



MIT 3.071

Amorphous Materials

10: Electrical and Transport Properties

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After-class reading list

- Fundamentals of Inorganic Glasses
 - Ch. 14, Ch. 16
- Introduction to Glass Science and Technology
 - Ch. 8
- 3.024 band gap, band diagram, engineering conductivity

Basics of electrical conduction

■ Electrical conductivity

$$\sigma = \sum_i n_i Z_i e \mu_i$$

$$v_i = \mu_i E$$

■ Einstein relation

$$D_i = \frac{\mu_i k_B T}{Z_i e}$$

$$\Rightarrow \sigma = \frac{1}{k_B T} \cdot \sum_i (Z_i e)^2 n_i D_i$$

σ : electrical conductivity

n : charge carrier density

Z : charge number

e : elementary charge

μ : carrier mobility

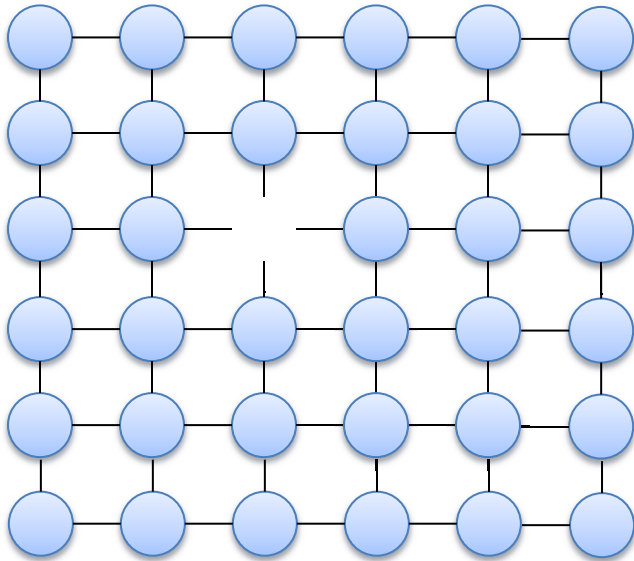
D : diffusion coefficient

v : carrier drift velocity

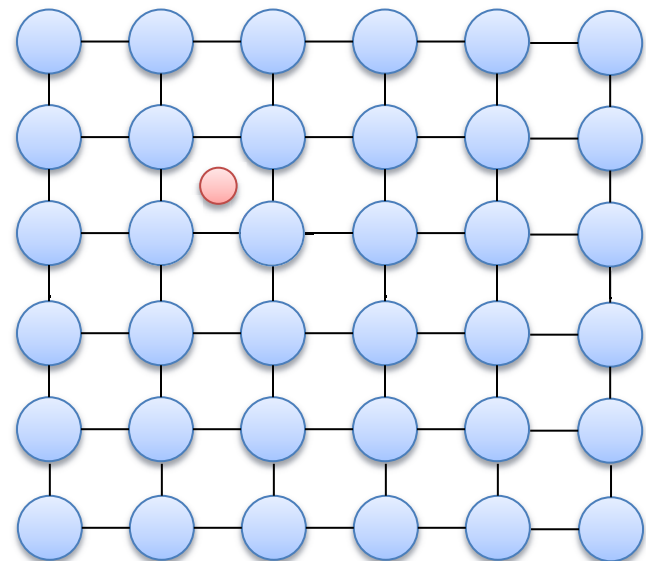
E : applied electric field

Both ions and electrons
contribute to electrical
conductivity in glasses

Ionic conduction in crystalline materials



Vacancy mechanism



Interstitial mechanism

Ionic conduction pathway in amorphous solids

[Comparison of Li transport pathways](#) figure removed due to copyright restrictions. See: Figure 8: Adams, S. and R. Prasada Rao. "Transport Pathways for Mobile Ions in Disordered Solids from the Analysis of Energy-scaled Bond-valence Mismatch Landscapes." *Phys. Chem. Chem Phys.* 11 (2009): 3210-3216.

- There are low energy “sites” where ions preferentially locate
- Ionic conduction results from ion transfer between these sites
- Ionic conduction is thermally activated

2-D slices of regions with Li site energies below a threshold value in $\text{Li}_2\text{O-SiO}_2$ glasses

Phys. Chem. Chem. Phys. **11**, 3210 (2009)

A tale of two valleys

- Assuming completely random hops, the average total distance an ion moves after M hops in 1-D is:

$$r = d \cdot \sqrt{M}$$

- Average diffusion distance:

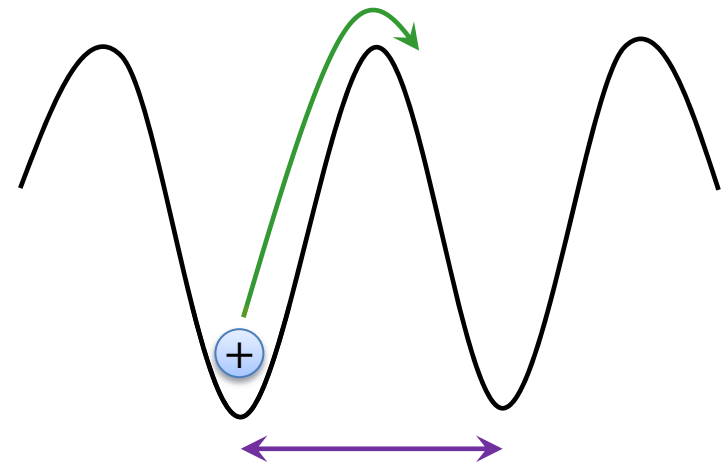
$$r = \sqrt{2\tau D} \quad (1\text{-D})$$

$$r = \sqrt{6\tau D} \quad (3\text{-D})$$

$$\Rightarrow D = \frac{1}{2} \nu d^2 \quad (1\text{-D})$$

$$\Rightarrow D = \frac{1}{6} \nu d^2 \quad (3\text{-D})$$

Electric field $E = 0$



Average spacing between adjacent sites: d

- For correlated hops:

