T) = \sum_i \nu_i v_{ii}(T, p_o)
$$
\n
$$
\Delta h^o(T) = \Delta u^o(T) + p_o \Delta v^o(T)
$$

at STP:
$$
T_o = 298.15 \text{ K}
$$

\nand $p_o = 1 \text{ atm}$:
\n
$$
\Delta u^o = \sum_i \nu_i u_{ii}(T_o, p_o)
$$
\n
$$
\Delta s^o = \sum_i \nu_i s_{ii}(T_o, p_o)
$$
\n
$$
\Delta h^o = \sum_i \nu_i h_{ii}(T_o, p_o)
$$
\n
$$
\Delta g^o = \sum_i \nu_i g_{ii}(T_o, p_o)
$$
\n
$$
\Delta g^o = \Delta h^o - T_o \Delta s^o
$$
\n
$$
\Delta v^o = \sum_i \nu_i v_{ii}(T_o, p_o)
$$
\n
$$
\Delta h^o = \Delta u^o + n \Delta v^o
$$

using the Hess relation and the properties of formation at STP:

$$
\Delta u^o = \sum_i \nu_i \, \Delta u_f^o|_i
$$

\n
$$
\Delta s^o = \sum_i \nu_i \, \Delta s_f^o|_i
$$

\n
$$
\Delta h^o = \sum_i \nu_i \, \Delta h_f^o|_i
$$

\n
$$
\Delta g^o = \sum_i \nu_i \, \Delta g_f^o|_i
$$

\n
$$
\Delta v^o = \sum_i \nu_i \, \Delta v_f^o|_i
$$

Enthalpy of formation of H **and** H-H **bond energy**

van der Waals forces vs covalent bonds

van der Waals forces have a significant impact on the physical properties of substances, such as boiling points, melting points, and solubility. But they are very weak compared to ionic or covalent bonds.

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Enthalpy of formation and bond energies

 $2 * \epsilon(H-H) + \epsilon(O=O) - 4 * \epsilon(H-O)$ -478

Enthalpy of formation of 2 moles of water -484

$$
V_{\text{LJ}}(r) = 4\epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^{6} \right]
$$

$$
= \epsilon \left[\left(\frac{r_{\min}}{r} \right)^{12} - 2\left(\frac{r_{\min}}{r} \right)^{6} \right]
$$

 $V(r)$ in kJ/mol vs r in pm (10⁻¹² m)

Properties of formation and Hess relation

Define a complete set of independent elemental constituents A_i^e .

Completeness: for every molecule A_i there must be a reaction of formation from the A_j^e 's:

 $\sum_{i} \alpha_{ij} A_i^e = A_i$ or in our notation $-\sum_{i} \alpha_{ij} A_i^e + A_i = 0$ or also $A_i = \sum_{i} \alpha_{ij} A_i^e$ **Independence**: the A_j^e 's cannot be transformed into one another:

$$
\textstyle{\sum_j}\beta_jA_j^\text{e}=0\quad\Rightarrow\quad\beta_j=0\ \forall j
$$

Practical choice: $\{A_i^e\}$ = the species with nuclei of only one type that are most stable at STP, one for each atomic number Z_j^e in the periodic table (isotopes are chemically equivalent). **Property of formation at STP** of a generic chemical species A_i :

 $\Delta u_{\rm f}^o|_i = -\sum_{j} \alpha_{ij} u_{jj}^{\rm e}(T_o, p_o) + u_{ii}(T_o, p_o)$

We can safely select $u_{ij}^e(T_o, p_o) = 0$: in fact, for any mechanism $\sum_i \nu_i A_i = 0$ we can write:

$$
\sum_{i} \nu_i A_i = \sum_{i} \nu_i \sum_{j} \alpha_{ij} A_j^e = \sum_{j} (\sum_{i} \nu_i \alpha_{ij}) A_j^e = 0 \Rightarrow \sum_{i} \nu_i \alpha_{ij} = 0 \forall j
$$

$$
\Delta u^o = \sum_{i} \nu_i u_{ii} (T_o, p_o) = \sum_{i} \nu_i [\Delta u_i^o|_i + \sum_{j} \alpha_{ij} u_{jj}^e (T_o, p_o)] = \sum_{i} \nu_i \Delta u_i^o|_i + \sum_{j} \sum_{i} \nu_i \alpha_{ij} u_{jj}^e (T_o, p_o)
$$

