8.05 Quantum Physics II, Fall 2012 FINAL EXAM Tuesday December 18, 1:30 pm -4:30 pm You have 3 hours.

Answer all problems in the white books provided. Write YOUR NAME and YOUR SECTION on your white book(s).

THIS TEST <u>MUST BE RETURNED</u> WITH YOUR WHITE BOOK(S)

YOUR NAME:

The formula sheet need not be returned

There are six questions, totalling 100 points.

None of the problems requires extensive algebra.

No books, notes, or calculators allowed.

TIME MANAGEMENT:

One hour to answer questions 1,2, and 3. 30 minutes for problem 4.

45 minutes for each problem 5 and 6. The last two questions in problem 5 and the last in problem 6 may be challenging

- 1. **True or false questions** [20 points] No explanations required. Just indicate T or F for true or false, respectively.
 - (1) The wavefunction of a particle in a three-dimensional central potential must vanish at the origin.
 - (2) The energy spectrum of a particle on a three-dimensional spherical infinite well $(V = 0 \text{ for } r \leq a, \text{ and } V = \infty \text{ otherwise})$ has a number of degeneracies that are not explained by angular momentum conservation.
 - (3) $[H, L_z] = 0$ and $[H, L_x] = 0$ guarantee the full rotational invariance of the Hamiltonian.
 - (4) When the spin angular momentum takes a half-integer value, the orbital angular momentum must also take a half-integer value in order that the total angular momentum takes an integer value.
 - (5) $\mathbf{r} \times \mathbf{p} = -\mathbf{p} \times \mathbf{r}$.
 - (6) $\mathbf{r} \times \mathbf{L} = -\mathbf{L} \times \mathbf{r}$.
 - (7) If two Schrödinger operators commute, the corresponding Heisenberg operators may fail to commute if the Schrödinger operators have explicit time dependence.
 - (8) The Heisenberg Hamiltonian is given by the Schrödinger Hamiltonian with each operator replaced by its Heisenberg version.
 - (9) For a simple harmonic oscillator the Heisenberg number operator is identical to the Schrödinger number operator.
 - (10) The knowledge of the time evolution operator $\mathcal{U}(t, t_0)$ does not suffice to construct the Hamiltonian of the system.
- 2. Two short problems [10 points]
 - (a) A spin-1/2 particle with $\vec{\mu} = \gamma \vec{S}$ is in the up state along the z axis at t = 0. Give a time-independent magnetic field \vec{B} that will make the spin trace a cone in which the spin points successively along the x, y, and z axis, in a motion with period T.
 - (b) Consider the state built by combining two spin-1/2 particles as follows

$$\frac{1}{\sqrt{2}} \left(|\vec{n};+\rangle |\vec{n};-\rangle - |\vec{n};-\rangle |\vec{n};+\rangle \right).$$

Demonstrate that this is a singlet of the total angular momentum for any choice of the direction \vec{n} by writing this state in terms of $|+\rangle$ and $|-\rangle$ states.

3. Runge Lenz vector [10 points]

- (a) [4 pt] Is it true that $\mathbf{r} \cdot \mathbf{L} = 0$ and $\mathbf{p} \cdot \mathbf{L} = 0$? Explain your answer. (Their analogs are true in classical mechanics!).
- (b) [6 pt] The quantum Runge-Lenz vector can be written as

$$\mathbf{R} = \frac{1}{me^2} \left(\mathbf{p} \times \mathbf{L} - i\hbar \mathbf{p} \right) - \frac{\mathbf{r}}{r}.$$

Prove that

$$\mathbf{R} \cdot \mathbf{L} = 0$$

This identity was used to show that the two hidden angular momenta in hydrogen have the same magnitude.