$$D = \frac{1}{6} f \nu d^2 \quad 0 < f < 1$$

Ion hopping frequency: ν

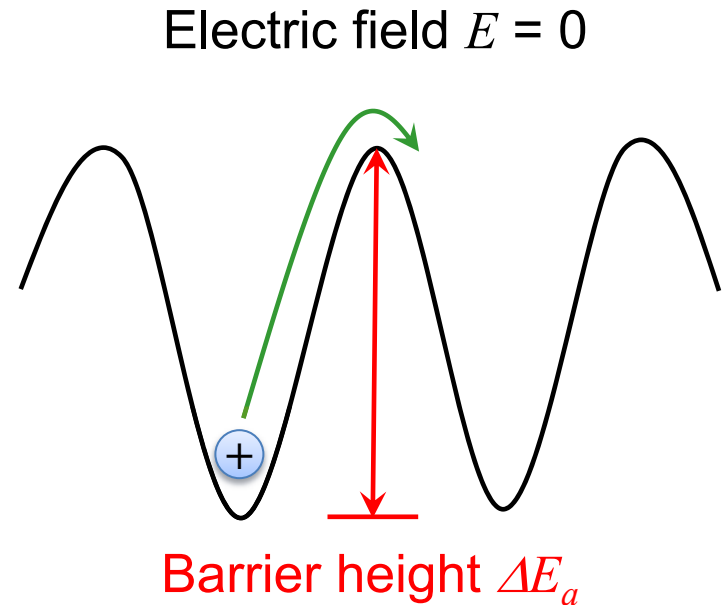
A tale of two valleys

- Attempt (vibration) frequency: ν_0
- Frequency of successful hops (ion hopping frequency):

$$\nu = \nu_0 \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

$$\Rightarrow D = \frac{1}{6} f \nu_0 d^2 \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

- Equal probability of hopping along all directions: zero net current



A tale of two valleys

- Energy difference between adjacent sites: $ZeEd$
- Hopping frequency \rightarrow :

$$v_{\rightarrow} = \frac{1}{2} v_0 \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

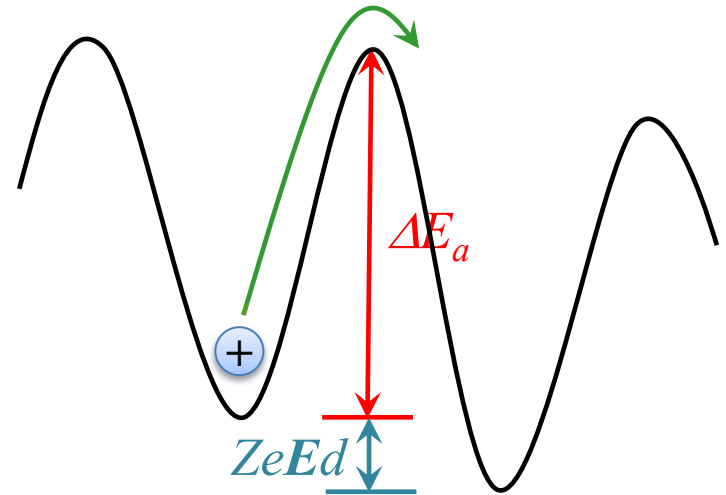
- Hopping frequency \leftarrow :

$$v_{\leftarrow} = \frac{1}{2} v_0 \exp\left(-\frac{\Delta E_a + ZeEd}{k_B T}\right)$$

- Net ion drift velocity:

$$v = (v_{\rightarrow} - v_{\leftarrow}) \cdot d = \frac{v_0 ZeEd^2}{2k_B T} \cdot \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

Electric field $E > 0$



A tale of two valleys

- Ion mobility

$$\mu = \frac{v_0 Zed^2}{2k_B T} \cdot \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

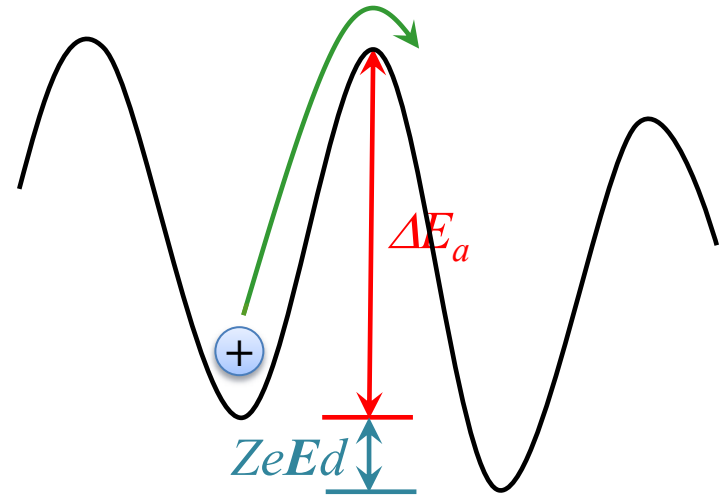
- Electrical conductivity (1-D, random hop)

$$\sigma = \frac{nv_0 (Zed)^2}{2k_B T} \cdot \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

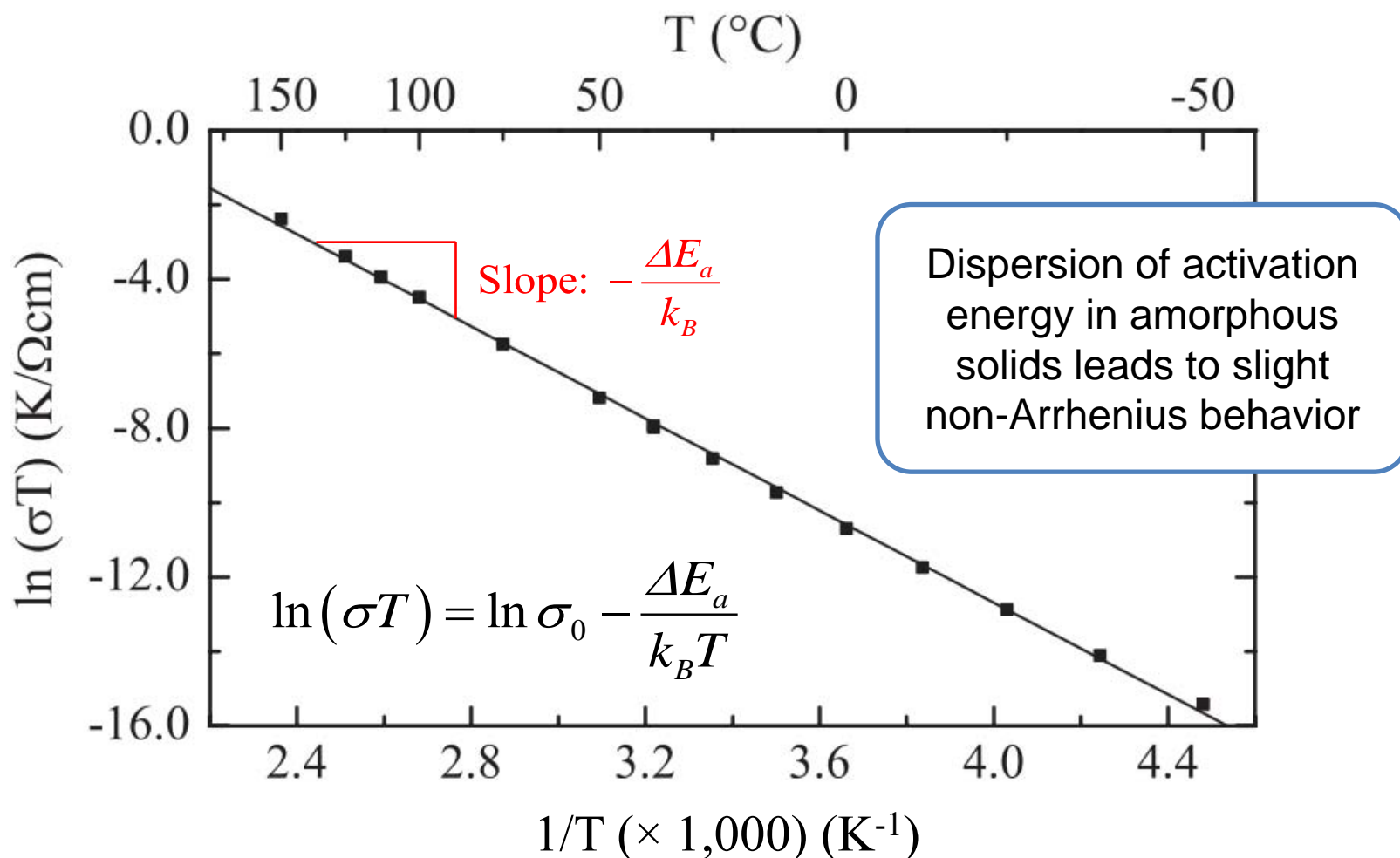
- Einstein relation (3-D, correlated hops)

