

7. Entropy as a Thermodynamic Variable

$$\left(\frac{\partial S}{\partial E}\right)_{dW=0} \equiv \frac{1}{T} \quad \text{gives us } T$$

Other derivatives give other thermodynamic variables.

$$dW = \left\{ \begin{array}{l} -PdV \\ SdA \\ \mathcal{F}dL \end{array} \right\} + HdM + \mathcal{E}d\mathcal{P} + \dots \equiv \sum_i X_i dx_i$$

We chose to use the extensive external variables (a complete set) as the constraints on Ω . Thus

$$S \equiv k \ln \Omega = S(E, V, M, \dots)$$

Now solve for E .

$$S(E, V, M, \dots) \leftrightarrow E(S, V, M, \dots)$$

We know

$$dE|_{dW=0} = dQ \quad \text{from the 1}^{ST} \text{ law}$$

$$dE|_{dW=0} \leq TdS \quad \text{utilizing the 2}^{ND} \text{ law}$$

Now include the work.

$$dE = \delta Q + \delta W$$

$$dE \leq TdS + \delta W$$

$$dE \leq TdS + \left\{ \begin{array}{c} -PdV \\ SdA \\ \mathcal{F}dL \end{array} \right\} + HdM + \mathcal{E}d\mathcal{P} + \dots$$

The last line expresses the combined
1ST and 2ND laws of thermodynamics.

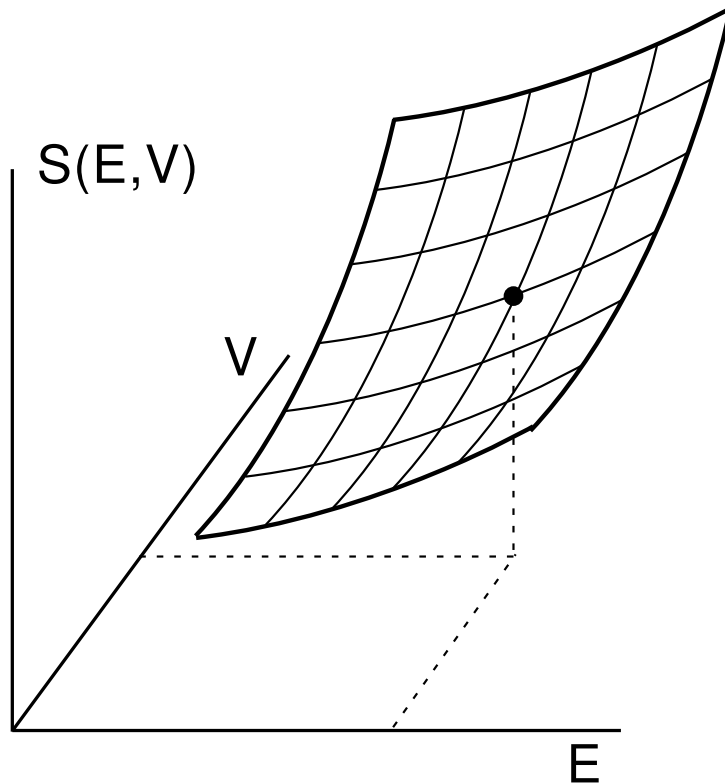
Solve for dS .

$$dS = \frac{1}{T}dE + \frac{P}{T}dV - \frac{H}{T}dM - \frac{\mathcal{E}}{T}d\mathcal{P} + \dots$$

Examine the partial derivatives of S .

$$\begin{aligned} \left(\frac{\partial S}{\partial E}\right)_{V,M,\mathcal{P}} &= \frac{1}{T} & \left(\frac{\partial S}{\partial M}\right)_{E,V,\mathcal{P}} &= -\frac{H}{T} \\ \left(\frac{\partial S}{\partial V}\right)_{E,M,\mathcal{P}} &= \frac{P}{T} & \left(\frac{\partial S}{\partial x_j}\right)_{E,x_i \neq x_j} &= -\frac{X_j}{T} \end{aligned}$$

INTERPRETATION



$$\begin{aligned} dS &= \left(\frac{\partial S}{\partial E} \right)_V dE + \left(\frac{\partial S}{\partial V} \right)_E dV \\ &= \frac{1}{T} dE + \frac{P}{T} dV \end{aligned}$$

UTILITY

Internal Energy

$$\left(\frac{\partial S(E, V, N)}{\partial E}\right)_V = \frac{1}{T} \rightarrow T(E, V, N) \leftrightarrow E(T, V, N)$$