4. Squeezing a Hamiltonian [15 points]

The squeezing operator $S(\gamma)$ of the harmonic oscillator is unitary

$$S(\gamma) = \exp\left(-\frac{\gamma}{2}(\hat{a}^{\dagger}\hat{a}^{\dagger} - \hat{a}\hat{a})\right), \ \gamma \in \mathbb{R}$$

Given the Hamiltonian

$$H(m,k) = \frac{\hat{p}^2}{2m} + \frac{1}{2}k\,\hat{x}^2, \quad \omega^2 \equiv \frac{k}{m},$$

it is then natural to ask what is

$$S^{\dagger}(\gamma) H(m,k) S(\gamma) = \cdots$$
 ?

- (a) Show that the right-hand side above is H(m', k') with some new parameters m' and k' that you should determine.
- (b) How are the energy levels of H(m, k) and H(m', k') related ? Given an eigenstate $|E\rangle$ of H(m, k), how would you build an energy eigenstate of H(m', k')?
- (c) Would you say that the action of $S(\gamma)$ is a symmetry of the harmonic oscillator? Explain.

5. A Three Dimensional Harmonic Oscillator [25 points]

Consider a particle of mass m in a three-dimensional isotropic harmonic oscillator:

$$\hat{H} = \hbar\omega \left(\hat{N}_1 + \hat{N}_2 + \hat{N}_3 + \frac{3}{2}\right) = \hbar\omega \left(\hat{N} + \frac{3}{2}\right),$$
$$\hat{N}_1 = \hat{a}_1^{\dagger}\hat{a}_1, \quad \hat{N}_2 = \hat{a}_2^{\dagger}\hat{a}_2, \quad \hat{N}_3 = \hat{a}_3^{\dagger}\hat{a}_3, \quad \hat{N} \equiv \hat{a}_i^{\dagger}\hat{a}_i, \quad [\hat{a}_i, \hat{a}_j^{\dagger}] = \delta_{ij}, \quad i, j = 1, 2, 3.$$
As usual \hat{a}_i^{\dagger} and \hat{a}_i are raising and lowering operators, \hat{N}_i are number operators and \hat{N} is the total number. Denoted in diagram constrained come. Consider the following price \hat{N}_i and \hat{N}_i are number operators.

is the *total* number. Repeated indices are summed over. Consider the following nine operators defined by

$$T_{ij} \equiv \hat{a}_i^{\dagger} \hat{a}_j, \quad i, j = 1, 2, 3.$$

- (a) Some T's are familiar. Write \hat{H} and the number operators \hat{N}_i in terms of T's.
- (b) The T's form a closed algebra under commutation (commutators of T's give linear combinations of T's). Confirm this by evaluating the commutator $[T_{ij}, T_{pq}]$.
- (c) Are the T operators symmetries? Answer this by computing $[\hat{H}, T_{pq}]$.
- (d) Write the state $\hat{a}_1^{\dagger}(\hat{a}_3^{\dagger})^2 |0\rangle$ as the action of some T acting on the state $(\hat{a}_3^{\dagger})^3 |0\rangle$.

We consider redefined creation and annihilation operators

$$\hat{a}_L \equiv \frac{1}{\sqrt{2}} (\hat{a}_1 + i\hat{a}_2), \quad \hat{a}_R \equiv \frac{1}{\sqrt{2}} (\hat{a}_1 - i\hat{a}_2), \quad [\hat{a}_L, \hat{a}_L^{\dagger}] = [\hat{a}_R, \hat{a}_R^{\dagger}] = 1,$$
$$\hat{N}_L = \hat{a}_L^{\dagger} \hat{a}_L, \quad \hat{N}_R = \hat{a}_R^{\dagger} \hat{a}_R, \quad \hat{N}_3 = \hat{a}_3^{\dagger} \hat{a}_3.$$