$$\sigma = \frac{1}{k_B T} (Ze)^2 nD = \frac{fnv_0 (Zed)^2}{6k_B T} \cdot \exp\left(-\frac{\Delta E_a}{k_B T}\right) = \frac{\sigma_0}{T} \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

Electric field $E > 0$



Temperature dependence of ionic conductivity



Phys. Rev. Lett. **109**, 075901 (2012)

Theoretical ionic conductivity limit in glass

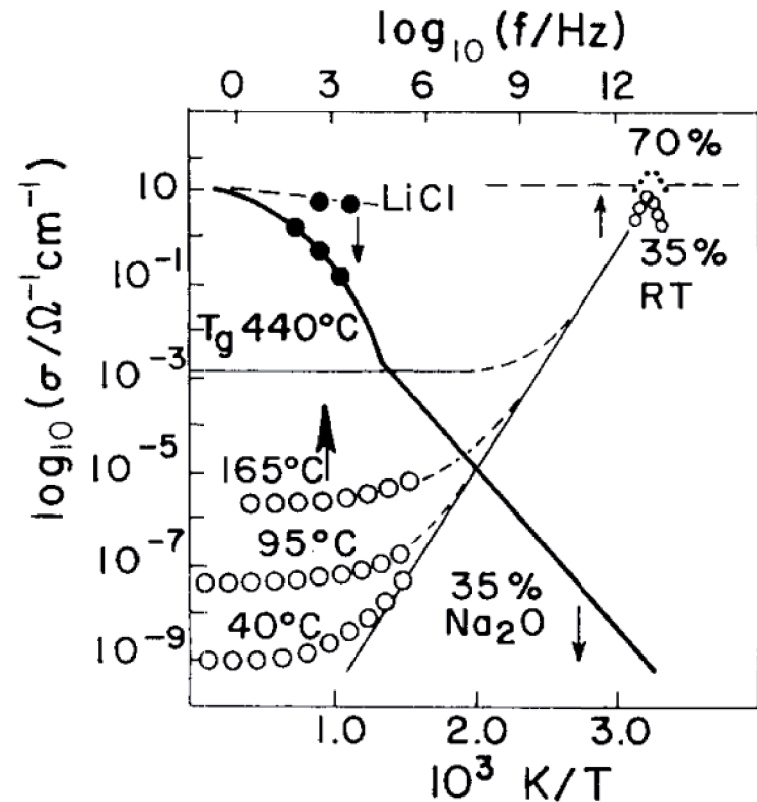
$$\sigma = \frac{\sigma_0}{T} \cdot \exp\left(-\frac{\Delta E_a}{k_B T}\right)$$

$$\sigma_0 = \frac{f n v_0 (Z e d)^2}{6 k_B}$$

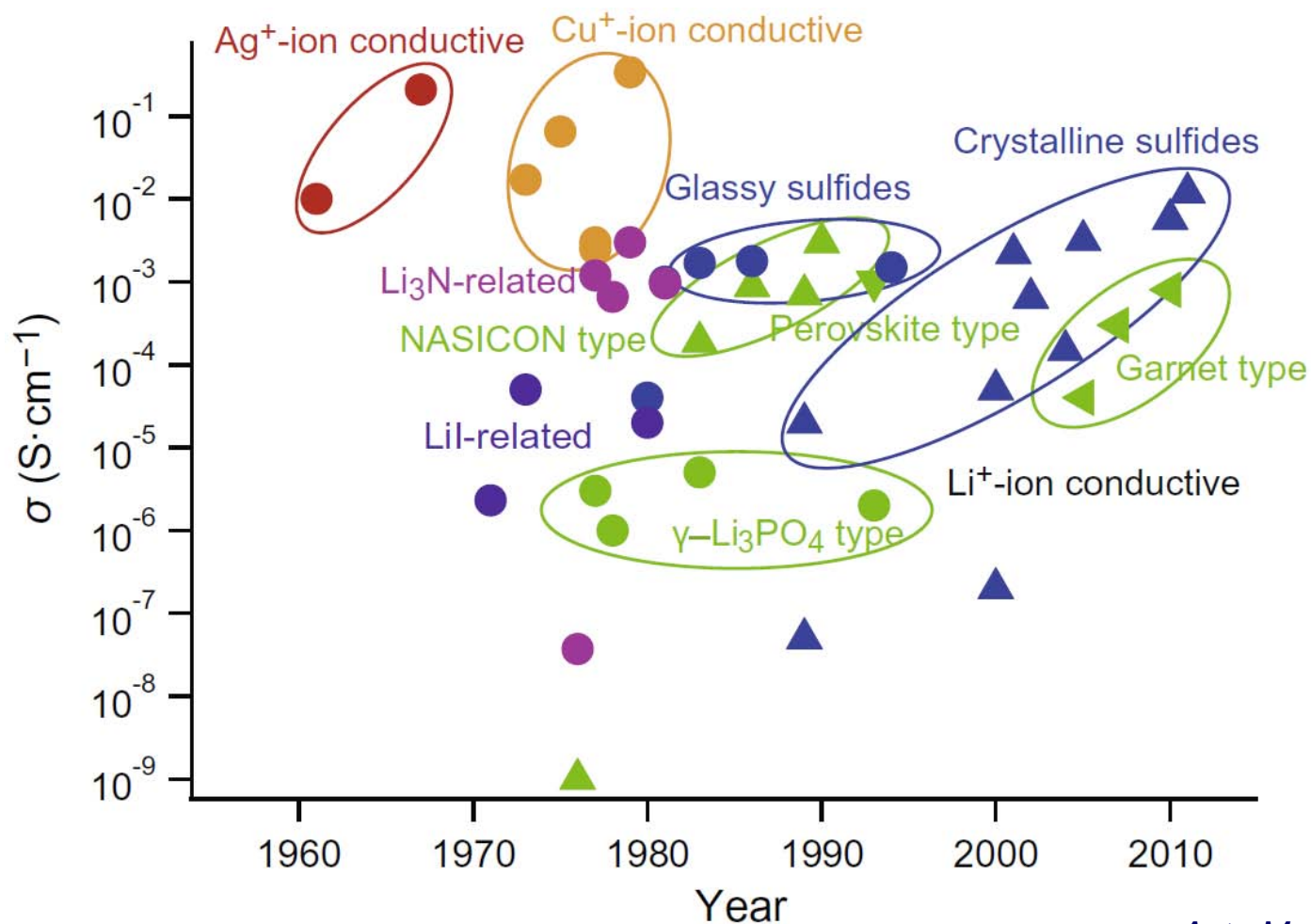
Note that σ_0 has a unit of $\Omega \cdot \text{cm}/\text{K}$

