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6.642 Continuum Electromechanics
Fall 2008

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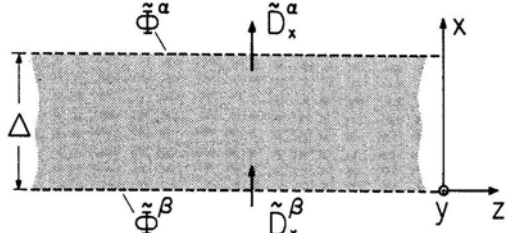
**Lecture 2: Continuum Electromechanics (Melcher) – Section 2.16
 Flux-Potential Relations for Laplacian Fields**

I. Planar Layers

A. Electric Fields

Flux-potential transfer relations for planar layer in terms of electric potential and normal displacement (Φ, D_x) . To obtain magnetic relations, substitute $(\Phi, D_x, \epsilon) \rightarrow (\Psi, B_x, \mu)$.

Planar layer



$$\begin{bmatrix} \tilde{D}_x^\alpha \\ \tilde{D}_x^\beta \end{bmatrix} = \epsilon \gamma \begin{bmatrix} -\coth(\gamma \Delta) & \frac{1}{\sinh(\gamma \Delta)} \\ \frac{-1}{\sinh(\gamma \Delta)} & \coth(\gamma \Delta) \end{bmatrix} \begin{bmatrix} \tilde{\Phi}^\alpha \\ \tilde{\Phi}^\beta \end{bmatrix} \quad (a)$$

$$\begin{bmatrix} \tilde{\Phi}^\alpha \\ \tilde{\Phi}^\beta \end{bmatrix} = \frac{1}{\epsilon \gamma} \begin{bmatrix} -\coth(\gamma \Delta) & \frac{1}{\sinh(\gamma \Delta)} \\ \frac{-1}{\sinh(\gamma \Delta)} & \coth(\gamma \Delta) \end{bmatrix} \begin{bmatrix} \tilde{D}_x^\alpha \\ \tilde{D}_x^\beta \end{bmatrix} \quad (b)$$

$$\Phi = \text{Re } \tilde{\Phi}(x, t) e^{-j(k_y y + k_z z)}$$

$$\gamma \equiv \sqrt{k_y^2 + k_z^2}$$

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$$\nabla \cdot \bar{\mathbf{E}} = 0$$

$$\nabla \times \bar{\mathbf{E}} = 0 \Rightarrow \bar{\mathbf{E}} = -\nabla \Phi$$

$$\nabla^2 \Phi = 0$$

$$\Phi(x, y, z, t) = \text{Re} \left[\tilde{\Phi}(x, t) e^{-j(k_y y + k_z z)} \right]$$

$$\nabla^2 \Phi = \frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0$$

$$\frac{d^2 \tilde{\Phi}}{dx^2} - (k_y^2 + k_z^2) \tilde{\Phi} = 0 \quad ; \quad \gamma = \sqrt{k_y^2 + k_z^2}$$

$$\tilde{\Phi} = \tilde{\Phi}_1 \sinh \gamma x + \tilde{\Phi}_2 \cosh \gamma x$$

$$\tilde{\Phi}(x = \Delta) = \tilde{\Phi}^\alpha \quad ; \quad \tilde{\Phi}(x = 0) = \tilde{\Phi}^\beta$$

$$\tilde{\Phi}(x) = \frac{\tilde{\Phi}^\alpha \sinh \gamma x - \tilde{\Phi}^\beta \sinh \gamma (x - \Delta)}{\sinh \gamma \Delta}$$

$$D_x = -\varepsilon \frac{\partial \Phi}{\partial x} \Rightarrow \tilde{D}_x = \frac{-\varepsilon \gamma \left[\tilde{\Phi}^\alpha \cosh \gamma x - \tilde{\Phi}^\beta \cosh \gamma (x - \Delta) \right]}{\sinh \gamma \Delta}$$

$$\tilde{D}_x^\alpha = \tilde{D}_x(x = \Delta) = \frac{-\varepsilon \gamma}{\sinh \gamma \Delta} \left[\tilde{\Phi}^\alpha \cosh \gamma \Delta - \tilde{\Phi}^\beta \right]$$

$$\tilde{D}_x^\beta = \tilde{D}_x(x = 0) = \frac{-\varepsilon \gamma}{\sinh \gamma \Delta} \left[\tilde{\Phi}^\alpha - \tilde{\Phi}^\beta \cosh \gamma \Delta \right]$$