Therefore, we conventionally choose these references: and use the **Hess relation** $u_{jj}^{e}(T_o, p_o) = 0$ for the A_j^{e} 's and $\Delta u^o = \sum_i \nu_i \, \Delta u_i^o|_i$ $u_{ii}(T_o, p_o) = \Delta u_{\rm f}^o|_i$ for all other species

Careful! Properties of formation at STP are listed for the **fictitious ideal gas state** by extrapolating the ideal-gas correlations for $c_{p,ii}(T)$ from the T, p range where they hold, down to T_o , p_o : $h_{ii}(T_o, p_o) = h_{ii}(T, p_o) - \int_{T_o}^{T} c_{p,ii}(T') dT' \qquad s_{ii}(T_o, p_o) = s_{ii}(T, p_o) - \int_{T_o}^{T} \frac{c_{p,ii}(T')}{T'} dT' - R \ln \frac{p_o}{p}$

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Slide 18.09

Properties of formation and Hess relation

Hess relation $\Delta h^o = \sum_i \nu_i \Delta h_f^o|_i$ For the reaction $\text{CH}_4 + 2\text{O}_2 = \text{CO}_2 + 2\text{H}_2\text{O}$ $M_{\text{CH}_4} = 16 \text{ kg/kmol}$ $-1(-74.9) - 2(0) + 1(-393.8) + 2(-242) = -802.9 \,\text{MJ/kmol}_{\text{CH}_4} = -50.2 \,\text{MJ/kg}_{\text{CH}_4}$ For the reaction $C_8H_{18} + 12.5O_2 = 8 CO_2 + 9 H_2O$ $M_{C_8H_{18}} = 114 \text{ kg/kmol}$ $-1(-224.3) - 12.5(0) + 8(-393.8) + 9(-242) = -5120 \,\text{MJ/kmol}_{\text{CaH}_{18}} = -44.9 \,\text{MJ/kg}_{\text{CaH}_{18}}$

Adapted from Table 29.4 of Gyftopoulos, Beretta, Thermodynamics. Foundations and applications, Dover, 2005

Systems with chemical reactions

what determines the direction of spontanueos reaction and how temperature affects it

Exergy Maximum work obtainable Minimum work required

When is an isothermobaric reaction product- or reactant-favored?

$$
\begin{array}{c}\nT, p_o \\
\hline\n\hat{H}_{\rm R}, \dot{S}_{\rm R}\n\end{array}\n\qquad\n\begin{array}{c}\n\text{steady state reactor} \\
\sum_{i=1}^{r} \nu_i A_i = 0 \\
\dot{\epsilon} \quad \dot{S}_{\rm irr} \\
\dot{H}_{\rm P}, \dot{S}_{\rm P}\n\end{array}\n\qquad\n\begin{array}{c}\n\hat{H}_{\rm P}, \dot{S}_{\rm P} \\
\dot{\theta}^+,\hat{H}_{\rm P}, \dot{S}_{\rm P}\n\end{array}
$$
\n
$$
0 = \dot{H}_{\rm R}(T) - \dot{H}_{\rm P}(T) + \dot{Q}^+
$$
\n
$$
0 = \dot{S}_{\rm R}(T) - \dot{S}_{\rm P}(T) + \frac{\dot{Q}^+}{T_Q} + \dot{S}_{\rm irr}
$$
\n
$$
T_Q \dot{S}_{\rm irr} = -\dot{\epsilon} \underbrace{[\Delta h^o(T) - T_Q \Delta s^o(T)]}_{\text{this must be negative}} \geq 0
$$
\n
$$
\text{for the reaction to be}
$$
\n
$$
\text{product-favored}
$$

Assume:

- \bullet steady-state
- heat interaction at T_Q maintains inlet and outlet temperatures at T \bullet in and out streams are either pure component or ideal GD mixtures

• neglect mixing terms in the entropies

$$
\dot{H}_P(T) - \dot{H}_R(T) = \dot{\epsilon} \Delta h^o(T)
$$

$$
\dot{S}_P(T) - \dot{S}_R(T) = \dot{\epsilon} \Delta s^o(T)
$$

$$
\frac{\dot{S}_{irr}}{R \dot{\epsilon}} = -\frac{\Delta h^o(T) - T_Q \Delta s^o(T)}{R T_Q}
$$

When is an isothermobaric reaction product- or reactant-favored?