- When $T \rightarrow \infty$, $\sigma \rightarrow \sigma_0/T$
- Extrapolation of the Arrhenius plot agrees with infrared spectroscopic measurements in ionic liquids (molten salts)

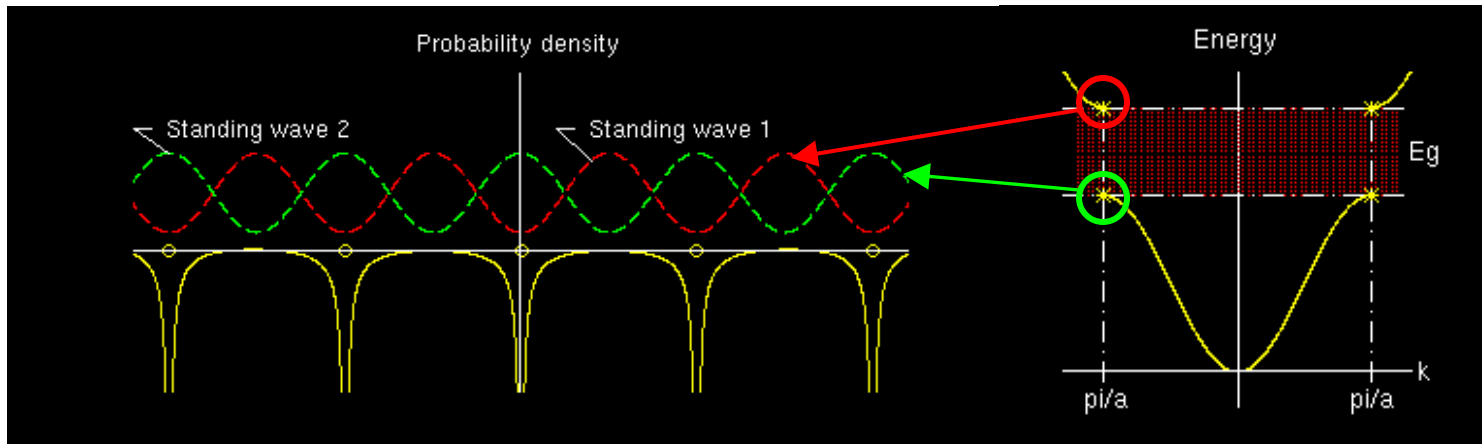
Solid State Ionics **18&19**, 72 (1986)
Annu. Rev. Phys. Chem. **43**, 693 (1992)



Fast ion conductors / superionic conductors



Band structures in defect-free crystalline solids



- ✓ All electronic states are labeled with **real** Bloch wave vectors k signaling translational symmetry
- ✓ All electronic states are **extended states**
- ✓ No extended states exist in the band gap

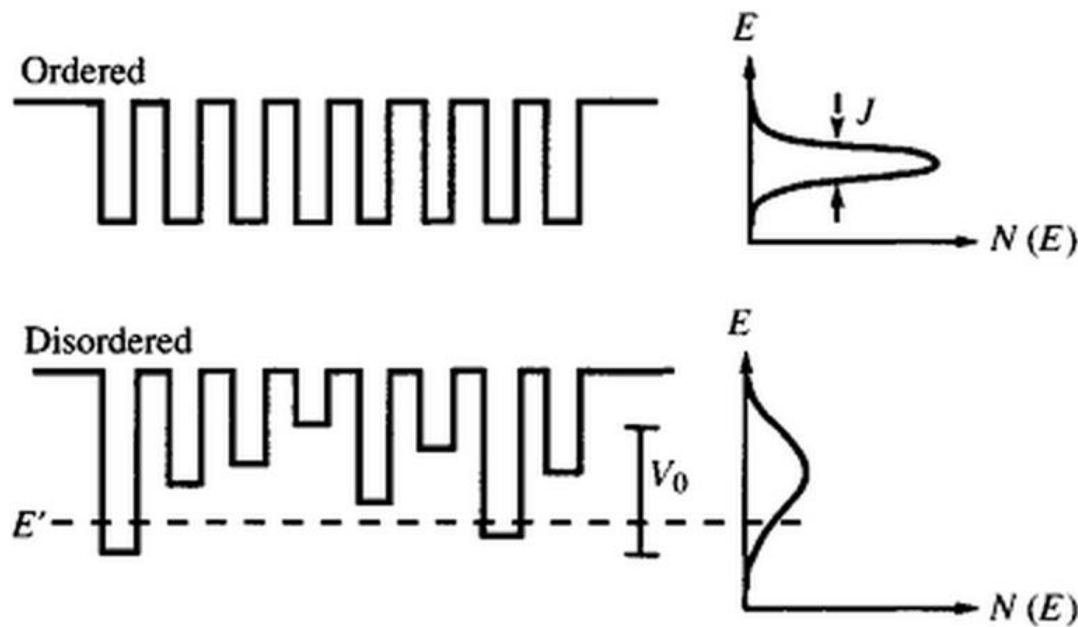
Band structures in defect-free crystalline solids

Figure removed due to copyright restrictions. See Figure 12, Chapter 7:
Kittel, Charles. *Introduction to Solid State Physics*. Wiley, 2005.