Equation of State

$$\left(\frac{\partial S(E, V, N)}{\partial V}\right)_E = \frac{P}{T} \rightarrow P(E, T, V, N) \rightarrow P(T, V, N)$$

Example Ideal Gas

$$S(E, N, V) = k \ln \Phi = kN \ln \left\{ V \left(\frac{4}{3} \pi e m \left(\frac{E}{N} \right) \right)^{3/2} \right\}$$

$$\left(\frac{\partial S}{\partial V} \right)_{E, N} = \frac{kN}{\{ \}} \frac{\{ \}}{V} = \frac{kN}{V} = \frac{P}{T}$$

$$\underline{PV = NkT}$$

COMBINATORIAL FACTS

different orderings (permutations) of K distinguishable objects = $K!$

of ways of choosing L from a set of K :

$$\frac{K!}{(K-L)!} \quad \text{if order matters}$$

$$\frac{K!}{L!(K-L)!} \quad \text{if order does not matter}$$

EXAMPLE Dinner Table, 5 Chairs (places)

Seating, 5 people

$$5 \cdot 4 \cdot 3 \cdot 2 \cdot 1 = 5! = 120$$

Seating, 3 people

$$5 \cdot 4 \cdot 3 = \frac{5!}{2!} = 60$$

Place settings, 3 people

$$5 \cdot 4 \cdot 3 / 6 = \frac{5!}{2!} \frac{1}{3!} = 10$$

EXAMPLE 2 Level System

Ensemble of N "independent" systems

ENERGY ↑

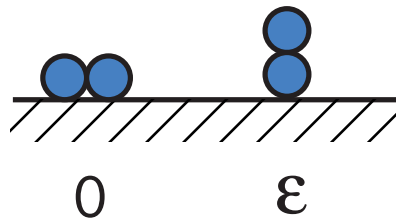
$|1\rangle$ ϵ

$|0\rangle$ 0

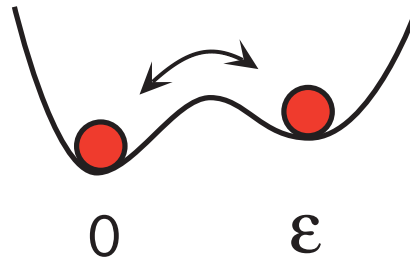
$$N = N_0 + N_1$$

$$E = \epsilon N_1$$

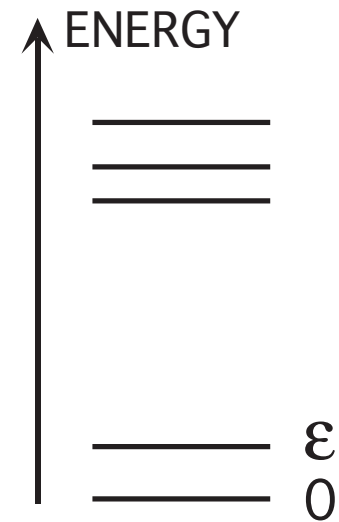
SURFACE MOLECULES



IONS IN A CRYSTAL



LOWEST LYING STATES



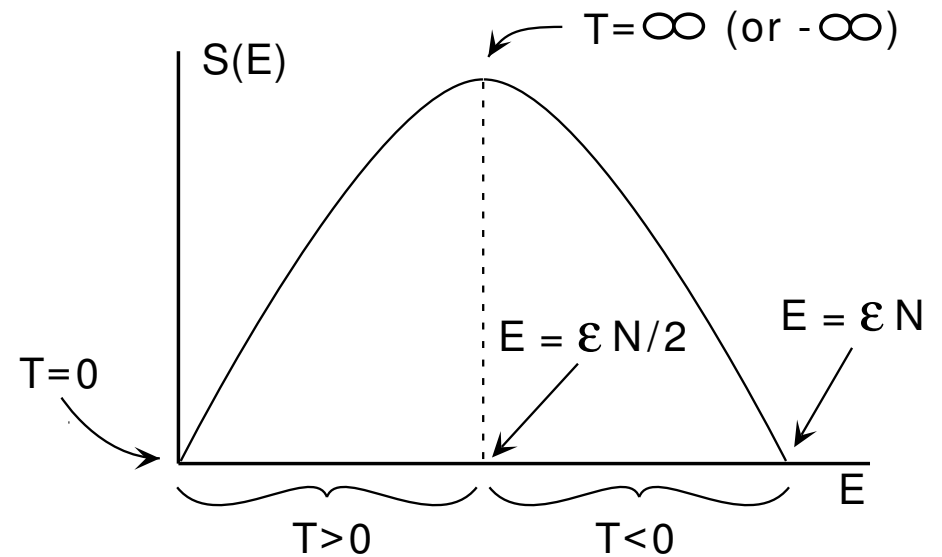
- $E \leftrightarrow N_1$
- NO WORK POSSIBLE (JUST HEAT FLOW)