B. Magnetic Fields

$$\nabla \times \vec{H} = 0 \Rightarrow \vec{H} = -\nabla \Psi$$

$$\nabla \cdot \vec{H} = 0 \Rightarrow \nabla^2 \Psi = 0 = \frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2}$$

$$\Psi(x, y, z, t) = \operatorname{Re} \left[\tilde{\Psi}(x, t) e^{-j(k_y y + k_z z)} \right]$$

$$\frac{d^2 \tilde{\Psi}}{dx^2} - \gamma^2 \tilde{\Psi} = 0 \quad ; \quad \gamma = \sqrt{k_y^2 + k_z^2}$$

$$\tilde{\Psi}(x) = \frac{\tilde{\Psi}^\alpha \sinh \gamma x - \tilde{\Psi}^\beta \sinh \gamma (x - \Delta)}{\sinh \gamma \Delta}$$

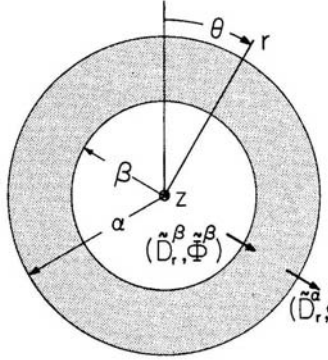
$$B_x = -\mu \frac{\partial \Psi}{\partial x} \Rightarrow \tilde{B}_x(x) = \frac{-\mu \gamma \left[\tilde{\Psi}^\alpha \cosh \gamma x - \tilde{\Psi}^\beta \cosh \gamma (x - \Delta) \right]}{\sinh \gamma \Delta}$$

$$\tilde{B}_x^\alpha = \tilde{B}_x(x = \Delta) = \frac{-\mu \gamma}{\sinh \gamma \Delta} \left[\tilde{\Psi}^\alpha \cosh \gamma \Delta - \tilde{\Psi}^\beta \right]$$

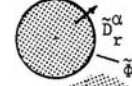
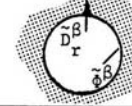
$$\tilde{B}_x^\beta = \tilde{B}_x(x = 0) = \frac{-\mu \gamma}{\sinh \gamma \Delta} \left[\tilde{\Psi}^\alpha - \tilde{\Psi}^\beta \cosh \gamma \Delta \right]$$

II. Cylindrical Annulus

Flux-potential relations for cylindrical annulus in terms of electric potential and normal displacement (ϕ, D_r) . To obtain magnetic relations, substitute $(\phi, D_r, \epsilon) \rightarrow (\psi, B_r, \mu)$.