In the band gap, wave equation solutions have **complex** wave vectors k

Anderson localization in disordered systems

- Localization criterion: $V_0 / J > 3$

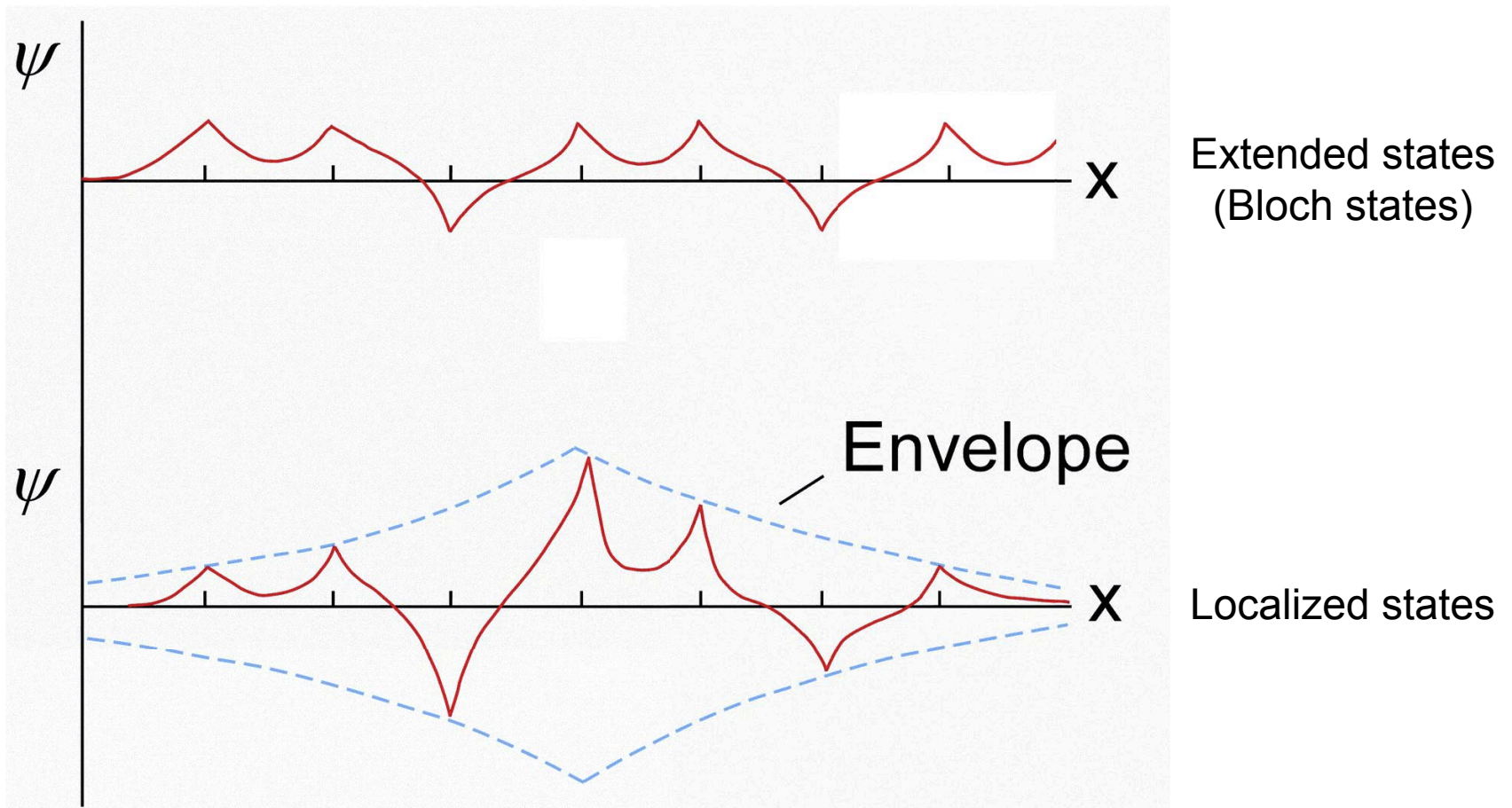


P. W. Anderson

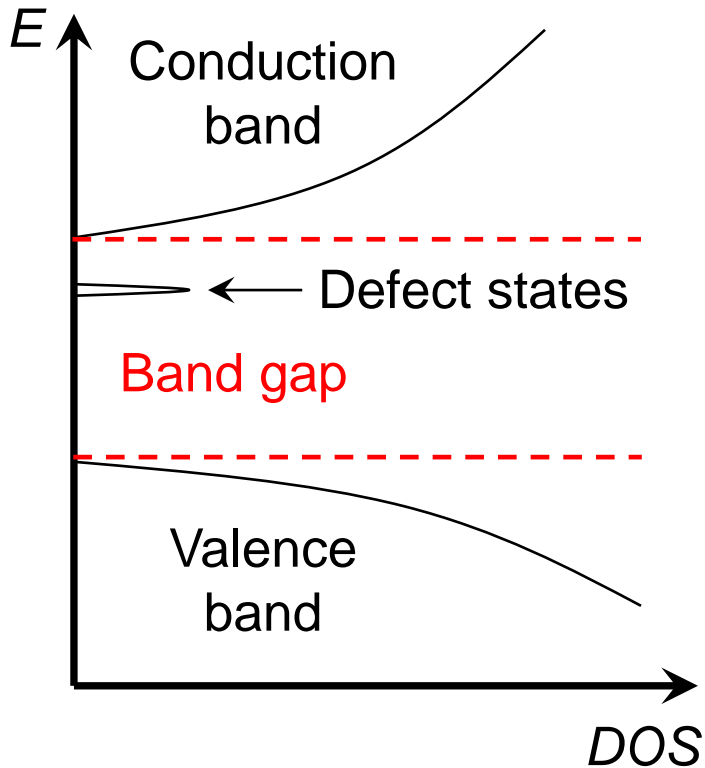
Image is in the public domain.
Source: [Wikimedia Commons](#).