$$\Omega(E) = \frac{N!}{N_1!(N-N_1)!}$$

1 when $N_1 = 0$ or N

Maximum when $N_1 = N/2$

$$S(E) = k \ln \Omega(E)$$



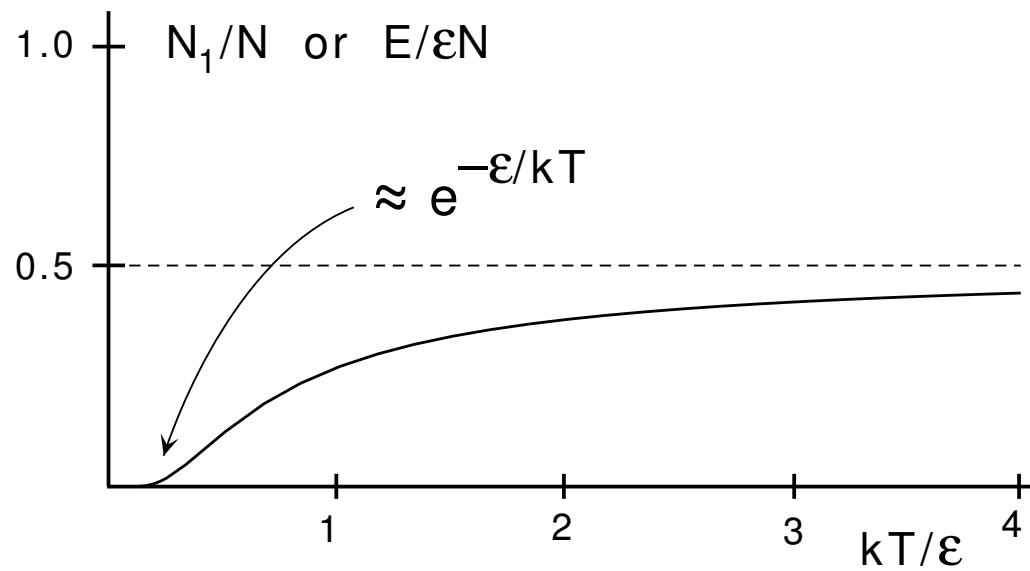
$$\ln N! \approx N \ln N - N$$

$$S(E) = k[N \ln N - N_1 \ln N_1 - (N - N_1) \ln(N - N_1) - N + N_1 + N - N_1]$$

$$\begin{aligned} \frac{1}{T} &= \left(\frac{\partial S}{\partial E} \right)_N = \frac{\partial S}{\partial N_1} \underbrace{\frac{\partial N_1}{\partial E}}_{1/\epsilon} = \frac{k}{\epsilon} [-1 - \ln N_1 + 1 + \ln(N - N_1)] \\ &= \frac{k}{\epsilon} \ln \left(\frac{N - N_1}{N_1} \right) = \frac{k}{\epsilon} \ln \left(\frac{N}{N_1} - 1 \right) \end{aligned}$$

