## Massachusetts Institute of Technology Department of Physics Physics 8.022 – Fall 2003 Useful Formulae

<u>Conservative Field:</u>  $\int_C \vec{F} \cdot d\vec{r} = 0$  for any closed path  $C, W_{ab,C} = W_{ab,C'}$  for any C, C' connecting a and  $b, \vec{F} = -\vec{\nabla}U, \vec{\nabla} \times \vec{F} = 0$ 

<u>Coulomb Law:</u>  $\vec{F}_{21} = \frac{q_1 q_2}{r^2} \hat{r}_{21}$  for two point charges at distance r.  $\vec{F}_{12} = -\vec{F}_{21}$ , and for charges  $dq_1$  and  $dq_2$  that make part of continuous charge distributions 1 and 2,  $d\vec{F}_{21} = \frac{dq_1 dq_2}{r^2} \hat{r}_{21}$ 

<u>Electric Field</u>: at point 2 due to  $q_1 \vec{E}_1 = \frac{q_1}{r^2} \hat{r}_{21}$ . If  $q_1$  is not a point charge but part of a continuous distribution,  $d\vec{E} = \frac{dq}{r^2} \hat{r}$ 

<u>Principle of Superposition</u>: Two or more electric fields acting at a given point P add vectorially:  $\vec{E}_P = \vec{E}_1 + \vec{E}_2 + \ldots + \vec{E}_n$ 

<u>Electrostatic Field is Conservative</u>:  $\vec{\nabla} \times \vec{E} = 0$  and thus there exists scalar function  $\phi$  such that  $\vec{E} = -\vec{\nabla}\phi$  where  $d\phi = \frac{dq}{r}$ 

Electrostatic potential: The potential at  $\vec{x}$  with respect to a *ref* point is  $\phi(\vec{x}) - \phi(ref) = -\int_{ref}^{\vec{x}} \vec{E} \cdot d\vec{r} = -\frac{W_{ref \to \vec{x}}}{q}$ 

<u>Gauss Law:</u>  $\int_S \vec{E} \cdot d\vec{a} = 4\pi \int_V \rho dV$  where S is a closed surface and V is its corresponding volume (integral form) or  $\vec{\nabla} \cdot \vec{E} = 4\pi\rho$  (differential form).

 $\begin{array}{l} \underline{\operatorname{Poisson Eqn:}} \nabla^2 \phi = -4\pi\rho, \, \underline{\operatorname{Laplace Eqn:}} \nabla^2 \phi = 0 \\ \underline{\operatorname{Energy:}} \, U = \frac{1}{2} \int_V dV \int_{V'} dV' \frac{\rho(\vec{x})\rho(\vec{x}')}{|\vec{x}-\vec{x}'|} = \frac{1}{2} \int_V \rho \phi dV = \frac{1}{8\pi} \int_V E^2 dV \\ \underline{\operatorname{Electric Force on Conductors:}} \, \frac{dF}{da} = 2\pi\sigma^2 = \frac{E^2}{8\pi} \\ \underline{\operatorname{Current Density:}} \, \vec{J}(\vec{x}) = \rho(\vec{x})\vec{v}(\vec{x}), \, \underline{\operatorname{Conservation Law/Continuity:}} \, \vec{\nabla} \cdot \vec{J} = -\frac{\partial\rho}{\partial t} \\ \underline{\operatorname{Currents:}} \, I = \frac{dq}{dt}, I = \int_S \vec{J} \cdot d\vec{a} \\ \underline{\operatorname{Capacitance:}} \, Q = CV, \, (\operatorname{energy}) \, U = \frac{1}{2}CV^2 \\ \underline{\operatorname{Capacitor networks:}} \, \operatorname{Parallel:} \, C = C_1 + C_2, \, \operatorname{series:} \, \frac{1}{C} = \frac{1}{C_1} + \frac{1}{C_2} \\ \underline{\operatorname{Ohm's Law:}} \, \vec{J} = \sigma \vec{E}, V = IR \\ \underline{\operatorname{Resistor networks:}} \, \operatorname{Series:} \, R = R_1 + R_2, \, \operatorname{parallel:} \, \frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2} \\ \underline{\operatorname{Magnetic charges:}} \, \vec{\nabla} \cdot \vec{B} = 0 \\ \underline{\operatorname{Biot-Savart's Law:}} \, d\vec{B} = \frac{Id\vec{l} \times \vec{r}}{cr^2} \\ \underline{\operatorname{Lorentz Force:}} \, \vec{F} = q\vec{E} + \frac{q}{c}\vec{v} \times \vec{B} \\ \underline{\operatorname{Force on current:}} \, \vec{F} = \frac{1}{c}d\vec{l} \times \vec{B} \end{array}$ 