Disorder leads to (electron, photon, etc.) wave function localization

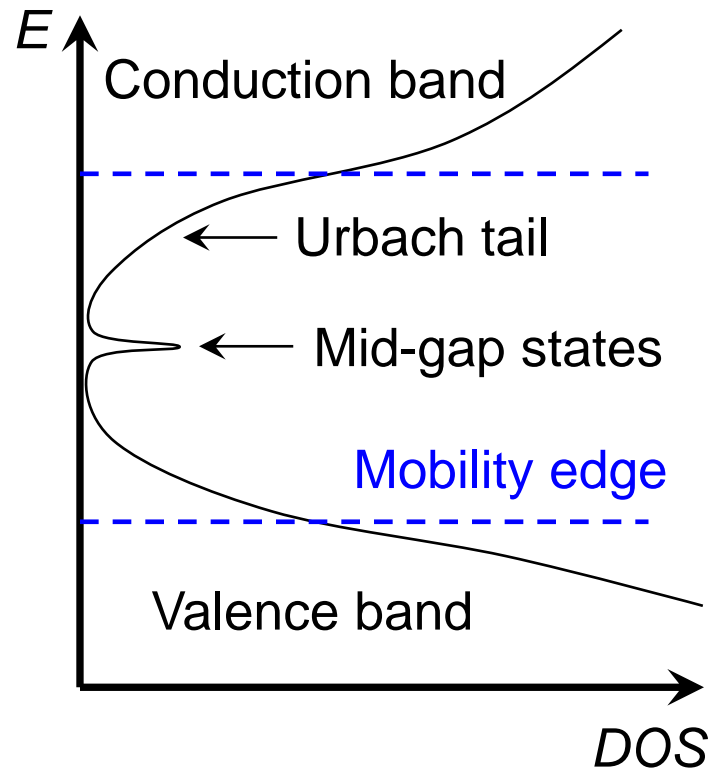
Anderson localization in disordered systems



Density of states (DOS) in crystalline and amorphous solids



Crystalline solids



Amorphous solids

Extended state conduction

- Extended state conductivity:

$$\sigma_{ex} = ne\mu_{ex}$$

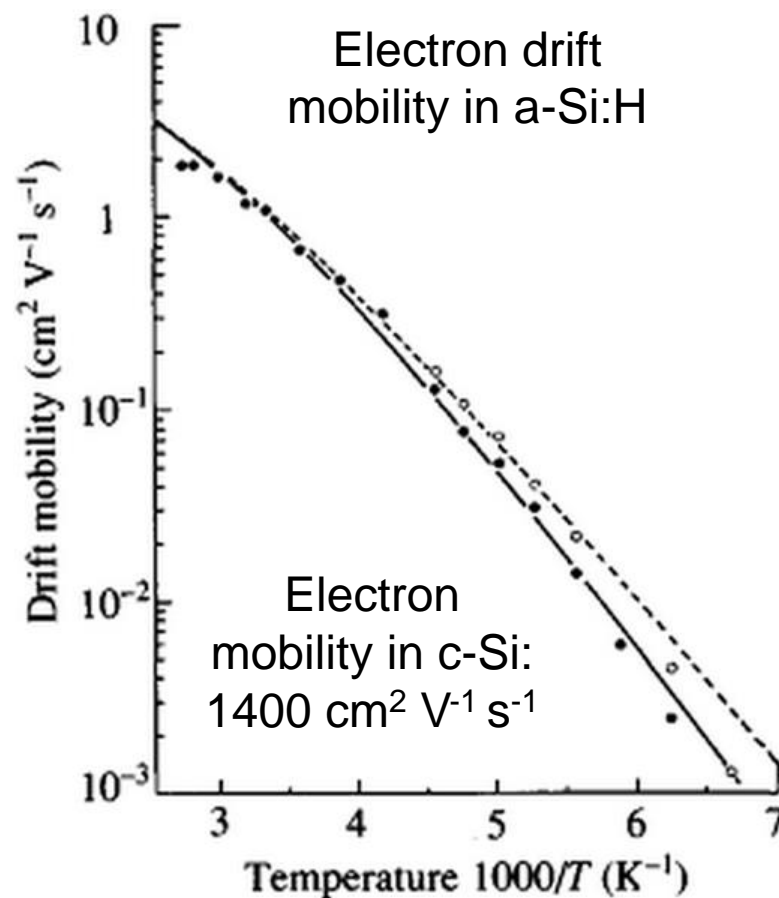
$$\mu_{ex} = (1 - f_{trap}) \mu_0$$

μ_0 : free mobility

f_{trap} : fraction of time in trap states

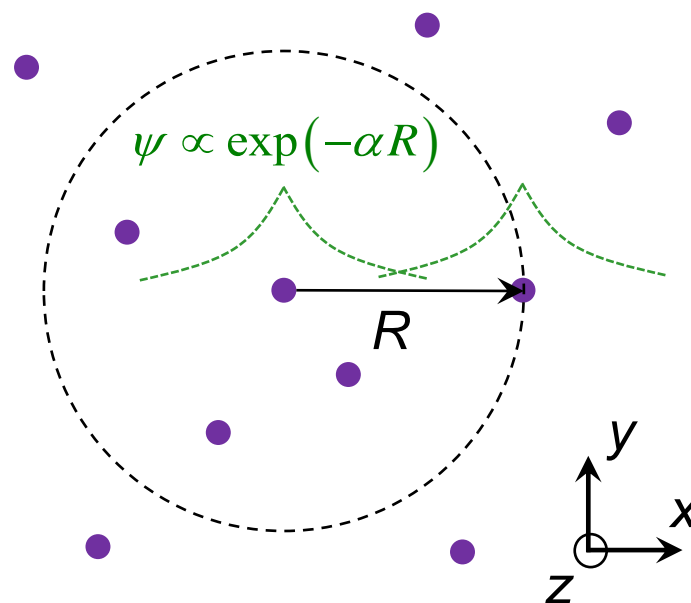
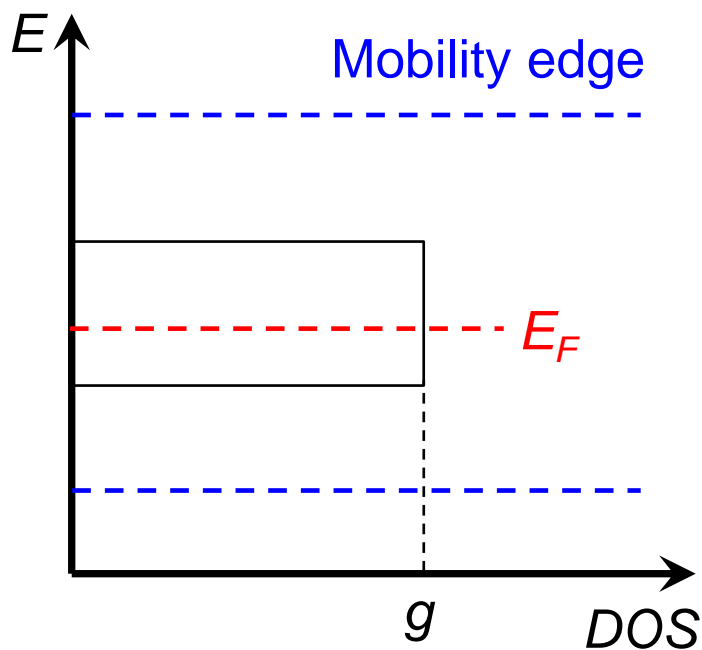
- Drift mobility μ_{ex} increases with temperature ($T \rightarrow \infty, f_{trap} \rightarrow 0$)
- Extended state conductivity follows Arrhenius dependence