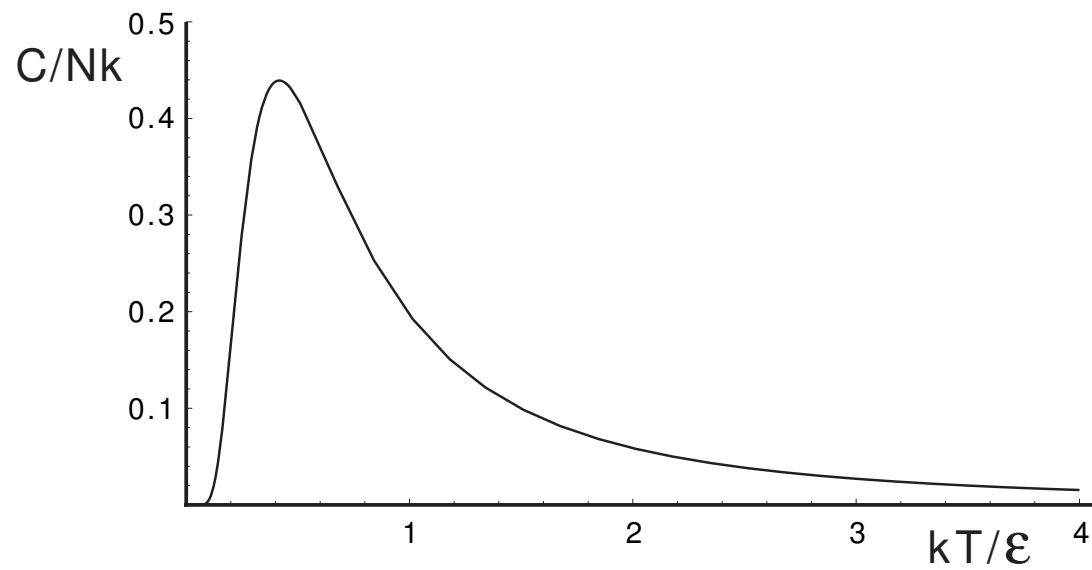
$$\frac{N}{N_1} - 1 = e^{\epsilon/kT} \rightarrow N_1 = \frac{N}{e^{\epsilon/kT} + 1}$$

$$E = \epsilon N_1 = \frac{\epsilon N}{e^{\epsilon/kT} + 1}$$



$$C \equiv \frac{\partial E}{\partial T} = Nk \left(\frac{\epsilon}{kT} \right)^2 \frac{e^{\epsilon/kT}}{(e^{\epsilon/kT} + 1)^2}$$

$$\rightarrow Nk \left(\frac{\epsilon}{kT} \right)^2 e^{-\epsilon/kT} \text{ low } T, \quad \rightarrow \frac{Nk}{4} \left(\frac{\epsilon}{kT} \right)^2 \text{ high } T$$



$$p(n) = ? \quad n = 0, 1 \qquad p(n) = \frac{\Omega'}{\Omega}$$

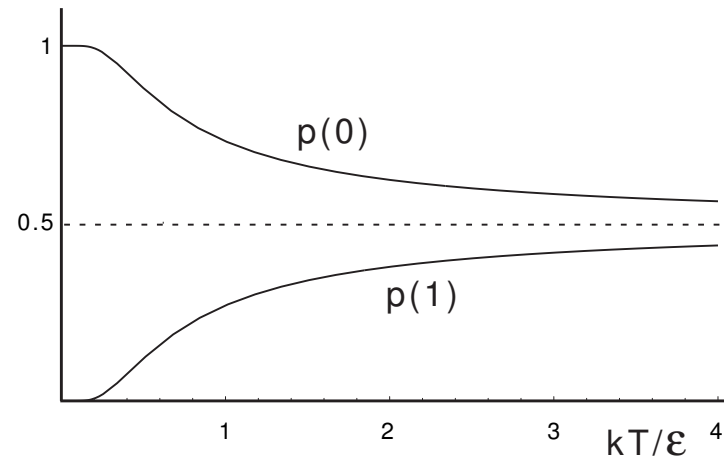
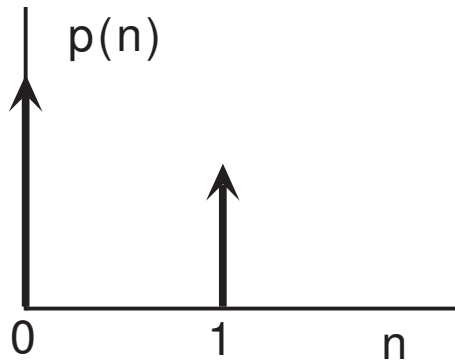
$$\text{In } \Omega' \quad N \rightarrow N - 1 \quad \text{and} \quad N_1 \rightarrow N_1 - n$$