Ampere's Law: 
$$\oint_C \vec{B} \cdot d\vec{l} = \frac{4\pi}{c} I_{encl} = \frac{4\pi}{c} \int_S \vec{J} \cdot d\vec{a}, \ \vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{J}$$

**Relativistic Transformations:** 

All primed quantities measured in the frame F' which is moving in the positive x direction with velocity  $u = \beta c$  (and  $\gamma = 1/\sqrt{1-\beta^2}$ ) as seen from F. Position and momentum 4-vectors:

$$x' = \gamma(x - \beta ct) \quad p' = \gamma(p - \beta \frac{E}{c})$$
  
$$t' = \gamma(t - \beta \frac{x}{c}) \quad E' = \gamma(E - \beta cp)$$

Electric and magnetic field components:

$$E'_{x} = E_{x} \quad E'_{y} = \gamma(E_{y} - \beta B_{z}) \quad E'_{z} = \gamma(E_{z} + \beta B_{y})$$
$$B'_{x} = B_{x} \quad B'_{y} = \gamma(B_{y} + \beta E_{z}) \quad B'_{z} = \gamma(B_{z} - \beta E_{y})$$

Relativistic Mass, Energy:  $m = \gamma m_0, E = mc^2$ <u>Faraday's Law:</u>  $\mathcal{E} = \oint_C \vec{E} \cdot d\vec{l} = -\frac{1}{c} \frac{d\Phi}{dt}, \vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$ <u>Mutual Inductance</u>:  $M_{12} = M_{21} = \frac{\Phi_{21}}{cI_1}, \mathcal{E}_{21} = -M_{21}\frac{dI_1}{dt}$ <u>Self Inductance</u>:  $L = \frac{\Phi}{cI}, \mathcal{E} = -L\frac{dI}{dt}$ Magnetic Field Energy Density:  $\frac{dU_B}{dv} = u_B = \frac{B^2}{8\pi}$ Impedance:  $V = ZI, Z_R = R, Z_C = -\frac{i}{\omega C}, Z_L = i\omega L$ <u>Admittance:</u> Y = 1/ZComplex notation:  $e^{i\theta} = \cos\theta + i\sin\theta$ ,  $z = \alpha + ib = |z|e^{i\theta}$ ,  $|z| = \sqrt{\alpha^2 + b^2}$ ,  $\tan\theta = b/\alpha$ Displacement: current  $I_d = \frac{1}{4\pi} \frac{d\Phi_E}{dt}$ , density  $\vec{J}_D = \frac{1}{4\pi} \frac{\partial \vec{E}}{\partial t}$ Maxwell's equations:  $\vec{\nabla} \cdot \vec{E} = 4\pi\rho, \vec{\nabla} \cdot \vec{B} = 0, \vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}, \vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$  $\overline{\omega = 2\pi/T, c = \omega/k, \vec{E} \times \vec{B}}$  points along propagation direction  $(\hat{k})$ <u>Wave "mechanics"</u>: Poynting vector  $\vec{S} = \frac{c}{4\pi}\vec{E}\times\vec{B}$ , momentum density  $\vec{g} = \frac{1}{4\pi c}\vec{E}\times\vec{B}$ , energy density  $u_T = \frac{E^2}{8\pi} + \frac{B^2}{8\pi}$ , radiation pressure  $\frac{F}{A} = u_T$ Waveguides:  $\frac{\partial I}{\partial r} = -C_0 \frac{\partial V}{\partial t}, \frac{\partial V}{\partial r} = -L_0 \frac{\partial I}{\partial t}, \frac{\partial^2 I}{\partial r^2} - L_0 C_0 \frac{\partial^2 I}{\partial t^2} = 0, \frac{\partial^2 V}{\partial r^2} - L_0 C_0 \frac{\partial^2 V}{\partial t^2} = 0$  $\begin{array}{l} \underline{\text{Gradient:}} \text{ in cartesian } \vec{\nabla}f = \frac{\partial f}{\partial x} \hat{x} + \frac{\partial f}{\partial y} \hat{y} + \frac{\partial f}{\partial z} \hat{z} \text{ , in cylindrical } \vec{\nabla}f = \frac{\partial f}{\partial \rho} \hat{\rho} + \frac{1}{\rho} \frac{\partial f}{\partial \phi} \hat{\phi} + \frac{\partial f}{\partial z} \hat{z} \\ \text{, in spherical } \vec{\nabla}f = \frac{\partial f}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial f}{\partial \theta} \hat{\theta} + \frac{1}{rsin\theta} \frac{\partial f}{\partial \phi} \hat{\phi} \end{array}$