# Classical Mechanics III (8.09) Fall 2014 Assignment 3 

Massachusetts Institute of Technology
Physics Department
Due September 29, 2014
September 22, 2014
6:00pm

## Announcements

This week we continue our discussion of Rigid Bodies, including the general example of a heavy symmetric top precessing and nutating.

## Reading Assignment for this week

- The main readings for Rigid Bodies are Goldstein Ch.4, sections 4.1, 4.2, 4.4, 4.6, 4.9. (Note that 4.3 is a linear algebra review.) The Ch. 4 reading is on rigid body kinematics where many physics topics will be familiar to you. Goldstein emphasizes vector notation and discusses rotations as matrices. Next read Ch.5, sections 5.1, 5.3-5.7. Ch. 5 is on rigid body dynamics. We will not cover Poinsot's construction, so you may skip this material on pages 201-205 of section 5.6. (Other sections from Ch. 4 and Ch. 5 may be of interest, but the above are the most important ones.)
- If you find the reading in Goldstein too dense, you should consider reading Thornton \& Marion Ch.11. I particularly recommend Ch. 11 section 10 on the force-free motion of a symmetric top.
- Our next item to discuss will be principal axes for oscillating motion. If you would like to read ahead, the reading for this will be Goldstein Ch. 6 sections 6.1-6.4.


## Problem Set 3

In these five problems you will study the dynamics of rigid bodies and the use of rotating coordinate systems. If you can obtain a result using symmetry you should!

1. Rotation Angle in the Euler Theorem [10 points]

In lecture we demonstrated that a general rotation can be thought of as a simple rotation about some fixed axis. In this problem you will fill in the last step of this analysis, namely showing that the angle $\Phi$ that appeared in those calculations is a rotation angle.

Lets first recall the setup. Consider a general rotation $U$ so that $\vec{r}^{\prime}=U \vec{r}$ with $U$ an orthogonal matrix. Let $\vec{\xi}_{i}$ be the eigenvectors satisfying $U \vec{\xi}_{i}=\lambda_{i} \vec{\xi}_{i}$ where

$$
\lambda_{1}=e^{i \Phi}, \quad \lambda_{2}=e^{-i \Phi}, \quad \lambda_{3}=1, \quad \vec{\xi}_{i}^{\dagger} \cdot \vec{\xi}_{j}=\delta_{i j}
$$

Here $\Phi \neq 0$, and $\dagger$ means the complex conjugate and transpose. Also $\vec{\xi}_{3}$ is real, $\vec{\xi}_{3}^{*}=\vec{\xi}_{3}$, whereas the other two eigenvectors are complex satisfying $\vec{\xi}_{1}^{*}=\vec{\xi}_{2}$. Finally, the matrix $X=\left(\vec{\xi}_{1} \vec{\xi}_{2} \vec{\xi}_{3}\right)$ is unitary, $X^{\dagger}=X^{-1}$ and

$$
X^{\dagger} U X=\left(\begin{array}{ccc}
\lambda_{1} & 0 & 0 \\
0 & \lambda_{2} & 0 \\
0 & 0 & \lambda_{3}
\end{array}\right)
$$

In order to demonstrate that $U$ describes a simple rotation with rotation angle $\Phi$ we will find a set of coordinates that make this obvious. Since $U \vec{\xi}_{3}=\vec{\xi}_{3}$, the rotation is about the axis given by the real vector $\vec{\xi}_{3}$, so we will pick $\vec{\xi}_{3}$ as our new $z$ axis.
(a) [4 points] A natural choice for the other two axes might be $\vec{\xi}_{1}$ and $\vec{\xi}_{2}$. In this case the transformation to the new coordinates would be $\vec{r}=X \vec{s}$, where in the new coordinate system $\vec{\xi}_{1}=\left(\begin{array}{l}1 \\ 0 \\ 0\end{array}\right), \vec{\xi}_{2}=\left(\begin{array}{l}0 \\ 1 \\ 0\end{array}\right), \vec{\xi}_{3}=\left(\begin{array}{l}0 \\ 0 \\ 1\end{array}\right)$. BUT in our original coordinates $\vec{\xi}_{1}$ and $\vec{\xi}_{2}$ are complex so they are NOT valid choices for the axes. Find two real vectors, $\vec{\xi}_{a}$ and $\vec{\xi}_{b}$, that are linear combinations of $\vec{\xi}_{1}$ and $\vec{\xi}_{2}$, and are such that $\left\{\vec{\xi}_{a}, \vec{\xi}_{b}, \vec{\xi}_{3}\right\}$ are a normalized orthogonal set of vectors. This is the basis that we use as the axes for our new coordinates.
(b) [2 points] Write the relation between components in the basis $\left\{\vec{\xi}_{a}, \vec{\xi}_{b}, \vec{\xi}_{3}\right\}$ and $\left\{\vec{\xi}_{1}, \vec{\xi}_{2}, \vec{\xi}_{3}\right\}$ as a matrix $W$, and show that $W$ is unitary.
(c) [4 points] Consider the new coordinates whose components are described by the transformation $\vec{u}=W \vec{s}=W X^{\dagger} \vec{r}$. Write the rotation $\vec{r}^{\prime}=U \vec{r}$ as a relation $\vec{u}^{\prime}=\widetilde{U} \vec{u}$, and show that the matrix $\widetilde{U}$ has the form of a standard rotation matrix with rotation angle $\Phi$.

