

# 3.020 Lecture 29

Prof. Rafael Jaramillo

# 1 Maximum entropy condition and the Boltzmann distribution

- from last time
  - optimizing  $S'$  for isolated system
  - $dS' = 0$  subject to constraints  $dU' = 0, dn_{TOT} = 0$
  - method of Lagrange multipliers:  $\alpha, \beta$

$$\sum_i \underbrace{(\dots)}_{\text{set coefficients to 0}} dn_i = 0$$

- determine  $\alpha$  by normalization

$$\frac{n_i}{n_{TOT}} = \frac{e^{\beta \epsilon_i / k_B}}{Q}, \quad Q = \sum_i e^{\beta \epsilon_i / k_B}$$

- Determine  $\beta$  by considering microscopic, reversible process  $dn_i$

$$\begin{aligned} dS' &= -k_B \sum_i \ln\left(\frac{n_i}{n_{TOT}}\right) dn_i = -k_B \sum_i \ln\left(\frac{e^{\beta \epsilon_i / k_B}}{Q}\right) dn_i \\ &= -\beta \sum_i \epsilon_i dn_i + k_B \ln Q \sum_i dn_i \\ &= -\beta dU' + k_B \ln Q dn_{TOT} = \mathbf{0} \text{ by construction} \end{aligned}$$

- Compare this to

$$dS' = \frac{1}{T} dU' + \frac{P}{T} dV' - \frac{\mu}{T} dn$$

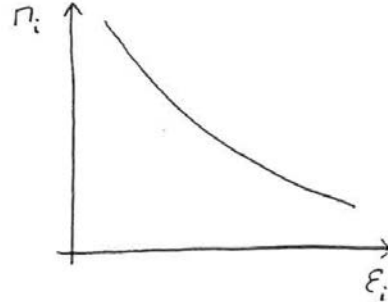
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$$\beta = -1/T$$

- With  $\beta = -1/T$ , we have the Boltzmann distribution

$$\frac{n_i}{n_{TOT}} = \frac{e^{\epsilon_i/k_B T}}{Q}$$

$$Q = \sum_i e^{-\epsilon_i/k_B T}$$



⇒ Distribution that maximizes entropy for an isolated system with  $n_{TOT}$  particles distributed over states with energy  $\epsilon_i$  according to  $\{n_i\}$

- Describes equilibrium state

- Maxwell distribution and likelihood  
e.g. How much less likely is it to find particle in state with energy  $\epsilon_m = \epsilon_l + \Delta\epsilon$  than in state with energy  $\epsilon_l$

$$\frac{n_m}{n_l} = \frac{e^{-(\epsilon_l + \Delta\epsilon)/k_B T}}{e^{-\epsilon_l/k_B T}} = e^{-\Delta\epsilon/k_B T}$$

- Depends on  $\Delta\epsilon/k_B T$
- Becomes less likely with energy splitting  $\Delta\epsilon$
- Becomes more likely with temperature  $T$

- Thermal energy  $k_B T$

$$k_B = R/N_A = 1.380 \times 10^{-23} \text{ J/k} = 8.617 \times 10^{-5} \text{ eV/k}$$

- sets energy scale for likely fluctuations  $\Delta\epsilon$
- $\Delta\epsilon \approx k_B T$  are likely
- $\Delta\epsilon \gg k_B T$  are unlikely
- at  $T = 298 \text{ K}$ ,  $k_B T = 4.112 \times 10^{-21} \text{ J} = 0.0257 \text{ eV} = 25.7 \text{ meV}$
- foundation for Arrhenius rate equation (hello 3.091 !!)