T, p_o	steady state reactor	T, p_o
$\overrightarrow{H_R}, \dot{S_R}$	$\overrightarrow{\epsilon} \dot{S_{irr}}$	$\dot{H_P}, \dot{S_P}$
$\dot{n}_i \mathbf{p} = \dot{n}_{iR} + \nu_i \dot{\epsilon}$	$\dot{Q} \leftarrow \mathbf{F_Q} = T$	
$0 = \dot{H_R}(T) - \dot{H_P}(T) + \dot{Q} \leftarrow$		
$0 = \dot{S_R}(T) - \dot{S_P}(T) + \frac{\dot{Q} \leftarrow \dot{S_{irr}}$		
$T_Q \dot{S_{irr}} = -\dot{\epsilon} [\Delta h^o(T) - T \Delta s^o(T)] = -\dot{\epsilon} \Delta g^o(T) \ge 0$		
$\Delta g^o(T)$ must be negative for the reaction to be product-favored		

Assume:

- \bullet steady-state
- heat interaction at T_Q maintains inlet and outlet temperatures at T \bullet in and out streams are either pure component or ideal GD mixtures • neglect mixing terms in the entropies

$$
\dot{H}_P(T) - \dot{H}_R(T) = \dot{\epsilon} \Delta h^o(T)
$$

$$
\dot{S}_P(T) - \dot{S}_R(T) = \dot{\epsilon} \Delta s^o(T)
$$

$$
\frac{\dot{S}_{irr}}{R \dot{\epsilon}} = -\frac{\Delta g^o(T)}{RT}
$$

$$
\frac{d}{dT} \frac{\dot{S}_{irr}}{R \dot{\epsilon}} = -\frac{d}{dT} \frac{\Delta g^o(T)}{RT} = \frac{\Delta h^o(T)}{RT^2}
$$

When is an isothermal reaction product- or reactant-favored?

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When is an adiabatic reaction product- or reactant-favored?

$$
\overrightarrow{T}, p_o
$$
 steady state reactor

$$
\Sigma_{i=1}^r \nu_i A_i = 0
$$

$$
\overrightarrow{H}_{\text{R}}, \overrightarrow{S}_{\text{R}}
$$

$$
\overrightarrow{i}_{i\text{P}} = \overrightarrow{n}_{i\text{R}} + \nu_i \overrightarrow{\epsilon}
$$

$$
\overrightarrow{H}_{\text{P}}, \overrightarrow{S}_{\text{P}}
$$

$$
0 = \dot{H}_R(T) - \dot{H}_P(T_{\text{ad}}) \Rightarrow \text{ solve for } T_{\text{ad}}
$$

$$
0 = \dot{S}_R(T) - \dot{S}_P(T_{\text{ad}}) + \dot{S}_{\text{irr}}
$$

$$
T \dot{S}_{\text{irr}} = -\dot{\epsilon} \Delta g^o(T) - \dot{E}_{\text{X}P}(T_{\text{ad}}) \ge 0
$$

For $\dot{\epsilon} > 0$, $\Delta g^o(T)$ must be negative enough that $-\dot{\epsilon} \Delta g^o(T)$ is greater than the exergy^{*} needed to heat/cool the products from environmental T to T_{ad}

Assume:

- \bullet steady-state
- \bullet adiabatic, no heat interactions
- in and out streams are either pure component or ideal GD mixtures
- neglect mixing terms in the entropies

$$
\dot{H}_P(T) - \dot{H}_R(T) = \dot{\epsilon} \Delta h^o(T)
$$

$$
\dot{S}_P(T) - \dot{S}_R(T) = \dot{\epsilon} \Delta s^o(T)
$$

$$
\dot{E}_{\text{XP}}(T_{\text{ad}}) = \dot{H}_P(T_{\text{ad}}) - \dot{H}_P(T) - T \left[\dot{S}_P(T_{\text{ad}}) - \dot{S}_P(T) \right]
$$

Exergy / Maximum work obtainable / Minimum work required

Assume:

• steady-state, with a work interaction • heat interaction at T_Q maintains

• in and out streams are either pure

• neglect mixing terms in the entropies

in let and outlet temperatures at T

component or ideal GD mixtures

$$
\begin{array}{c}\nT, p_o \\
\hline\n\hat{H}_{\text{R}}, \dot{S}_{\text{R}} \\
\hline\n\hat{n}_i \mathbf{p} = \dot{n}_{i\text{R}} + \nu_i \dot{\epsilon} \\
\hline\n\hat{W}^+ \end{array}\n\quad\n\begin{array}{c}\n\text{steady state reactor} \\
\hline\n\hat{L}_{i=1}^r \nu_i A_i = 0 \\
\hline\n\hat{H}_{\text{P}}, \dot{S}_{\text{P}} \\
\hline\n\hat{H}_{\text{P}}, \dot{S}_{\text{P}} \\
\hline\n\hat{W}^+ \end{array}
$$

 $0 = H_R(T) - H_P(T) + Q^{\leftarrow} - W^{\rightarrow}$ $\dot{H}_P(T) - \dot{H}_R(T) = \dot{\epsilon} \Delta h^o(T)$ $0 = \dot{S}_R(T) - \dot{S}_P(T) + \frac{\dot{Q}^{\leftarrow}}{T_O} + \dot{S}_{irr}$ $\dot{S}_P(T) - \dot{S}_R(T) = \dot{\epsilon} \Delta s^o(T)$ $\dot{W}^{\rightarrow} = -\dot{\epsilon} \left[\Delta h^o(T) - T_Q \Delta s^o(T) \right] - T_Q \dot{S}_{irr}$ this must be negative $\dot{W}_{\text{max}}^{\rightarrow} = -\dot{\epsilon} [\Delta h^o(T) - T_Q \Delta s^o(T)]$ for $\dot{\epsilon} > 0$ and $\dot{W} \dot{\tau} > 0$ Equivalently, $\dot{W}^{\leftarrow} = \dot{\epsilon} \left[\Delta h^o(T) - T_Q \, \Delta s^o(T) \right] + T_Q \dot{S}_{irr}$ if this is positive $\dot{W}_{\min}^{\leftarrow} = \dot{\epsilon} [\Delta h^o(T) - T_O \Delta s^o(T)]$ $\dot{\epsilon} > 0$ requires $\dot{W}^{\leftarrow} > 0$

Systems with chemical reactions

Oxidation of hydrocarbons Heating values of fuels Exergy of a fuel Adiabatic flame temperature

Oxidation of hydrocarbons

Overall basic reaction mechanism, assuming full oxidation

$$
C_kH_{\ell} + \left(k + \frac{\ell}{4}\right)O_2 = k\,CO_2 + \frac{\ell}{2}H_2O
$$

If oxygen is taken from dry air with mole fraction $y_{O_2} = 0.2095$, 1 mol of O_2 requires 4.77 mol of air. From this, the **stoichiometric air-fuel ratio** is defined **as**

$$
\left(\frac{n_{\text{air}}}{n_{\text{fuel}}}\right)_{\text{s}} = 4.77 \left(k + \frac{\ell}{4}\right) \qquad \left(\frac{m_{\text{air}}}{m_{\text{fuel}}}\right)_{\text{s}} = 4.77 \left(k + \frac{\ell}{4}\right) \frac{28.96}{M_{\text{fuel}}}
$$

Moreover, in order to increase the probability of having the oxidation of all the molecules, an excess of air is usually used i.e. $n_{air} = (n_{air})_s + (n_{air})_{excess}$, so that the **equivalence ratio** is defined

$$
\lambda = \frac{n_{\text{air}}/n_{\text{fuel}}}{(n_{\text{air}}/n_{\text{fuel}})_{\text{s}}} = \frac{m_{\text{air}}/m_{\text{fuel}}}{(m_{\text{air}}/m_{\text{fuel}})_{\text{s}}} > 1
$$

For most hydrocarbons (except H_2)

$$
-2.4\% < \frac{\Delta h^0 - \Delta g^0}{\Delta g^0} < +2.4\%
$$

and

$$
40\,\mathrm{MJ/kg} < \mathrm{LHV} = -\Delta h^o < 50\,\mathrm{MJ/kg}
$$

Properties of the reaction of oxidation of hydrocarbons

Adapted from Figure 31.7 of Gyftopoulos, Beretta, Thermodynamics. Foundations and applications, Dover, 2005

Heating values of fuels

Source: Data mainly from J. B. Heywood, Internal Combustion Engine Fundamentals, McGraw-Hill, New York, 1988.