$$\phi = \text{Re } \tilde{\phi}(r, t) e^{-j(m\theta + kz)}$$

$\begin{bmatrix} \tilde{D}_r^\alpha \\ \tilde{D}_r^\beta \end{bmatrix} = \epsilon \begin{bmatrix} f_m(\beta, \alpha) & g_m(\alpha, \beta) \\ g_m(\beta, \alpha) & f_m(\alpha, \beta) \end{bmatrix} \begin{bmatrix} \tilde{\phi}^\alpha \\ \tilde{\phi}^\beta \end{bmatrix} \quad (a)$ <p>$k = 0, m = 0$</p> $f_0(x, y) = \frac{1}{y} \ln\left(\frac{x}{y}\right); \quad g_0(x, y) = \frac{1}{x} \ln\left(\frac{x}{y}\right)$ <p>$k = 0, m = 1, 2, \dots$</p> $f_m(x, y) = \frac{m}{y} \frac{[\left(\frac{x}{y}\right)^m + \left(\frac{y}{x}\right)^m]}{[\left(\frac{x}{y}\right)^m - \left(\frac{y}{x}\right)^m]}$ $g_m(x, y) = \frac{2m}{x} \frac{1}{[\left(\frac{x}{y}\right)^m - \left(\frac{y}{x}\right)^m]}$ <p>$k \neq 0, m = 0, 1, 2, \dots^*$</p> $f_m(x, y) = \frac{jk[H_m(jkx)J'_m(jky) - J_m(jkx)H'_m(jky)]}{[J_m(jkx)H'_m(jky) - J_m(jky)H'_m(jkx)]}$ $g_m(x, y) = \frac{-2j}{\pi x [J_m(jkx)H'_m(jky) - J_m(jky)H'_m(jkx)]}$ $f_m(x, y) = \frac{k[K'_m(kx)I'_m(ky) - I'_m(kx)K'_m(ky)]}{[I'_m(kx)K'_m(ky) - I'_m(ky)K'_m(kx)]}$ $g_m(x, y) = \frac{1}{x[I'_m(kx)K'_m(ky) - I'_m(ky)K'_m(kx)]}$	$\begin{bmatrix} \tilde{\phi}^\alpha \\ \tilde{\phi}^\beta \end{bmatrix} = \frac{1}{\epsilon} \begin{bmatrix} F_m(\beta, \alpha) & G_m(\alpha, \beta) \\ G_m(\beta, \alpha) & F_m(\alpha, \beta) \end{bmatrix} \begin{bmatrix} \tilde{D}_r^\alpha \\ \tilde{D}_r^\beta \end{bmatrix} \quad (b)$ <p>$k = 0, m = 0$</p> <p>No inverse</p> <p>$k = 0, m = 1, 2, \dots$</p> $F_m(x, y) = \frac{y}{m} \frac{[\left(\frac{x}{y}\right)^m + \left(\frac{y}{x}\right)^m]}{[\left(\frac{x}{y}\right)^m - \left(\frac{y}{x}\right)^m]}$ $G_m(x, y) = \frac{2y}{m} \frac{1}{[\left(\frac{x}{y}\right)^m - \left(\frac{y}{x}\right)^m]}$ <p>$k \neq 0, m = 0, 1, 2, \dots^*$</p> $F_m(x, y) = \frac{1}{jk} \frac{[J'_m(jkx)H'_m(jky) - H'_m(jkx)J'_m(jky)]}{[J'_m(jky)H'_m(jkx) - J'_m(jkx)H'_m(jky)]}$ $G_m(x, y) = \frac{-2}{j\pi k(kx) [J'_m(jky)H'_m(jkx) - J'_m(jkx)H'_m(jky)]}$ $F_m(x, y) = \frac{1}{k} \frac{[I'_m(kx)K'_m(ky) - K'_m(kx)I'_m(ky)]}{[I'_m(ky)K'_m(kx) - I'_m(kx)K'_m(ky)]}$ $G_m(x, y) = \frac{1}{k(kx) [I'_m(ky)K'_m(kx) - I'_m(kx)K'_m(ky)]}$
<p>$\beta \rightarrow 0$</p>  $\tilde{D}_r^\alpha = \epsilon f_m(0, \alpha) \tilde{\phi}^\alpha; \quad f_m(0, \alpha) = -\frac{kI'_m(k\alpha)}{I_m(k\alpha)} \quad (c)$	<p>$\alpha \rightarrow \infty$</p>  $\tilde{D}_r^\beta = \epsilon f_m(\infty, \beta) \tilde{\phi}^\beta; \quad f_m(\infty, \beta) = -\frac{kK'_m(k\beta)}{K_m(k\beta)} \quad (d)$
<p>* See Prob. 2.17.2 for proof that $H_m(jkx)J'_m(jky) - J_m(jkx)H'_m(jky) = -2/(\pi kx)$ and $K_m(kx)I'_m(kx) - I'_m(kx)K'_m(kx) = 1/kx$ incorporated into g_m and G_m.</p>	

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$$\nabla^2 \Phi = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \Phi}{\partial \theta^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0$$

$$\Phi(r, \theta, z, t) = \operatorname{Re} \left[\tilde{\Phi}(r, t) e^{-j(m\theta + kz)} \right]$$

$$\frac{1}{r} \frac{d}{dr} \left(r \frac{d\tilde{\Phi}}{dr} \right) - \frac{m^2}{r^2} \tilde{\Phi} - k^2 \tilde{\Phi} = 0$$

