

MASSACHUSETTS INSTITUTE OF TECHNOLOGY
Physics Department

Physics 8.07: Electromagnetism II
Prof. Alan Guth

September 5, 2012

PROBLEM SET 1

DUE DATE: Friday, September 14, 2012. Either hand it in at the lecture, or by 5:00 pm in the 8.07 homework box.

READING ASSIGNMENT: Chapter 1 of Griffiths: *Vector Analysis*.

PROBLEM 1: VECTOR IDENTITIES INVOLVING CROSS PRODUCTS
(20 points)

In manipulating cross products, it is useful to define ε_{ijk} (the Levi-Civita antisymmetric symbol) to be:

$$\varepsilon_{ijk} = \begin{cases} +1 & \text{if } ijk = (123, 231, 312) \\ -1 & \text{if } ijk = (213, 321, 132) \\ 0 & \text{otherwise .} \end{cases} \quad (1.1.1)$$

That is, ε_{ijk} is nonzero only when all three indices are different; it is then equal to +1 if ijk is a cyclic permutation of 123, and -1 if ijk is an anti-cyclic permutation. Note that ε_{ijk} is totally antisymmetric, in the sense that it changes sign if any two indices are interchanged:

$$\varepsilon_{ijk} = -\varepsilon_{ikj} = \varepsilon_{kij} . \quad (1.1.2)$$

With this definition, the i^{th} component of the cross product of two vectors \vec{A} and \vec{B} can be written as

$$\left(\vec{A} \times \vec{B}\right)_i = \varepsilon_{ijk} A_j B_k , \quad (1.1.3)$$

where we have used the summation convention that repeated indices are summed over (that is, $\varepsilon_{ijk} A_j B_k = \sum_{j=1}^3 \sum_{k=1}^3 \varepsilon_{ijk} A_j B_k$). For the rest of this problem set, we will always assume that this summation convention is implied, unless explicitly stated otherwise.

(a) From the definition in Eq. (1.1.1), show that

$$\varepsilon_{ijk} \varepsilon_{inm} = \delta_{jn} \delta_{km} - \delta_{jm} \delta_{kn} , \quad (1.1.4)$$

where of course there is an implied sum over the i index in Eq. (1.1.4), but the indices j , k , n , and m are free.

(b) Using Eqs. (1.1.3) and (1.1.4), show that for any vectors \vec{A} , \vec{B} , and \vec{C} ,

$$\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B} (\vec{A} \cdot \vec{C}) - \vec{C} (\vec{A} \cdot \vec{B}) \quad (1.1.5)$$

(c) Using Eqs. (1.1.3) and (1.1.4), show that for any vectors \vec{A} and \vec{B} ,

$$\vec{\nabla} \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\vec{\nabla} \times \vec{A}) - \vec{A} \cdot (\vec{\nabla} \times \vec{B}) . \quad (1.1.6)$$

(d) Using Eqs. (1.1.3) and (1.1.4), show that for any vector \vec{A} ,

$$\vec{A} \times \vec{\nabla} \times \vec{A} = \frac{1}{2} \vec{\nabla} A^2 - (\vec{A} \cdot \vec{\nabla}) \vec{A} . \quad (1.1.7)$$

(e) Using Eqs. (1.1.3) and (1.1.4), show that for any vectors \vec{A} and \vec{B} ,

$$\vec{\nabla} \times (\vec{A} \times \vec{B}) = (\vec{B} \cdot \vec{\nabla}) \vec{A} - (\vec{A} \cdot \vec{\nabla}) \vec{B} + \vec{A} (\vec{\nabla} \cdot \vec{B}) - \vec{B} (\vec{\nabla} \cdot \vec{A}) . \quad (1.1.8)$$

PROBLEM 2: TRIPLE CROSS PRODUCTS (10 points)

Griffiths Problem 1.2 (p. 4), Griffiths Problem 1.6 (p. 8).

PROBLEM 3: PROPERTIES OF THE ROTATION MATRIX R (15 points)

Griffiths Eq. (1.31), p. 11, is

$$\bar{A}_i = \sum_{j=1}^3 R_{ij} A_j .$$

If we use the convention that repeated indices are summed over, then this can be written as

$$\bar{A}_i = R_{ij} A_j . \quad (1.2.1)$$

(a) Show that the elements (R_{ij}) of the three-dimensional rotation matrix must satisfy the constraint

$$R_{ij} R_{ik} = \delta_{jk} \quad (1.2.2)$$

in order to preserve the length of \vec{A} for all \vec{A} . Matrices satisfying Eq. (1.2.2) are called *orthogonal*. Here δ_{jk} is the Kronecker delta (δ_{jk} is 1 if $j = k$ and 0 otherwise), and we use the summation convention above.