R. Street, *Hydrogenated Amorphous Silicon*, Ch. 7



Hopping conduction via localized states

- Fixed range hopping: hopping between nearest neighbors
 - Hopping between dopant atoms at low temperature
- Variable range hopping (VRH)
 - Hopping between localized states near E_F



Variable range hopping

- Hopping probability $P \propto \exp\left(-2\alpha R - \frac{\Delta E}{k_B T}\right) \propto \sigma_{VRH}$

- Within distance R , the average minimal energy difference ΔE is:

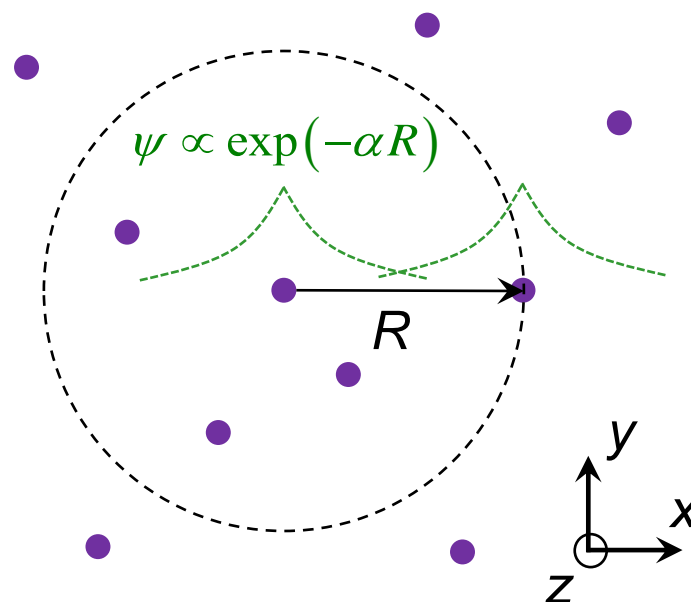
$$\Delta E = \frac{3}{4\pi R^3 g}$$

- Optimal hopping distance:

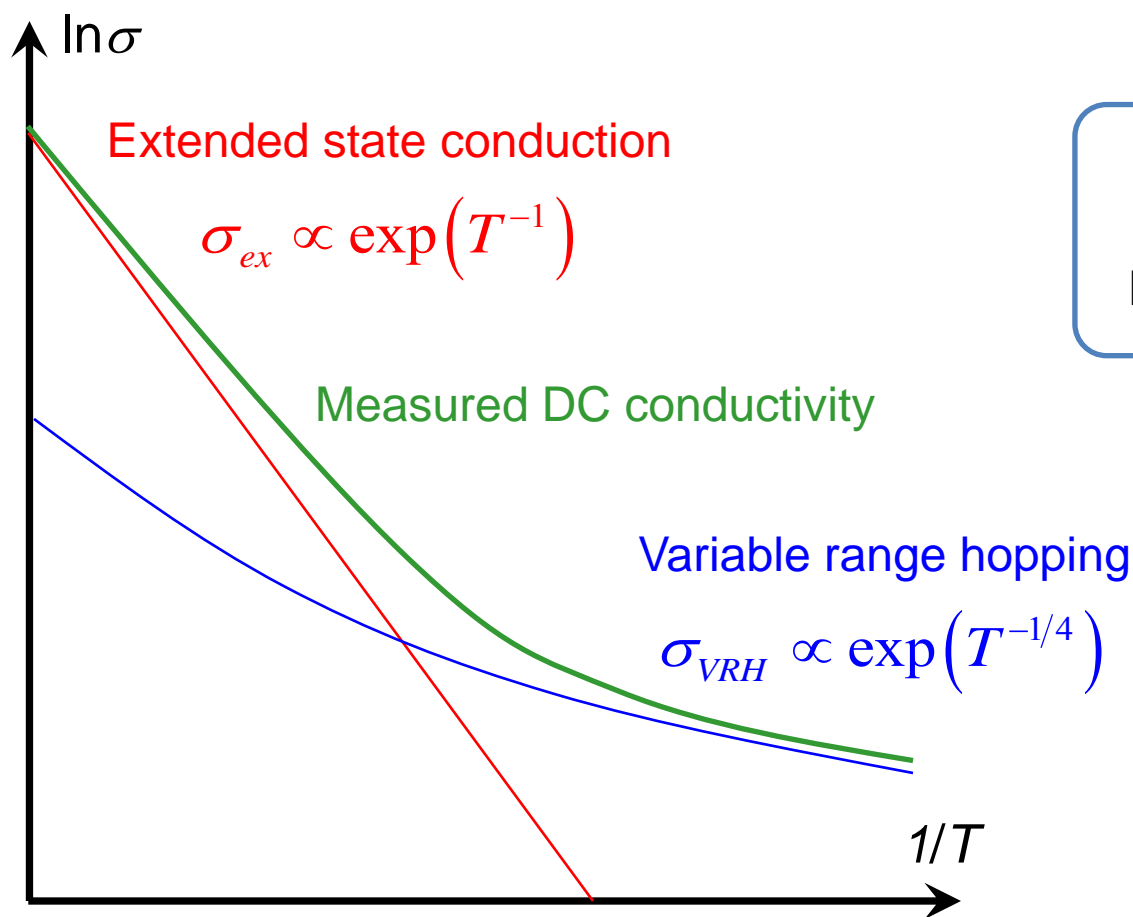
$$R = (8\pi g \alpha k_B T / 9)^{-1/4}$$

$$\left(2\alpha R + \frac{\Delta E}{k_B T}\right)_{\min} = 4\alpha^{3/4} \left(\frac{2}{9\pi g k_B T}\right)^{1/4}$$