^a Values of h_{fg} at 25°C for liquid fuels. Values of ρ and $(c_p)_f$ at 25°C and 1 atm for liquid fuels. Values of $(c_p)_g$ at 25°C and saturation pressure for liquid fuels, and at 1 atm and 25°C for gaseous fuels.

Adapted from Figure 31.1 of Gyftopoulos, Beretta, Thermodynamics. Foundations and applications, Dover, 2005

Lower and higher heating values of fuels

In combustion, water is present both in reactants (moist air) and in products, this complicates the stoichiometry. The highest possible mole fraction of water vapor in the gaseous mixture of products of combustion is given by Raoult's law $(1 = H₂O)$

 $y_{1g}^{\text{max}} p = p_{\text{sat},11}^{\text{fg}}(T)$ where $y_{1f} = 1$ (condensate is pure liquid water)

Let $\phi = y_{1g}/y_{1g}^{\text{max}}$ be the relative humidity of the products of combustion in the gaseous phase. Since $p_{\text{sat,11}}^{\text{fg}}(25\text{°C}) = 3.169 \text{ kPa}$, at STP $(T_o = 25 \text{°C} \text{ and } p_o = 1 \text{ atm})$, the maximum amount of vapor is $y_{1g}^{\text{max}} = 3.169/101.325 = 0.0316$, so that by splitting water vapor and condensate, i.e., $n_{\rm H_2O}=n_{\rm H_2O,g}+n_{\rm H_2O,f}$ we get

$$
\frac{n_{\text{H}_2\text{O}}}{n_{\text{C}_k\text{H}_\ell}} = \frac{n_{\text{H}_2\text{O},\text{f}}}{n_{\text{C}_k\text{H}_\ell}} + \frac{n_{\text{H}_2\text{O},\text{g}}}{n_{\text{C}_k\text{H}_\ell}} = \frac{\ell}{2}
$$
\n
$$
\frac{n_{\text{H}_2\text{O},\text{g}}}{n_{\text{C}_k\text{H}_\ell}} \le \frac{0.0316\left[4.77\left(k + \frac{\ell}{4}\right) - \frac{\ell}{4}\right]}{1 - 0.0316} = 0.123\,\ell + 0.156\,k
$$
\n
$$
\frac{n_{\text{tot,prod}}}{n_{\text{C}_k\text{H}_\ell}} = \frac{\ell}{4} + 4.77\,\left(k + \frac{\ell}{4}\right)
$$
\n
$$
\Rightarrow \quad \text{so that when } \ell/2 \text{ exceeds this amount, i.e., when } \ell/k > 0.414, \text{ then part of the water condenses}
$$
\n
$$
\frac{n_{\text{H}_2\text{O},g}}{n_{\text{tot,prod}} - n_{\text{H}_2\text{O},\text{f}}} = y_{1\text{g}} \le 0.0316
$$
\n
$$
\frac{n_{\text{H}_2\text{O},g}}{n_{\text{C}_k\text{H}_\ell}} = \frac{\frac{\ell}{2} - 0.0316\left[\frac{\ell}{2} + 4.77\left(k + \frac{\ell}{4}\right)\right]}{1 - 0.0316} = 0.469\,\ell - 0.156\,k
$$

meaning that the $\dot{Q}_{o}^{\rightarrow}$ extracted is not only given by $\dot{n}_{C_kH_\ell}\Delta h^o$ but also by the condensation enthalpy

$$
\dot{Q}_{o}^{\rightarrow} = \dot{H}_{R}^{o} - \dot{H}_{P}^{o} = \dot{n}_{C_{k}H_{\ell}}HV_{\text{actual}} \quad \text{where} \quad \text{HV}_{\text{actual}} = -\Delta h^{o} + \frac{\dot{n}_{H_{2}O,f}}{\dot{n}_{C_{k}H_{\ell}}}h_{H_{2}O,fg}(T_{o})
$$

With less details, in the common practice one defines the lower and higher heating values

$$
LHV = -\Delta h^o \qquad \text{and} \qquad HHV = -\Delta h^o + \frac{\ell}{2} h_{H_2O,fg}(T_o)
$$

For example, for methane, LHV, HV_{actual} , and HHV are, respectively, 50.1, 54.7, and 55.5 MJ/kg.

