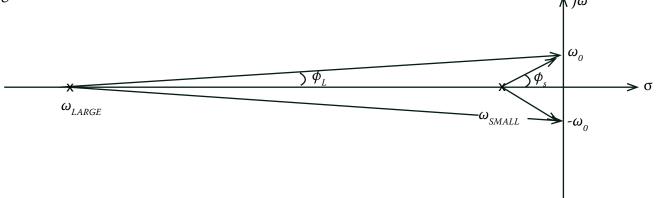
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In lecture, Prof. Roberge spoke of building elaborate systems using op-amp circuits. When we do this, we habitually make statements like, "...and we'll choose the dynamics such that the poles contributed by the op-amps themselves are negligible." Let's discuss.

Consider a two-pole system whose behavior interests us only in the range from DC to ω_0 . A pole-zero diagram:



The angle ϕ_L is small compared to $\phi_{s'} \Rightarrow \text{pole} @ \omega_{LARGE}$ contributes very little phase shift.

For the frequency response evaluated at $s=j\omega_{a}$:

$$\frac{1}{\left(\frac{s}{\omega_{L}}+1\right)\left(\frac{s}{\omega_{S}}+1\right)} = \frac{1}{\left(\frac{j\omega_{0}}{\omega_{L}}+1\right)\left(\frac{j\omega_{0}}{\omega_{S}}+1\right)}$$

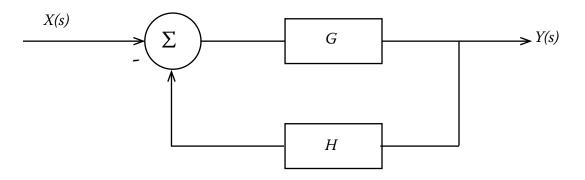
$$\frac{\omega_0}{\omega_L} <<1 \qquad \qquad \frac{\omega_0}{\omega_S} \quad \text{is on the order of unity (at least)!}$$

So this transfer function is well-approximated by the single-pole transfer function $\frac{1}{\left(\frac{s}{\omega_{S}}+1\right)}$ for frequencies from DC to ω_{0} .

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But what determines ω_0 ? Depends on the context, but for feedback systems we often look at the loop transmission.



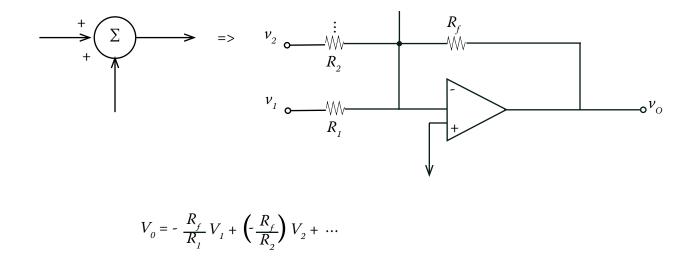
When $|L(s)| = |GH| \ll 1$, we arguably have an open-loop system:

$$\frac{Y(s)}{X(s)} = \frac{G}{1+GH} \approx G$$

So when we're looking for dynamics to ignore, we will often discard poles that are large compared to the loop crossover frequency, or the frequency at which |L(S)| = 1.

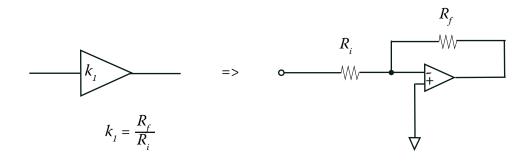
Op-amp circuits for modeling our systems

We need an integrator, a summer, an inversion, and a gain.

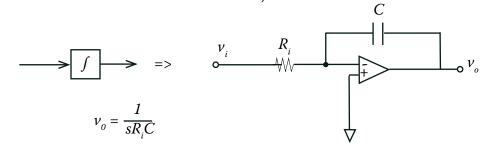


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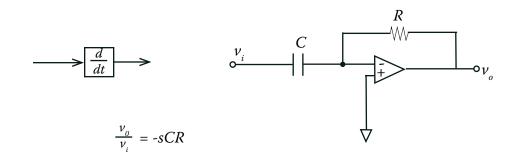
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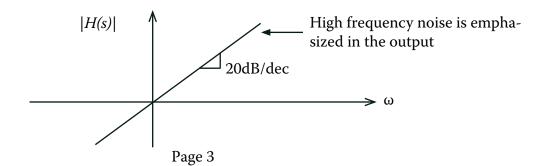
(Note that this also covers us for an inversion.)



And, a block that we do not use:



Difficult to manage in a noisy world:



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Also against the differentiator: it's hard enough to get high gain at DC. High gain at high frequencies? Forget it.