thermally activated processes

$$k = A e^{-E_a/k_B T}$$

k: rate;  $E_a$  : activation energy

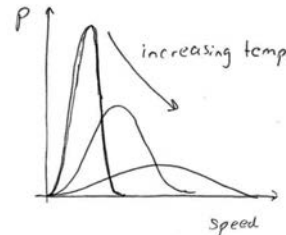
## 2 Maxwell-Boltzmann distribution

- Consider particles with kinetic energy  $\mathbf{p}^2/2m = (p_x^2 + p_y^2 + p_z^2)/2m$
- Distribution of energy  $f_E$  depends on both
  - Maxwell distribution  $n_i/n_{TOT} = Q^{-1} \exp(-\mathbf{p}_i^2/2mk_B T)$
  - Density of states in momentum space

$$d^3\mathbf{p} = 4\pi p^2 dp = 4\pi m \sqrt{2mE} dE$$

$$f_E = 2\sqrt{\frac{E}{\pi}} \left(\frac{1}{kT}\right)^{3/2} e^{-E/kT}$$

- more faster particles at higher temp
- vanishing number of particles with  $p = 0$



## 3 Ensembles

A group of things or people acting or taken together as a whole

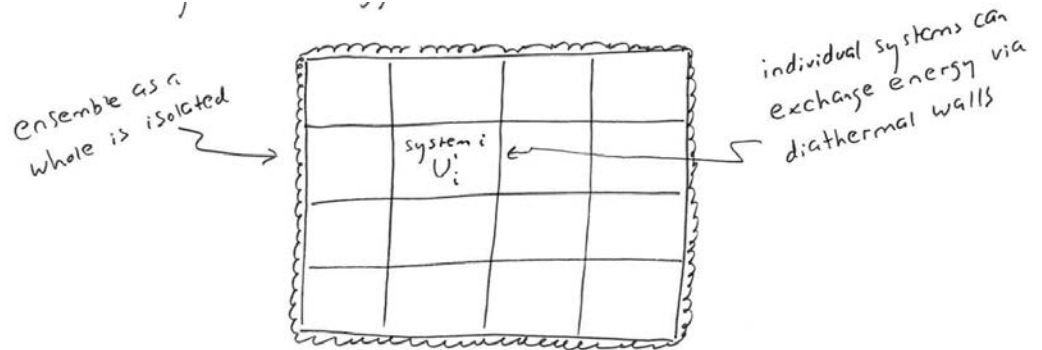
- Microcanonical ensemble: Set of all possible microstates of an isolated system with fixed  $U', V', n$ 
  - Microcanonical partition function  $Q = \sum_i e^{\epsilon_i/k_B T}$ , (sum over all possible single particle states)
  - Equilibrium (S maximized) for Boltzmann probability distribution

$$\mathcal{P}_i = \frac{n_i}{n_{TOT}} = \frac{e^{-\epsilon_i/k_B T}}{Q}$$

probability  $\mathcal{P}_i$  of finding single particles in state  $i$ , with energy  $\epsilon_i$

- Canonical ensemble: Set of all possible microstates of a system with fixed  $V', n$ 
  - System is closed and rigid, but can exchange energy with surroundings, i.e. diathermal walls

- System energy  $U'_\nu$  can fluctuate



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- canonical partition function  $Z = \sum_\nu e^{-U'_\nu/kT}$   
sum over possible states of the system
  - equilibrium (F minimized) for distribution

$$P_\nu = \frac{e^{-U'_\nu/k_B T}}{Z} \quad \text{Boltzmann distribution for canonical ensemble}$$

$P_\nu$ : prob. of finding system in state  $\nu$  with energy  $U'_\nu$

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## 4 The partition functions

- Describe equilibrium properties and fluctuations
- Basis for calculating thermodynamic potentials and properties

$$\begin{aligned} \text{e.g. } \langle U \rangle &= -\frac{\partial \ln Z}{\partial (1/kT)} \\ C_V &= \frac{1}{kT^2} \frac{\partial^2 \ln Z}{\partial (1/kT)^2} \\ S' &= \frac{\langle U \rangle}{T} + k_B \ln Z \\ F' &= -k_B T \ln Z \\ \langle U^2 \rangle - \langle U \rangle^2 &= kT^2 \frac{\partial \langle U \rangle}{\partial T} \\ &\dots \end{aligned}$$

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