$$p(n) = \frac{\frac{(N-1)!}{(N_1-n)!(N-1-N_1+n)!}}{\frac{N!}{N_1!(N-N_1)!}}$$

$$p(n) = \frac{(N-1)!}{N!} \frac{N_1!}{(N_1-n)!} \frac{(N-N_1)!}{(N-N_1-1+n)!}$$

$1/N$	$1 \quad n=0$	$N-N_1 \quad n=0$
	$N_1 \quad n=1$	$1 \quad n=1$

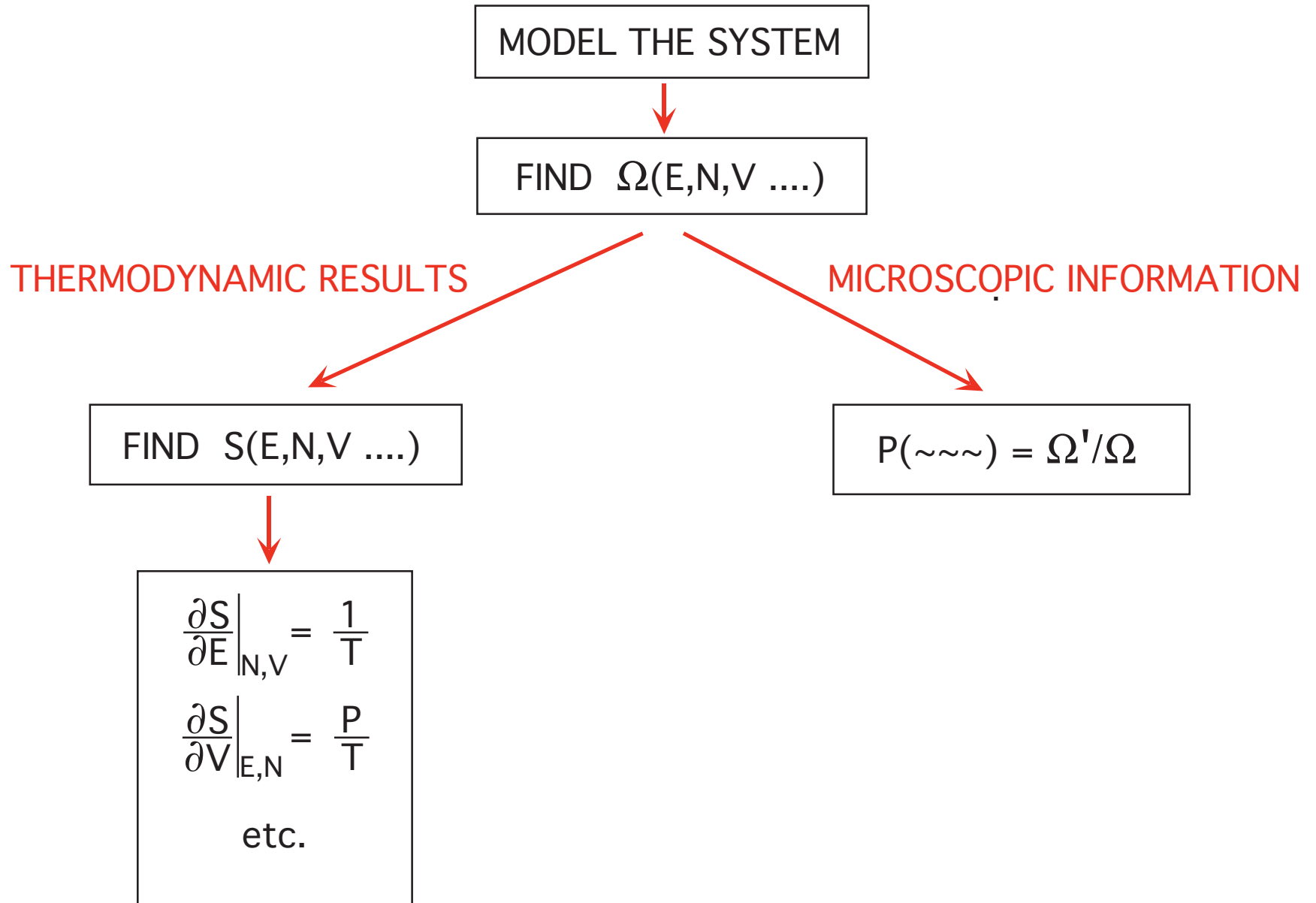
$$\left. \begin{aligned} p(0) &= \frac{N-N_1}{N} = 1 - \frac{N_1}{N} \\ p(1) &= \frac{N_1}{N} = [e^{\epsilon/kT} + 1]^{-1} \end{aligned} \right\} p(0) + p(1) = 1$$



$$E = (0)N p(0) + (\epsilon)N p(1) = \frac{\epsilon N}{e^{\epsilon/kT} + 1}$$

But we knew E , so we could have worked backwards to find $p(1)$.

MICROCANONICAL ENSEMBLE



The microcanonical ensemble is the starting point for Statistical Mechanics.

- We will no longer use it to solve problems.
- We will develop our understanding of the 2^{ND} law.
- We will derive the canonical ensemble, the real workhorse of S.M.

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