Any operator in $(\hat{a}_L, \hat{a}_L^{\dagger})$ commutes with any operator in $(\hat{a}_R, \hat{a}_R^{\dagger})$. With these operators the Hamiltonian and the (honest!) angular momentum operators are given by

$$\hat{H} = \hbar \omega \left(\hat{N}_L + \hat{N}_R + \hat{N}_3 + \frac{3}{2} \right), \hat{L}_z = \hbar \left(\hat{a}_R^{\dagger} \hat{a}_R - \hat{a}_L^{\dagger} \hat{a}_L \right) = \hbar (\hat{N}_R - \hat{N}_L), \hat{L}_+ = \sqrt{2} \hbar \left(\hat{a}_3^{\dagger} \hat{a}_L - \hat{a}_R^{\dagger} \hat{a}_3 \right), \hat{L}_- = \sqrt{2} \hbar \left(\hat{a}_L^{\dagger} \hat{a}_3 - \hat{a}_3^{\dagger} \hat{a}_R \right).$$

This basis of operators is best to analyze the angular momentum content of states.

- (e) Consider the set of degenerate states of total number N = n, with n some positive integer. These states organize into representations of angular momentum with various values of l. What is the highest value l_{*} of l? Explain. What are the other values of l that arise ? You need not explain this (we derived it in class). For the l = l_{*} multiplet construct explicitly the three states with largest value of L_z. For the multiplet with next highest value l_{**} of l, construct the state with maximal L_z. Don't bother to normalize any of the states.
- (f) Find an operator K that is T-like (namely, involves the product of one creation and one annihilation operator), commutes with the Hamiltonian, and acting on $|\ell_{**}, \ell_{**}\rangle$ gives $|\ell_*, \ell_*\rangle$. This operator "explains" the degeneracy of the ℓ_* and ℓ_{**} multiplets.

6. Conservation of Angular Momentum in Particle Decays. [20 points]

This problem deals with the decay of an unstable particle X to two daughter particles A and B. In the decay process $X \to AB$, total angular momentum is conserved. In the X rest frame, the total angular momentum $\mathbf{J} = \mathbf{S}_X$ is given just by the spin of particle X. After the decay, the total angular momentum receives three contributions:

$$\mathbf{J} = \mathbf{S}_A + \mathbf{S}_B + \mathbf{L},$$

where \mathbf{S}_A is the spin of particle A, \mathbf{S}_B is the spin of particle B, and \mathbf{L} is the relative orbital angular momentum between A and B. Because angular momentum is conserved in this decay, if the initial state is an eigenstate of \mathbf{J}^2 and \mathbf{J}_z , then the final state is also an eigenstate with the same eigenvalues.

- (a) Consider the case where X is spin-0 and both A and B are spin-1/2 (i.e. $s_X = 0$, $s_A = 1/2$, $s_B = 1/2$). What values of the orbital angular momentum ℓ are allowed, consistent with angular momentum conservation?
- (b) Repeat the above problem for $s_X = 3/2$, $s_A = 1/2$, and $s_B = 1$.
- (c) There are certain processes for which a two-body decay is forbidden. Explain why a neutron n cannot decay to a proton p and an electron e^- via the decay $n \to p e^-$, despite the fact that this decay is consistent with energy and charge conservation. Recall that both the proton and the neutron are spin-1/2 particles. [Note: the three-body decay $n \to p e^- \overline{\nu}$ is allowed, where $\overline{\nu}$ is a spin-1/2 antineutrino.]
- (d) A mystery particle X of unknown spin s_X is polarized such that $m_X = +s_X$. It decays via $X \to AB$, where $s_A = 1/2$, $s_B = 0$, but the relative orbital angular momentum ℓ is also unknown. (For simplicity, you may assume the decay yields a single value of ℓ .) After the decay, the z-component m_A of the spin of particle A is measured, and it is found to have probabilities:

$$P(m_A = 1/2) = 1/5, \qquad P(m_A = -1/2) = 4/5.$$

What is s_X and what is ℓ ? [Hint: you may want to consider the action of \mathbf{J}_+ on both the initial and final states.]

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