We've spent almost all of our time on linear system theory and its consequences. We are <u>lucky</u> that, although all systems in nature are nonlinear, linear system theory gets us so far.

What are we to do when forced to confront nonlinearity head on? Some of the options:

- $\bigcirc$  Linearize..and hope.  $\bigcirc$ 
  - 1.5) Linearize about a number of different points, and change our control strategy for each point.
- ② Use nonlinear analysis explicitly (nonlinear system theory is intensely formal. Insight, or general applicability, are hard to come by.)
- ③ Describing functions
- ④ Basic reasoning

In this recitation, we're going to talk about ③. But never forget ④: the concepts of control theory are bigger than any mathematical framework.

Before we go on, there's a concept that has crept into the class that observes careful treatment. With feedback systems, the word "stable" has a very definite meaning: all the poles of the system are in the left-half plane. But in other contexts, the word "stable" has a different meaning. A system is said to be in a stable "state" if the system tends to restore itself to that state in response to a perturbation. This is a general concept that shows up all over the place in engineering.

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EXAMPLE:







state: ball @ x = 0

perturbation: take my finger and move the ball

restorative force: remove disturbance, and system returns to x = 0

 $\Rightarrow$  this x = 0 is a stable equilibrium.

**OTHER EXAMPLES**:



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CLASS EXERCISE: Describe the stability of each of the following equilibria.



With describing function analysis, we make the following approximation, which I present here in three steps:



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Now, assume harmonics are small:

$$\underbrace{\underbrace{\text{E sin } \omega t}}_{\text{G}_{D}(\text{E},\omega)} \xrightarrow{} |G_{D}| \sin(\omega t + \measuredangle G_{D})$$

We only keep track of the fundamental.

Let's use this to look carefully at the describing function for a comparator:



In response to sine wave input Esin $\omega$ t (period, T =  $\frac{2\pi}{\omega}$ ), we get a square wave that looks as follows:



We put in a sine wave, and the nonlinear block returned that sine wave plus a bunch of harmonics. For describing function analysis, we only care about what happened to the fundamental.

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Using Fourier Analysis, we can extract the fundamental component from the output wave form.

$$\mathbf{v}_{0}(\mathbf{t}) = \mathbf{y}_{0}^{0} + \sum_{n=1}^{\infty} \mathbf{A}_{n} \sin n\omega \mathbf{t} + \sum_{n=1}^{\infty} \mathbf{B}_{n} \cos n\omega \mathbf{t}$$

( $a_0$  is zero by inspection: the output waveform clearly has no DC component.) We're intersted in  $A_1$  and  $B_1$ :

$$A_{1} = \frac{2}{T} \int_{0}^{T} v_{0}(t) \sin\omega t dt = \frac{2}{T} \int_{0}^{T/2} \sin\omega t dt - \frac{2}{T} \int_{T/2}^{T} \sin\omega t dt$$
$$= \frac{2}{T} \left[ -\frac{1}{\omega} \cos\omega t \right]_{0}^{T/2} - \frac{2}{T} \left[ -\frac{1}{\omega} \cos\omega t \right]_{T/2}^{T}$$
$$\vdots$$
$$= \frac{4}{\pi}$$

$$\mathbf{B}_{1} = \frac{2}{T} \int_{0}^{T} \mathbf{v}_{0}(t) \cos \omega t dt = \cdots = 0$$

Since  $B_1 = 0$ , there is no phase shift.

$$G_{D}(E) = \left(\frac{\sqrt{A_{1}^{2} + B_{1}^{2}}}{\text{input amplitude}}\right) = \frac{4}{\pi E}$$

Notice curious behavior: as the input amplitude gets larger, the gain gets smaller! This is useful in building constant-amplitude oscillators.  $\Rightarrow$ 

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Consider the 3-pole oscillator we looked at in lecture.



Ideally, we have two closed-loop poles sitting on  $j\omega$  axis:



We pick k to satisfy this condition. But what happens when we perturb the system and the amplitude gets bigger than intended? Does the system tend to restore itself?

Using our describing function,  $G_D(E) = \frac{4}{\pi E}$ , notice that as E gets bigger, the gain gets smaller. So in response to a perturbation that makes the amplitude bigger, the root locus changes as follows:

