

## MITOCW | 7.2 Relativistic Doppler Effect

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**MARKUS**

Welcome back to 8.20, Special Relativity. In this section, we're going to talk about the relativistic Doppler effect.

**KLUTE:**

And we make good use of our space-time diagrams, which we discussed earlier.

So the situation is as follows-- to simplify this, we have a source which is emitting pulses. So the waves are pulses. Every now and then there is a beep, and another beep, and another beep. And those pulses travel with their velocity-- with their wave velocity.

And they have a world line represented here in the space-time diagram. This is pulse number one, and this is pulse number two. The distance between those two pulses is our period, the period of our wave, which we call  $\tau$ . The question now is, how is this being observed by an observer which is moving with a relative velocity  $v$  with respect to the source?

So let's analyze this. So if we want to characterize or find our position  $x_1$  and  $x_2$ , we can do this by saying  $x_1$  is equal to  $ct_1$  or equal to  $x_0$ , which is the distance of the observer to the source plus  $c$  times  $t_1$ .  $v$  is the velocity in which the source is moving.

And similarly for  $t_2$ , we find  $c$  times  $t_2$  minus  $\tau$ . And that's also equal to  $x_0$  plus  $v$  times  $t_2$ . So the distance in time-- we're still in the reference frame as of the source-- is given by  $c$  times  $\tau$  over  $c$  minus  $v$ . And the distance in space is given by  $v$  times  $c$  times  $\tau$  over  $c$  minus  $v$ .

So the question is not how this observed-- how this is seen by the source but how this is being seen by the observer. So we have to apply Lorentz transformation. So in the  $s$  prime frame, which is the observer frame, we find  $\Delta t'$  is equal to  $\gamma \Delta t$  minus  $v$  over  $c$  squared  $\Delta x$ . And then we just fill in the information as we discussed before.

$\tau'$  is then  $\gamma$  times  $c \tau$  over  $c$  minus  $v$  times  $1$  minus  $v$  square over  $c$  square. And then you make use of  $\Delta x$  equal  $v \tau$  over  $c$ . And we make use of  $\gamma$  equals  $1$  over square root of  $1$  minus  $\beta$  square.

And we find then-- this is a little bit of an algebra exercise here-- that the period now is given by  $1$  plus  $\beta$  over  $1$  minus  $\beta$  square root of that times  $\tau$ . And the frequency is the inverse. We'll have  $1$  minus  $\beta$  over  $1$  plus  $\beta$  square root of that [ $\tau$  times  $\tau$ ] the frequency. So we just calculated relativistically how the period and the frequency of a wave is Lorentz transformed.