1. $m=0, k=0$ solutions

$$\frac{1}{r} \frac{d}{dr} \left(r \frac{d\tilde{\Phi}}{dr} \right) = 0 \Rightarrow r \frac{d\tilde{\Phi}}{dr} = C \Rightarrow \tilde{\Phi} = C \ln r + D$$

$$\tilde{\Phi}(r) = \frac{\tilde{\Phi}^\alpha \ln \frac{r}{\beta} - \tilde{\Phi}^\beta \ln \frac{r}{\alpha}}{\ln \frac{\alpha}{\beta}}$$

$$D_r = -\varepsilon \frac{\partial \Phi}{\partial r} \Rightarrow \tilde{D}_r = -\varepsilon \frac{\tilde{\Phi}^\alpha - \tilde{\Phi}^\beta}{r \ln \frac{\alpha}{\beta}}$$

$$\tilde{D}_r^\alpha = -\varepsilon \frac{\tilde{\Phi}^\alpha - \tilde{\Phi}^\beta}{\alpha \ln \frac{\alpha}{\beta}}$$

$$\tilde{D}_r^\beta = -\varepsilon \frac{\tilde{\Phi}^\alpha - \tilde{\Phi}^\beta}{\beta \ln \frac{\alpha}{\beta}}$$

$$\begin{bmatrix} \tilde{D}_r^\alpha \\ \tilde{D}_r^\beta \end{bmatrix} = \frac{-\varepsilon}{\ln \frac{\alpha}{\beta}} \begin{bmatrix} \frac{1}{\alpha} & -\frac{1}{\alpha} \\ \frac{1}{\beta} & -\frac{1}{\beta} \end{bmatrix} \begin{bmatrix} \tilde{\Phi}^\alpha \\ \tilde{\Phi}^\beta \end{bmatrix}$$

2. $k=0, m \neq 0$

$$r \frac{d}{dr} \left(r \frac{d\tilde{\Phi}}{dr} \right) - m^2 \tilde{\Phi} = 0 \Rightarrow \tilde{\Phi} = A_1 r^m + A_2 r^{-m}$$

$$\tilde{\Phi} = \tilde{\Phi}^\alpha \frac{\left[\left(\frac{\beta}{r} \right)^m - \left(\frac{r}{\beta} \right)^m \right]}{\left[\left(\frac{\beta}{\alpha} \right)^m - \left(\frac{\alpha}{\beta} \right)^m \right]} + \tilde{\Phi}^\beta \frac{\left[\left(\frac{r}{\alpha} \right)^m - \left(\frac{\alpha}{r} \right)^m \right]}{\left[\left(\frac{\beta}{\alpha} \right)^m - \left(\frac{\alpha}{\beta} \right)^m \right]}$$

3. $k \neq 0, m \neq 0$

$$\tilde{\Phi} = A_1 I_m(kr) + A_2 K_m(kr) \text{ [Modified Bessel Functions]}$$

$$I_m(r \rightarrow \infty) = \infty, I_m(r \rightarrow 0) \rightarrow \text{finite}$$

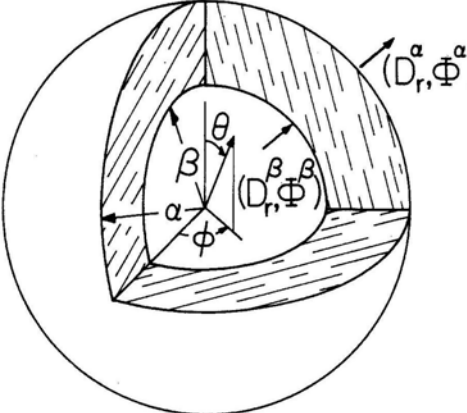
$$K_m(r \rightarrow \infty) \rightarrow 0, K_m(r \rightarrow 0) \rightarrow \infty$$

$$J_m(jkr) = j^m I_m(kr), H_m(jkr) = \frac{2}{\pi} j^{-(m+1)} K_m(kr)$$

$$\tilde{\Phi}(r) = \frac{\left\{ \tilde{\Phi}^\alpha [H_m(jk\beta)J_m(jkr) - J_m(jk\beta)H_m(jkr)] + \tilde{\Phi}^\beta [J_m(jk\alpha)H_m(jkr) - H_m(jk\alpha)J_m(jkr)] \right\}}{[H_m(jk\beta)J_m(jk\alpha) - J_m(jk\beta)H_m(jk\alpha)]}$$

