# 8.701

Introduction to Nuclear and Particle Physics

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8. Neutrinos

8.2 Neutrino Mass

### **Neutrino Masses**

Mass terms can be constructed in different ways by introducing sterile neutrinos

$$-\mathcal{L}_{M_{\nu}} = M_{Dij}\bar{\nu}_{si}\nu_{Lj} + \frac{1}{2}M_{Nij}\bar{\nu}_{si}\nu_{sj}^{c} + \text{h.c.}$$

The first **"Dirac"** terms is generated after spontaneous electroweak symmetry breaking from Yukawa interactions as we have seen for charged leptons

$$Y_{ij}^{\nu}\bar{\nu}_{si}\tilde{\phi}^{\dagger}L_{Lj} \Rightarrow M_{Dij} = Y_{ij}^{\nu}\frac{v}{\sqrt{2}}$$

It conserved lepton number but breaks lepton flavour number symmetries.

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## **Dirac Neutrinos**

Identify sterile neutrinos as right-handed component of four-spinor neutrino field

$$-\mathcal{L}_{M_{\nu}} = \sum_{k=1}^{3} m_k \bar{\nu}_{Dk} \nu_{Dk}$$

$$\nu_{Dk} = (V^{\nu \dagger} \vec{\nu}_L)_k + (V_R^{\nu \dagger} \vec{\nu}_s)_k$$

Weak-doublet components

$$\nu_{Li} = P_L \sum_{j=1}^3 V_{ij}^{\nu} \nu_{Dj} \,. \qquad i = 1, 2, 3$$

#### **Neutrino Masses**

$$-\mathcal{L}_{M_{\nu}} = M_{Dij}\bar{\nu}_{si}\nu_{Lj} + \frac{1}{2}M_{Nij}\bar{\nu}_{si}\nu_{sj}^{c} + \text{h.c.}$$

2nd term is a "Majorana" mass term.

It is a singlet of the SM gauge group and can appear as a bare mass term

It involved two neutrino (right-handed) fields and breaks lepton number.

Rewrite 
$$-\mathcal{L}_{M_{\nu}} = \frac{1}{2} (\overline{\vec{\nu}_{L}^{c}}, \overline{\vec{\nu}_{s}}) \begin{pmatrix} 0 & M_{D}^{T} \\ M_{D} & M_{N} \end{pmatrix} \begin{pmatrix} \vec{\nu}_{L} \\ \vec{\nu}_{s}^{c} \end{pmatrix} + \text{h.c.} \equiv \overline{\vec{\nu}^{c}} M_{\nu} \vec{\nu} + \text{h.c.}$$

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## See-saw Mechanism

to  $M_{\rm M}^{-1}$ 

If the mass eigenstate of  $M_N$  are much higher than the EW scale, the diagonalization of  $M_V$  leads to three light neutrinos  $v_1$  and potentially numerous heavy neutrinos N.

$$-\mathcal{L}_{M_{\nu}} = \frac{1}{2}\bar{\nu}_l M^l \nu_l + \frac{1}{2}\bar{N}M^h N$$

$$\begin{split} M^l \simeq -V_l^T M_D^T M_N^{-1} M_D V_l, \qquad M^h \simeq V_h^T M_N V_h \,, \\ V^\nu \simeq \begin{bmatrix} \left(1 - \frac{1}{2} M_D^{\dagger} M_N^{*-1} M_N^{-1} M_D\right) V_l & M_D^{\dagger} M_N^{*-1} V_h \\ -M_N^{-1} M_D V_l & \left(1 - \frac{1}{2} M_N^{-1} M_D M_D^{\dagger} M_N^{*-1}\right) V_h \end{bmatrix} \\ \end{split}$$
 Mass of heavier states proportional to  $M_N$  while light state

## Light sterile neutrinos

If the scale of some eigenvalues of MN are night higher than the EW scale, the low energy spectrum contains additional light states with large admixture of sterile components, i.e. more than 3 light neutrinos. Both light and heavy neutrinos are Majorana particles. MIT OpenCourseWare <u>https://ocw.mit.edu</u>

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