Massachusetts Institute of Technology
Department of Electrical Engineering & Computer Science
6.041/6.431: Probabilistic Systems Analysis
(Fall 2010)

Recitation 24: Solutions
December 7, 2010

1. (a) Normalization of the distribution requires:

$$1 = \sum_{k=0}^{\infty} p_K(k;\theta) = \sum_{k=0}^{\infty} \frac{e^{-k/\theta}}{Z(\theta)} = \frac{1}{Z(\theta)} \sum_{k=0}^{\infty} e^{-k/\theta} = \frac{1}{Z(\theta)\cdot(1-e^{-1/\theta})},$$

so $Z(\theta) = \frac{1}{1-e^{-1/\theta}}$.

(b) Rewriting $p_K(k;\theta)$ as:

$$p_K(k;\theta) = \left(e^{-1/\theta}\right)^k \left(1-e^{-1/\theta}\right), \qquad k = 0, 1, \ldots$$

the probability distribution for the photon number is a geometric probability distribution with probability of success $p = 1 - e^{-1/\theta}$, and it is shifted with 1 to the left since it starts with $k = 0$. Therefore the photon number expectation value is

$$\mu_K = \frac{1}{p} - 1 = \frac{1}{1-e^{-1/\theta}} - 1 = \frac{1}{e^{1/\theta}-1}$$

and its variance is

$$\sigma_K^2 = \frac{1-p}{p^2} = \frac{e^{-1/\theta}}{(1-e^{-1/\theta})^2} = \mu_K^2 + \mu_K.$$

(c) The joint probability distribution for the k_i is

$$p_K(k_1,...,k_n;\theta) = \frac{1}{Z(\theta)^n}\Pi_{i=1}^n e^{-k_i/\theta} = \frac{1}{Z(\theta)^n} e^{-\frac{1}{\theta}\sum_{i=1}^n k_i}.$$

The log likelihood is $-n \cdot \log Z(\theta) - 1/\theta \sum_{i=1}^n k_i$.
We find the maxima of the log likelihood by setting the derivative with respect to the parameter θ to zero:

$$\frac{d}{d\theta} \log p_K(k_1,...,k_n;\theta) = -n \cdot \frac{e^{-1/\theta}}{\theta^2(1-e^{-1/\theta})} + \frac{1}{\theta^2}\sum_{i=1}^n k_i = 0$$

or

$$\frac{1}{e^{1/\theta}-1} = \frac{1}{n}\sum_{i=1}^n k_i = s_n.$$

For a hot body, $\theta \gg 1$ and $\frac{1}{e^{1/\theta}-1} \approx \theta$, we obtain

$$\theta \approx \frac{1}{n}\sum_{i=1}^n k_i = s_n.$$

Thus the maximum likelihood estimator $\hat{\Theta}_n$ for the temperature is given in this limit by the sample mean of the photon number

$$\hat{\Theta}_n = \frac{1}{n}\sum_{i=1}^{n} K_i.$$

(d) According to the central limit theorem, the sample mean for large enough n (in the limit) approaches a Gaussian distribution with standard deviation our root mean square error

$$\sigma_{\hat{\Theta}_n} = \frac{\sigma_K}{\sqrt{n}}.$$

To allow only for 1% relative root mean square error in the temperature, we need $\frac{\sigma_K}{\sqrt{n}} < 0.01\mu_K$. With $\sigma_K^2 = \mu_K^2 + \mu_K$ it follows that

$$\sqrt{n} > \frac{\sigma_K}{0.01\mu_K} = 100\frac{\sqrt{\mu_K^2 + \mu_K}}{\mu_K} = 100\sqrt{1 + \frac{1}{\mu_K}}.$$

In general, for large temperatures, i.e. large mean photon numbers $\mu_K \gg 1$, we need about 10,000 samples.

(e) The 95% confidence interval for the temperature estimate for the situation in part (d), i.e.

$$\sigma_{\hat{\Theta}_n} = \frac{\sigma_K}{\sqrt{n}} = 0.01\mu_K,$$

is

$$[\hat{K} - 1.96\sigma_{\hat{K}}, \hat{K} + 1.96\sigma_{\hat{K}}] = [\hat{K} - 0.0196\mu_K, \hat{K} + 0.0196\mu_K].$$

2. (a) Using the regression formulas of Section 9.2, we have

$$\hat{\theta}_1 = \frac{\sum_{i=1}^{5}(x_i - \bar{x})(y_i - \bar{y})}{\sum_{i=1}^{5}(x_i - \bar{x})^2}, \qquad \hat{\theta}_0 = \bar{y} - \hat{\theta}_1\bar{x},$$

where

$$\bar{x} = \frac{1}{5}\sum_{i=1}^{5} x_i = 4.94, \qquad \bar{y} = \frac{1}{5}\sum_{i=1}^{5} y_i = 134.38.$$