III. Spherical Shell

Flux-potential transfer relations for spherical shell in terms of electric potential and normal displacement (Φ, D_r). To obtain magnetic relations, substitute $(\Phi, D_r, \epsilon) \rightarrow (\Psi, B_r, \mu)$.



$$\Phi = \text{Re } \tilde{\Phi}(r, t) P_n^m(\cos \theta) e^{-jm\phi}$$

$$P_n^m = (1-x^2)^{m/2} \frac{d^m P_n}{dx^m}$$

$$P_0 = 1, P_1 = x, P_2 = \frac{1}{2}(3x^2 - 1)$$

$$P_3 = \frac{1}{2}(5x^3 - 3x)$$

$$P_4 = \frac{1}{8}(35x^4 - 30x^2 + 3)$$

m	P_0^m	P_1^m	$P_1^m \cos m\phi$	P_2^m	$P_2^m \cos m\phi$	P_3^m	$P_3^m \cos m\phi$
0	1	$\cos \theta$	$\begin{matrix} + \\ - \end{matrix}$	$\frac{1}{2}(3 \cos^2 \theta - 1)$	$\begin{matrix} + \\ - \\ + \end{matrix}$	$\frac{1}{2}(5 \cos^3 \theta - 3 \cos \theta)$	$\begin{matrix} + \\ - \\ + \\ - \end{matrix}$
1	0	$\sin \theta$	$\begin{matrix} + & - & + \end{matrix}$	$3 \sin \theta \cos \theta$	$\begin{matrix} + & - & + \\ - & + & - \end{matrix}$	$\frac{3}{2} \sin \theta (5 \cos^2 \theta - 1)$	$\begin{matrix} + & - & + \\ - & + & - \\ + & - & + \end{matrix}$
2	0	0		$3 \sin^2 \theta$	$\begin{matrix} + & - & + & - \\ - & + & - & + \end{matrix}$	$15 \sin^2 \theta \cos \theta$	$\begin{matrix} + & - & + & - & + \\ - & + & - & + & - \end{matrix}$
3	0	0		0		$15 \sin^3 \theta$	$\begin{matrix} + & - & + & - & + & - \\ - & + & - & + & - & + \end{matrix}$

$$\begin{bmatrix} \tilde{D}_r^\alpha \\ \tilde{D}_r^\beta \end{bmatrix} = \epsilon \begin{bmatrix} f_n(\beta, \alpha) & g_n(\alpha, \beta) \\ g_n(\beta, \alpha) & f_n(\alpha, \beta) \end{bmatrix} \begin{bmatrix} \tilde{\Phi}^\alpha \\ \tilde{\Phi}^\beta \end{bmatrix} \quad (a)$$

$$f_n(x, y) = \frac{[n(\frac{y}{x})^n + (n+1)(\frac{x}{y})^{n+1}]}{[x(\frac{x}{y})^n - y(\frac{y}{x})^n]}$$

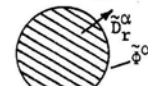
$$g_n(x, y) = \frac{(2n+1)}{x^2[\frac{1}{y}(\frac{x}{y})^n - \frac{1}{x}(\frac{y}{x})^n]}$$

$$\begin{bmatrix} \tilde{\Phi}^\alpha \\ \tilde{\Phi}^\beta \end{bmatrix} = \frac{1}{\epsilon} \begin{bmatrix} F_n(\beta, \alpha) & G_n(\alpha, \beta) \\ G_n(\beta, \alpha) & F_n(\alpha, \beta) \end{bmatrix} \begin{bmatrix} \tilde{D}_r^\alpha \\ \tilde{D}_r^\beta \end{bmatrix} \quad (b)$$

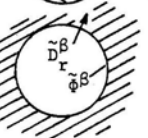
$$F_n(x, y) = \frac{y}{x} \frac{[\frac{1}{n}(\frac{y}{x})^n + \frac{1}{n+1}(\frac{x}{y})^{n+1}]}{[\frac{1}{y}(\frac{x}{y})^n - \frac{1}{x}(\frac{y}{x})^n]}$$

$$G_n(x, y) = \frac{y}{x} \frac{(2n+1)}{n(n+1)} \frac{1}{[\frac{1}{y}(\frac{x}{y})^n - \frac{1}{x}(\frac{y}{x})^n]}$$