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Heating value and Exergy of a fuel / Adiabatic flame temperature

$$
\begin{array}{c|c} T_o, p_o & \text{steady state reactor} & T_P, p_o \\ \hline \dot{H}_{\rm R}, \dot{S}_{\rm R} & \dot{\epsilon} = \dot{n}_{\rm C_kH_\ell} & \dot{S}_{\rm irr} & \dot{H}_{\rm P}, \dot{S}_{\rm P} \\ \dot{n}_{\rm i\rm P} = \dot{n}_{\rm i\rm R} + \nu_i \dot{\epsilon} & \ddot{\xi} = \dot{\eta}_{\rm C_kH_\ell} & \dot{\zeta}_{\rm i\rm F}, \dot{S}_{\rm P} \\ \hline \dot{W}^\rightarrow & \dot{\psi}^\rightarrow & \dot{Q}_o^\leftarrow, T_o \end{array}
$$

Heating value setup: $T_{\rm P} = T_o$, $\dot{W}^{\rightarrow} = 0$, $\dot{\epsilon} = \dot{n}_{\rm C_kH_{\ell}}$

 $\dot{Q}_{o}^{\rightarrow} = \dot{H}_{R}^{o} - \dot{H}_{P}^{o} = \dot{\epsilon} \text{HV}_{\text{actual}}$ where $\text{HV}_{\text{actual}} = -\Delta h^{o} + \frac{\dot{n}_{\text{H}_2\text{O,f}}}{\dot{n}_{\text{C,H}_2}} h_{\text{H}_2\text{O,fg}}(T_o)$

Exergy setup: $T_P = T_o$, $\dot{W}^{\rightarrow} = \dot{W}_{\text{max}}^{\rightarrow}$, $\dot{S}_{irr} = 0$ $\dot{W}_{\text{max}}^{\rightarrow} = \dot{H}_R^o - \dot{H}_P^o - T_o \left(\dot{S}_R^o - \dot{S}_P^o \right) = \dot{\epsilon} E_{X_{\text{C}_k\text{H}_\ell}}$ where $E_{X_{\text{C}_k\text{H}_\ell}} = -\Delta h^o + T_o \Delta s^o = -\Delta g^o$ Notice that condensation does not contribute because $h_{\text{H}_2\text{O},\text{fg}}(T_o) - T_o s_{\text{H}_2\text{O},\text{fg}}(T_o) = 0.$ **Adiabatic flame temperature** setup: $T_P = T_{ad}$, $\dot{Q}_{o}^{\rightarrow} = 0$, $\dot{W}^{\rightarrow} = 0$, assume ideal GD mix of

ideal gases

$$
0 = \dot{H}_R^o - \dot{H}_P(T_{ad}) = \dot{H}_R^o - \dot{H}_P^o + \dot{H}_P^o - \dot{H}_P(T_{ad}) = \dot{\epsilon} \text{HV}_{\text{actual}} - \dot{n}_P \int_{T_o}^{T_{ad}} \sum_i y_{iP} \, c_{p,ii}(T') \, dT' \Rightarrow \text{solve}_{\text{for } T_{ad}}
$$

$$
\dot{S}_{irr} = \dot{S}_P(T_{ad}) - \dot{S}_R^o = \dot{S}_P(T_{ad}) - \dot{S}_P^o + \dot{S}_P^o - \dot{S}_R^o = \dot{n}_P \int_{T_o}^{2} \frac{\sum_i y_{iP} c_{p,ii}(T')}{T'} dT' + \dot{\epsilon} \Delta s^o \approx -0.3 \dot{\epsilon} \frac{\Delta g^o}{T_o}
$$

Typical values of
$$
\frac{T_o S_{irr}}{-\dot{\epsilon} \Delta g^o}
$$
 are in the range 26% to 30%

Adiabatic flame temperature calculations

$$
c_{p,ii}(T) = a_{ii} + b_{ii}T^{1/4} + c_{ii}T^{1/2} + d_{ii}T^{3/4}
$$

\n
$$
\sum_{i} y_{iP} c_{p,ii}(T) = a + bT^{1/4} + cT^{1/2} + dT^{3/4}
$$

\nwhere $a = \sum_{i} y_{iP} a_{ii}$ $b = \sum_{i} y_{iP} b_{ii}$ $c = \sum_{i} y_{iP} c_{ii}$ $d = \sum_{i} y_{iP} d_{ii}$
\n
$$
\int_{T_o}^{T_{ad}} \sum_{i} y_{iP} c_{p,ii}(T') dT' = a (T_{ad} - T_o) + \frac{4}{5}b (T_{ad}^{5/4} - T_o^{5/4}) + \frac{2}{3}c (T_{ad}^{3/2} - T_o^{3/2}) + \frac{4}{7}d (T_{ad}^{7/4} - T_o^{7/4})
$$