(b) Using the orthogonality constraint (1.2.2), show that

$$A_i = R_{ji} \bar{A}_j . \quad (1.2.3)$$

Note that we can now show that $R_{ji} R_{ki} = \delta_{jk}$ using this relation, in a manner similar to the procedure in (a) (you do not have to show this).

(c) Using the chain rule for partial differentiation and the results of (b), show that if f is scalar function of $\vec{r} \equiv (x_1, x_2, x_3)$, then $\vec{\nabla} f(\vec{r})$ transforms as a vector; i.e., show that if

$$\bar{f}(\bar{x}_1, \bar{x}_2, \bar{x}_3) = f(x_1, x_2, x_3) , \quad (1.2.4)$$

where $\bar{x}_i = R_{ij} x_j$, then

$$\frac{\partial \bar{f}}{\partial \bar{x}_i} = R_{ij} \frac{\partial f}{\partial x_j} . \quad (1.2.5)$$

PROBLEM 4: USE OF THE GRADIENT (10 points)

Griffiths Problem 1.12 (p.15), Griffiths Problem 1.13 part (a) only (p.15).

PROBLEM 5: THE DIRAC DELTA FUNCTION AND $\nabla^2(1/4\pi r)$ (20 points)

One of the most used identities in this course is be the relation

$$-\nabla^2 \frac{1}{4\pi r} = -\vec{\nabla} \cdot \left[\frac{\vec{\nabla}}{4\pi r} \right] = \vec{\nabla} \cdot \left[\frac{\hat{r}}{4\pi r^2} \right] = \delta^3(r) = \delta(x) \delta(y) \delta(z) . \quad (1.5.1)$$

It turns out of course (see Griffiths 1.5.1, p. 45) that

$$-\nabla^2 \frac{1}{4\pi r}$$

is zero everywhere except at the origin, and ill-defined there. To get a better feel for the fact that

$$-\nabla^2 \frac{1}{4\pi r}$$

is a delta function, let's look at a different function which approaches $-(1/4\pi r)$ in some limit, but which is well-behaved everywhere. The function is

$$f_a(r) = -\frac{1}{4\pi} \frac{1}{\sqrt{r^2 + a^2}} . \quad (1.5.2)$$

For a nonzero, $f_a(r)$ is well-behaved everywhere, and

$$\lim_{a \rightarrow 0} f_a(r) = -\frac{1}{4\pi r} \quad (1.5.3)$$

- (a) Calculate $g_a(r) = \nabla^2 f_a(r)$ and show that it is also well behaved for all r . Sketch $g_a(r)$ for some value of a as a function of r/a .
- (b) Show that

$$\int_{\text{all space}} g_a(r) d^3x = 1 . \quad (1.5.4)$$

- (c) Show that

$$\lim_{a \rightarrow 0} g_a(r) = 0 \text{ if } r \neq 0 . \quad (1.5.5)$$

Thus in the limit that a goes to zero, our well-behaved function $g_a(r)$ exhibits the properties we expect of a three-dimensional delta function.

PROBLEM 6: EXERCISES WITH δ -FUNCTIONS (10 points)

- (a) A charge Q is spread uniformly over a spherical shell of radius R . Express the volume charge density using a delta function in spherical coordinates. Repeat for a ring of radius R with charge Q lying in the xy plane.
- (b) In cartesian coordinates, we can write $\delta^3(\vec{r} - \vec{r}') = \delta(x - x')\delta(y - y')\delta(z - z')$. How would one express $\delta^3(\vec{r} - \vec{r}')$ in cylindrical coordinates ($s, \phi, \text{and } z$).
- (c) A charge λ per unit length is distributed uniformly over a cylindrical surface of radius b . Give the volume charge density using a delta function in cylindrical coordinates
- (d) What is $\nabla^2 \ln r$ in two dimensions? (Here r is the radial coordinate, $r = \sqrt{x^2 + y^2}$.)

PROBLEM 7: COROLLARIES OF THE FUNDAMENTAL INTEGRAL THEOREMS (15 points)

This problem is closely related to Problem 1.60, p. 56 of Griffiths. You will find useful hints there— but try without hints first!! Show that:

- (a) $\int_V \vec{\nabla} \psi d^3x = \int_S \psi d\vec{a}$, where S is the surface bounding the volume V . Show that as a consequence of this, $\int_S d\vec{a} = 0$ for a closed surface S .
- (b) $\int_V \vec{\nabla} \times \vec{A} d^3x = - \int_S \vec{A} \times d\vec{a}$, where S is the surface bounding the volume V .
- (c) $\int_S \vec{\nabla} \psi \times d\vec{a} = - \oint_{\Gamma} \psi d\vec{l}$, where Γ is the boundary of the surface S .
- (d) For a closed surface S , one has $\int_S (\vec{\nabla} \times \vec{A}) \cdot d\vec{a} = 0$.

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