$\beta \rightarrow 0$



$\alpha \rightarrow \infty$



$$\tilde{D}_r^\alpha = -\frac{\epsilon n}{\alpha} \tilde{\Phi}^\alpha \quad (c)$$

$$\tilde{D}_r^\beta = \frac{\epsilon(n+1)}{\beta} \tilde{\Phi}^\beta \quad (d)$$

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$$\Phi(r, \theta, \phi) = \text{Re} \left[\tilde{\Phi}(r) \Theta(\theta) e^{-jm\phi} \right]$$

$$\nabla^2 \Phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \Phi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Phi}{\partial \phi^2} = 0$$

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left[\sin \theta \frac{\partial \Theta}{\partial \theta} \right] - \frac{m^2}{\sin^2 \theta} = -k^2$$

$$\frac{1}{\tilde{\Phi}} \frac{d}{dr} \left(r^2 \frac{d\tilde{\Phi}}{dr} \right) = k^2$$

$$u = \cos \theta, \sqrt{1 - u^2} = \sin \theta$$

$$(1 - u^2) \frac{d^2 \Theta}{du^2} - 2u \frac{d\Theta}{du} + \left(k^2 - \frac{m^2}{1 - u^2} \right) \Theta = 0$$

$$k^2 = n(n + 1), n \text{ an integer}$$

$$\Theta = P_n^m(u)$$

$$\tilde{\Phi} = \tilde{\Phi}^\alpha \frac{\left[\left(\frac{r}{\beta} \right)^n - \left(\frac{\beta}{r} \right)^{n+1} \right]}{\left[\left(\frac{\alpha}{\beta} \right)^n - \left(\frac{\beta}{\alpha} \right)^{n+1} \right]} + \tilde{\Phi}^\beta \frac{\left[\left(\frac{r}{\alpha} \right)^n - \left(\frac{\alpha}{r} \right)^{n+1} \right]}{\left[\left(\frac{\beta}{\alpha} \right)^n - \left(\frac{\alpha}{\beta} \right)^{n+1} \right]}$$