$$\Rightarrow \sigma_{VRH} \propto \exp(T^{-1/4})$$

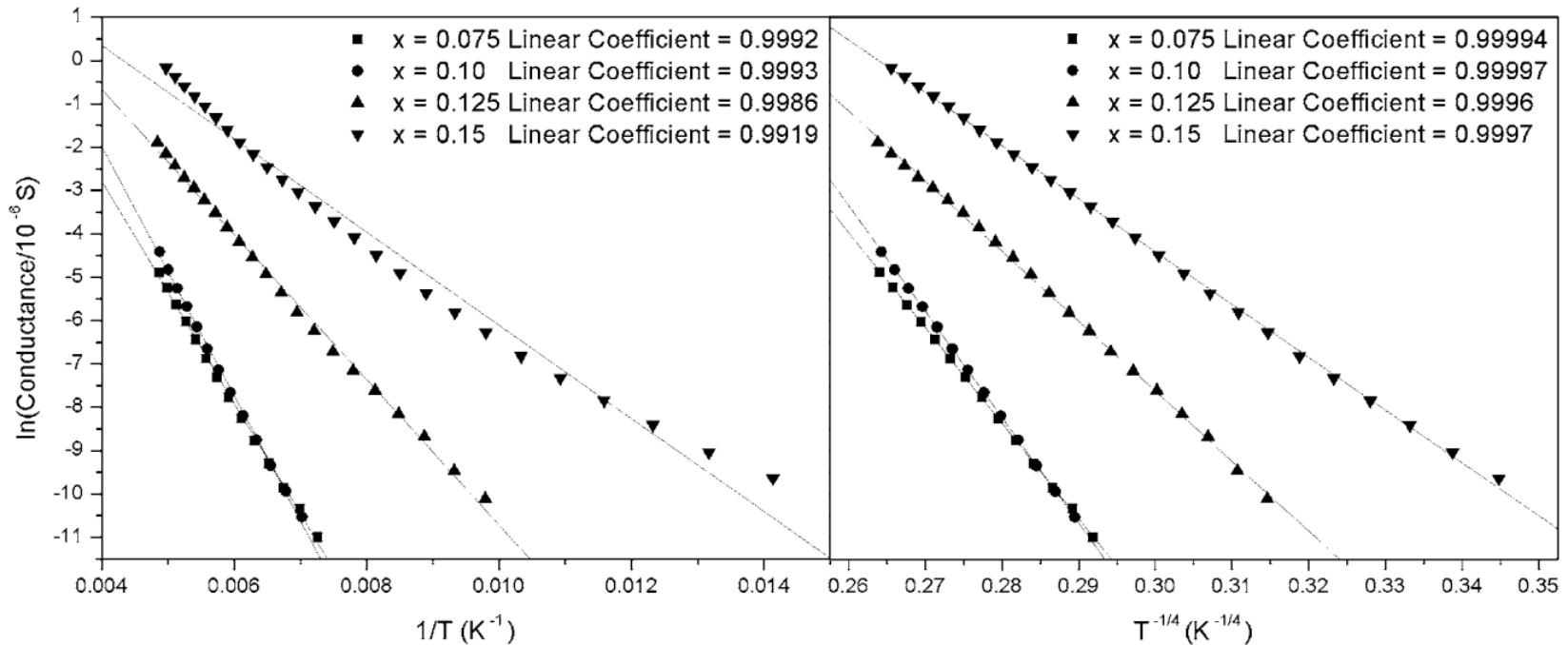


DC conductivity in amorphous semiconductors



VRH is most pronounced at low temperature

VRH in As-Se-Te-Cu glass



- ✓ Near room temperature, mixed ionic and extended state conduction
- ✓ At low temperature, variable range hopping dominates

Summary

■ Basics of electrical transport

- Conductivity: scalar sum of ionic and electronic contributions

$$\sigma = \sum_i n_i Z_i e \mu_i \quad v_i = \mu_i \mathbf{E}$$

- Einstein relation

$$D_i = \frac{\mu_i k_B T}{Z_i e} \Rightarrow \sigma = \frac{1}{k_B T} \cdot \sum_i (Z_i e)^2 n_i D_i$$

■ Ionic conductivity

- Occurs through ion hopping between different preferred “sites”
- Thermally activated process and non-Arrhenius behavior

$$\sigma = \frac{\sigma_0}{T} \cdot \exp\left(-\frac{\Delta E_a}{k_B T}\right) \quad \sigma_0 = \frac{f n v_0 (Z e d)^2}{6 k_B}$$

Summary

■ Electronic structure of amorphous semiconductors

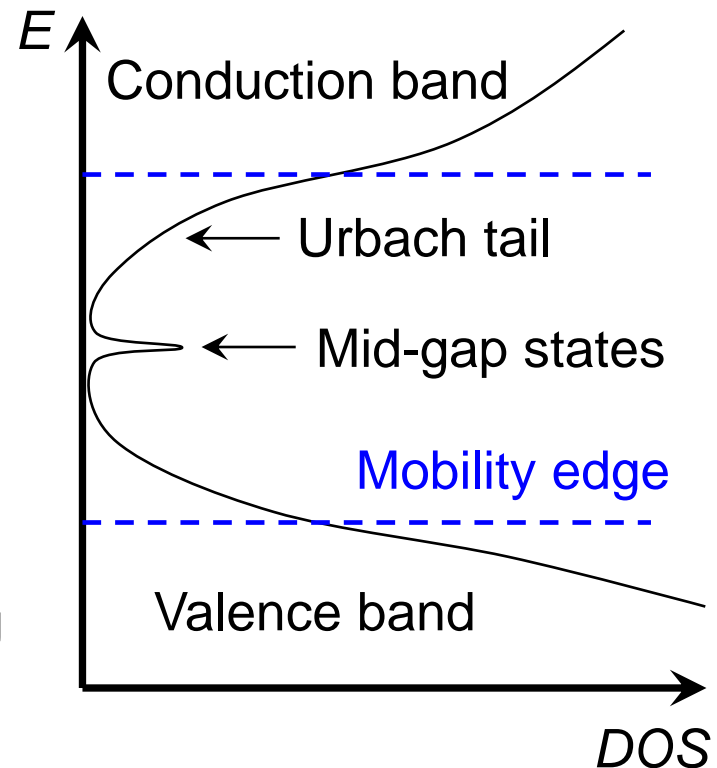
- Anderson localization: extended vs. localized states
- Density of states
- Mobility edge
- Band tail and mid-gap states

■ Extended state conduction

- Free vs. drift mobility
- Thermally activated process

■ Localized state conduction

- Fixed vs. variable range hopping
- Mott's $T^{-1/4}$ law of VRH



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