The resulting ML estimates are

$$\hat{\theta}_1 = 40.53, \qquad \hat{\theta}_0 = -65.86.$$

(b) Using the same procedure as in part (a), we obtain

$$\hat{\theta}_1 = \frac{\sum_{i=1}^{5}(x_i^2 - \bar{x})(y_i - \bar{y})}{\sum_{i=1}^{5}(x_i^2 - \bar{x})^2}, \qquad \hat{\theta}_0 = \bar{y} - \hat{\theta}_1\bar{x},$$

where

$$\bar{x} = \frac{1}{5}\sum_{i=1}^{5} x_i^2 = 33.60, \qquad \bar{y} = \frac{1}{5}\sum_{i=1}^{5} y_i = 134.38.$$

which for the given data yields

$$\hat{\theta}_1 = 4.09, \qquad \hat{\theta}_0 = -3.07.$$

Figure 1 shows the data points (x_i, y_i), $i = 1, \ldots, 5$, the estimated linear model

$$y = 40.53x - 65.86,$$

and the estimated quadratic model

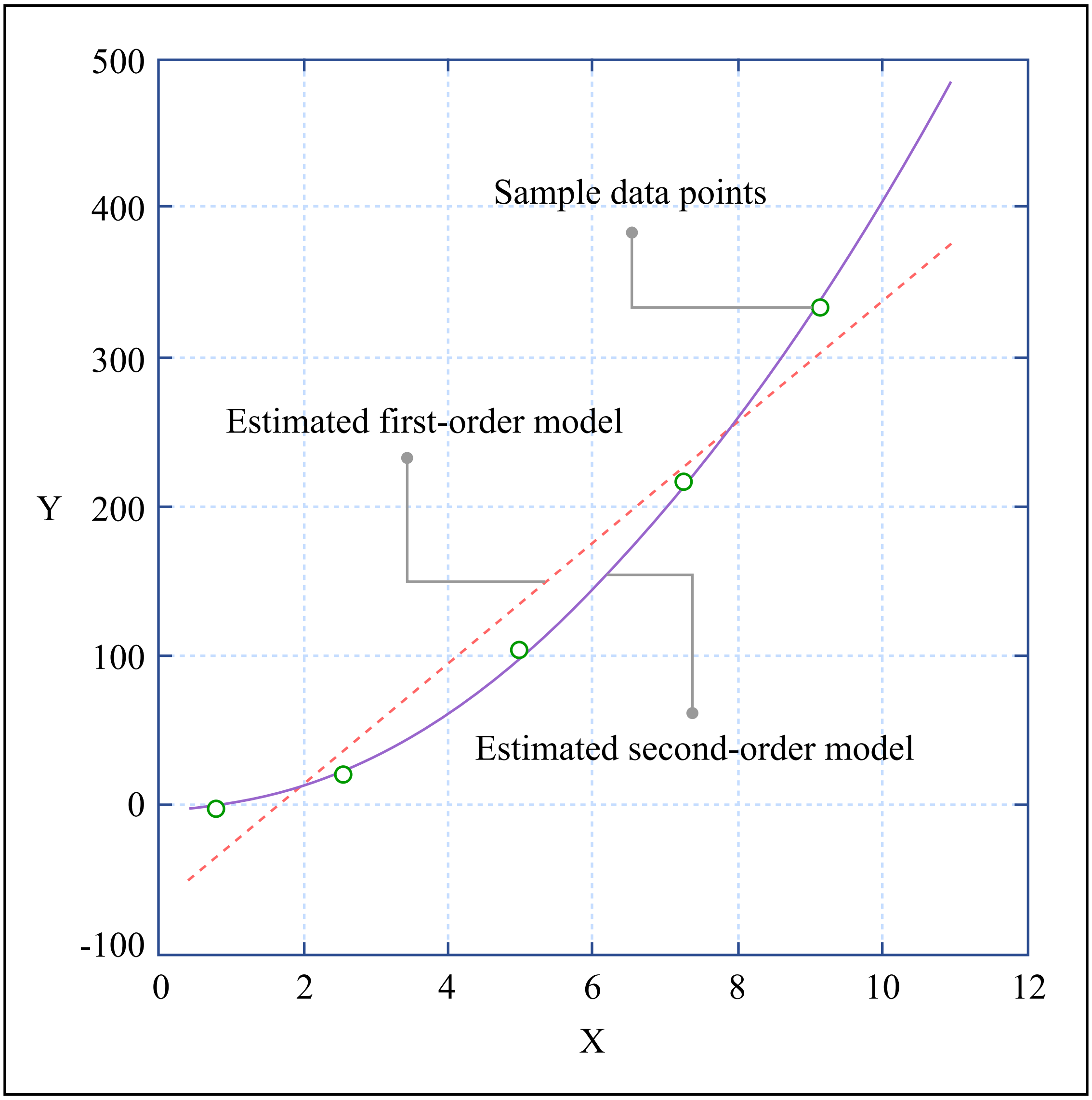

Figure 1: Regression Plot

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