Now, on to building analog computers. Suppose we have an all-pole system. It begins as a differential equation:

$$a_n \frac{d^n y}{dt^n} + a_{n-1} \frac{d^{n-1} y}{dt^{n-1}} + a_{n-2} \frac{d^{n-2} y}{dt^{n-2}} + \dots + a_0 y = x$$

Completely general procedure starts with taking the Laplace transform:

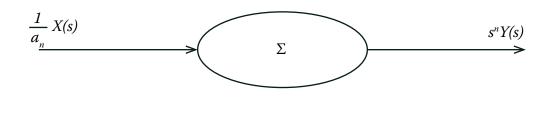
$$(a_{n}s^{n} + a_{n-1}s^{n-1} + a_{n-2}s^{n-2} + \dots + a_{n})Y(s) = x(s)$$

The system function, BTW, is: $\frac{Y(s)}{x(s)} = \frac{1}{a_{n}s^{n} + a_{n-1}s^{n-1} + a_{n-2}s^{n-2} + \dots + a_{n}}$

Solve for the highest order derivatve:

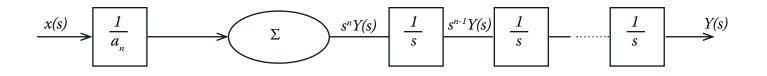
$$a_{n}s^{n}Y(s) = X(s) - (a_{n-1}s^{n-1} + a_{n-2}s^{n-2} + \dots + a_{0})Y(s)$$
$$s^{n}Y(s) = \frac{1}{a_{n}}X(s) - \left[\frac{a_{n-1}}{a_{n}}s^{n-1} + \frac{a_{n-2}}{a_{n}}s^{n-2} + \dots + \frac{a_{0}}{a_{n}}\right]Y(s)$$

Put down a big summing junction:

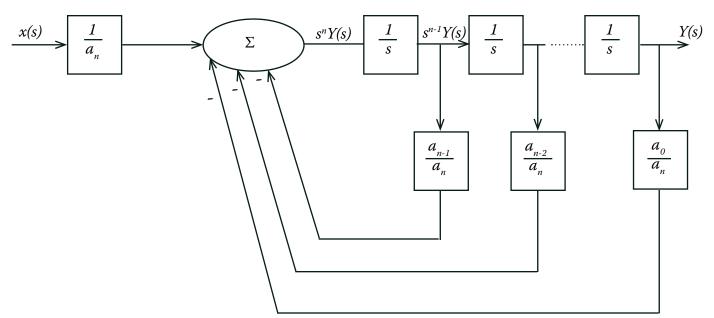


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Generate the derivatives that you need:



Complete the mapping:



EXAMPLE: First order system

$$\frac{Y(s)}{x(s)} = \frac{1}{\tau s + 1}$$

$$(\tau s + 1)Y(s) = x(s)$$

$$sY(s) = \frac{1}{\tau}x(s) - \frac{1}{\tau}Y(s)$$

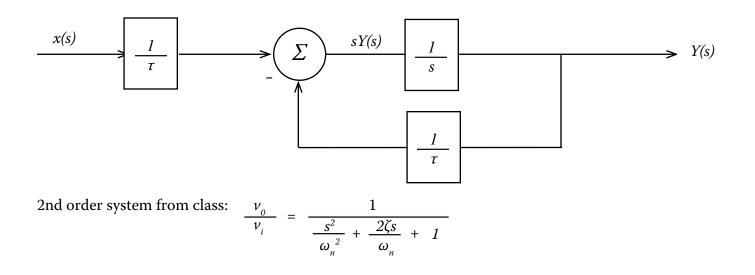
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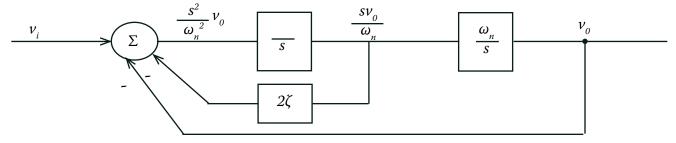


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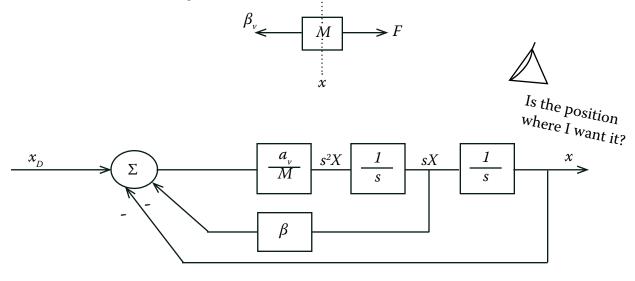
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With a little bit of manipulation, we can write the block diagram as



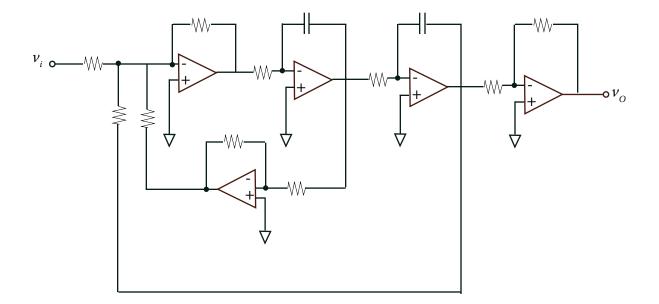
If we built this as an electronic circuit, it would be <u>analogous</u> to our mechanical system consisting of a mass, a viscous fluid, and a forcing mechanism:



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An analog computer might look something like this:

Make sure that ω_n is small compared to the parasitic poles of the op-amps. Then, we get a very good analog.