Note: $\operatorname{Tr}(U)=\operatorname{Tr}(\widetilde{U})=1+2 \cos \Phi$, so given a matrix $U$ we can easily determine $\Phi$ without going through transformations.

## 2. Foucault Pendulum and the Coriolis Effect [13 points]

Consider a pendulum consisting of a long massless rod of length $\ell$ attached to a mass $m$. The pendulum is hung in a tower that is at latitude $\lambda$ on the earth's surface, so it is natural to describe its motion with coordinates fixed to the rotating Earth. Let $\omega$ be the Earth's angular velocity. Use the spherical coordinates $(r, \theta, \phi)$ shown in the figure to investigate the Coriolis force. Here $\hat{z}$ is perpendicular to the Earth's surface and $\hat{y}$ is tangent to a circle of constant longitude that passes through the north pole.
a) [9 points] The velocity is given in terms of $\vec{v}$ in the rotating frame by $\vec{v}+\vec{\omega} \times\left(R_{e} \hat{z}+\vec{r}\right)$, so

$$
L=\frac{m}{2}\left[\vec{v}+\vec{\omega} \times\left(\vec{r}+R_{e} \hat{z}\right)\right]^{2}-V,
$$



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with $R_{e}$ the radius of the earth, and $V$ the potential energy due to gravity near the earth's surface (we neglect air resistance). Writing everything in terms of the variables $\theta$ and $\phi$, and the fixed angle $\lambda$, derive the equations of motion for the pendulum. From the start you should only keep terms up to first order in $\omega$. You can also drop the term $\propto \omega R_{e}$ since it is a total time derivative.
b) [4 points] Since $\ell$ is large, consider the small angle approximation for $\theta$ and simplify your equations of motion from a). Demonstrate that the pendulum undergoes precession with $\dot{\phi}=\omega \sin \lambda$.
3. Angular Velocity with Euler Angles [9 points]
(a) [2 points] Show that the components of angular velocity along the body axes $\left(x^{\prime}, y^{\prime}, z^{\prime}\right)$ are given in terms of Euler angles by

$$
\begin{aligned}
& \omega_{x^{\prime}}=\dot{\phi} \sin \theta \sin \psi+\dot{\theta} \cos \psi \\
& \omega_{y^{\prime}}=\dot{\phi} \sin \theta \cos \psi-\dot{\theta} \sin \psi \\
& \omega_{z^{\prime}}=\dot{\phi} \cos \theta+\dot{\psi}
\end{aligned}
$$

This is done in the text! I am asking you to go through the steps to ensure that you understand the calculation.
(b) [4 points] Show that the components of angular velocity along the fixed space set of axes, the inertial frame $(x, y, z)$, are given in terms of the Euler angles by

$$
\begin{aligned}
& \omega_{x}=\dot{\theta} \cos \phi+\dot{\psi} \sin \theta \sin \phi \\
& \omega_{y}=\dot{\theta} \sin \phi-\dot{\psi} \sin \theta \cos \phi \\
& \omega_{z}=\dot{\psi} \cos \theta+\dot{\phi}
\end{aligned}
$$

This problem is Goldstein Ch.4\#14. You may use results given in Goldstein.
(c) [3 points] Using the generalized coordinate $\psi$, derive an Euler equation of motion using the Euler-Lagrange equation. Use the form with generalized forces $Q_{j}$ :

$$
\frac{d}{d t} \frac{\partial T}{\partial \dot{q}_{j}}-\frac{\partial T}{\partial q_{j}}=Q_{j} .
$$

## 4. Point Mass on a Disk [12 points]

A thin uniform disk of radius $R$ and mass $M$ lies in the $x-y$ plane, and has a point mass $m=3 M / 8$ attached on its edge. (There is no gravity in this problem.)
(a) [4 points] Find the moment of inertia tensor of the disk about its center (ignoring the mass $m$ ). Then
 find the moment of inertia tensor of the combined system of the disk and point mass about the point $A$ in the figure.
(b) [4 points] Find the principal moments of intertia and the principal axes about $A$. (Recall that you may use mathematica or matlab. If you do then you should still write out intermediate steps.)
(c) [4 points] The disk is constrained to rotate about the $y$-axis with angular velocity $\omega$ by pivots at $A$ and $B$. What is the angular momentum about $A$ as a function of time?
5. A Rolling Cone [16 points] (Adapted from Goldstein Ch. 5 \#17)

A uniform right circular cone of height $h$, half-angle $\alpha$, and density $\rho$ rolls on its side without slipping on a uniform horizontal plane. It returns to its original position in a time $\tau$.
(a) [5 points] Find the CM of the cone. Find the moment of inertia tensor for body axes centered on the tip with the $y^{\prime}$-axis going through the CM.
(b) [2 points] What is the moment of inertia tensor if we move our axes in (a) so they are centered on the CM?
(c) [4 points] Now assume the cone rolls on a fixed plane. Pick a new set of body axes $(x, y, z)$ such that the $z$-axis is perpendicular to the plane, the $y$-axis coincides with the instantaneous line of contact, and the origin is the tip of the cone. Find the moment of inertia tensor for these axes.
(d) [5 points] Find the kinetic energy of the rolling cone. [There are two ways you could answer this, one uses your results from (b) and one uses those from (c).]

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