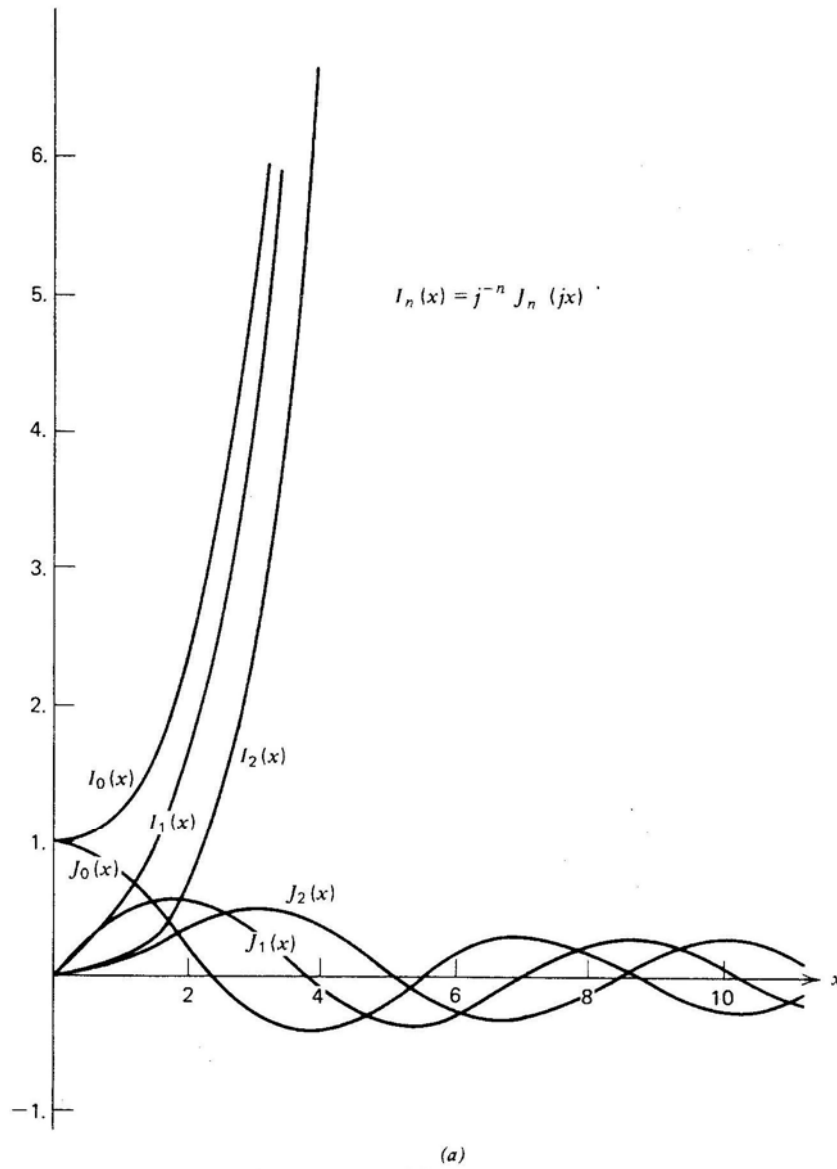
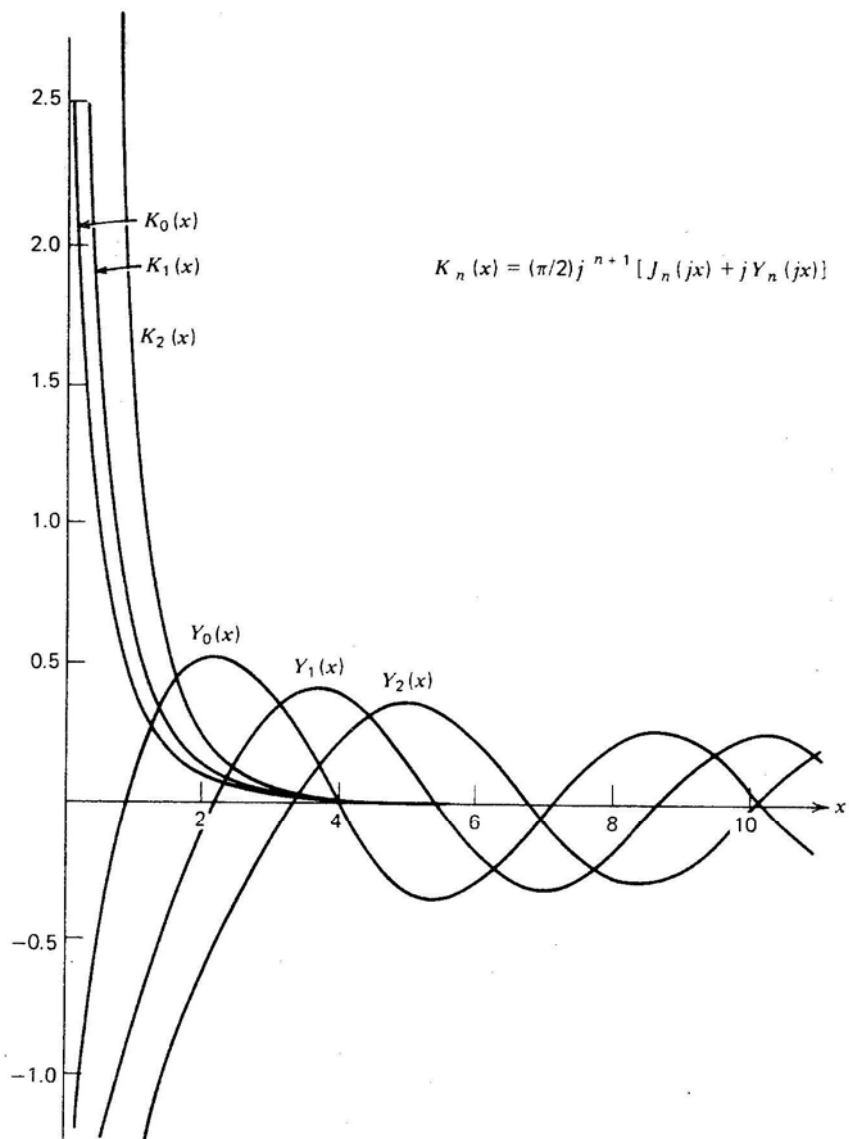


Figure 4-9a The Bessel functions (a) $J_n(x)$ and $I_n(x)$, and (b) $Y_n(x)$ and $K_n(x)$.

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(b)
Figure 4-9b

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