\n
$$
\int_{T_o}^{T_{ad}} \frac{\sum_{i} y_{iP} c_{p,ii}(T')}{T'} dT' = a \ln \frac{T_{ad}}{T_o} + 4b (T_{ad}^{1/4} - T_o^{1/4}) + 2c (T_{ad}^{1/2} - T_o^{1/2}) + \frac{4}{3}d (T_{ad}^{3/4} - T_o^{3/4})
$$

Adapted from Table 20.2 of Gyftopoulos, Beretta, Thermodynamics. Foundations and applications, Dover, 2005

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Adiabatic flame temperature calculations

E18-AdiabaticFlameTemperature(withEquilibrium).xlsx

Adiabatic flame temperature and pressure in a closed bomb

Volume balance (assuming GD mix of ideal gases): $V_P = \frac{n_P R T_P}{p_P} = V_R = \frac{n_R R T_o}{p_o}$
 Heating value setup: $W^{\rightarrow} = 0$, $\epsilon = n_{C_k H_{\ell}}$, $n_P = n_R + \nu \epsilon$, $\frac{p_P}{p_o} = 1 + \nu \frac{\epsilon}{n_R}$ $U_P - U_R^o = -Q_o^{\rightarrow} = -\epsilon \text{HV}_{\text{actual}}^{\text{bomb}}$ where $\text{HV}_{\text{actual}}^{\text{bomb}} = -\Delta u^o + \frac{n_{\text{H}_2\text{O,f}}}{n_{\text{C-H}}} u_{\text{H}_2\text{O,fg}}(T_o)$ **Adiabatic flame temperature** setup: $T_P = T_{ad}$, $Q_o^{\rightarrow} = 0$, $W^{\rightarrow} = 0$, $\frac{p_{ad}}{p_o} = \frac{T_{ad}}{T_o} \left(1 + \nu \frac{\epsilon}{n_R}\right)$ $0 = U_P(T_{\text{ad}}) - U_R^o = U_P(T_{\text{ad}}) - U_P^o + U_P^o - U_R^o = n_P \int_T^{T_{\text{ad}}} \sum_i y_{iP} c_{v,ii}(T') dT' - \dot{\epsilon} \text{HV}_{\text{actual}}^{\text{bomb}} \Rightarrow \frac{\text{solve}}{\text{for } T_{\text{ad}}}$ $S_{irr} = S_P(T_{ad}, p_{ad}) - S_R^o = S_P(T_{ad}, p_{ad}) - S_P^o + S_P^o - S_R^o$ $= n_P \int_{T_c}^{T_{\rm ad}} \frac{\sum_i y_{iP} c_{p,ii}(T')}{T'} dT' - n_P R \ln \frac{p_{\rm ad}}{p_o} + \epsilon \Delta s^o \approx -(0.22 \div 0.26) \epsilon \frac{\Delta g^o}{T_o}$ Typical values of $\frac{T_o \dot{S}_{irr}}{-\dot{\epsilon} \Delta q^o}$ are in the range 22% to 26%

Adiabatic flame temperature

Adapted from Figures 31.8 and 31.9 of Gyftopoulos, Beretta, Thermodynamics. Foundations and applications, Dover, 2005

Summary on an energy vs entropy diagram

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Is a reversible combustion possible?

Figures 3 and 4 from G.P. Beretta, A.M. Lezzi, A. Niro, and M. Silvestri, [On the concept of a reversible flame,](https://gianpaolo-beretta.unibs.it/Beretta-papers-online/ic22-BerettaLezziNiroSilvestri-Flowers92-165-1992.pdf) in Energy for the Transition Age, FLOWERS '92, Edited by S.S. Stecco and M.J. Moran, Nova Science Pu. Inc., New York, Additional Proceedings, pp. 165-177 (1992).

The question was first raised and discussed in Keenan (1941) and Obert (1948,1973), and later also by Richter & Knoche (1983), Dunbar & Lior (1994).

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Slide 28:

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