Assignment 6: MCMC Simulation and Review Problems

Preface: In this assignment, you will use computational software to implement the Markov Chain Monte Carlo (MCMC) simulation algorithm we simulated ourselves in class.

1. MCMC of Spins in a Magnetic Field

In this problem, we simulate spin system depicted in Fig. 1.



Figure 1: Particular microstate of the system we are simulating.

Using the partition function, we can show that for a collection of N spins each of which has magnetic dipole μ and subject to magnetic field H, the average spin is given by

$$\langle s \rangle = \tanh\left(\frac{\mu H}{k_B T}\right),$$
(1)

where $tanh(x) = (e^x - e^{-x})/(e^x + e^{-x})$. We will use a simulation to compare the theoretical result Eq.(1) to one obtains from computation.

We will be using *Mathematica* to complete this problem, so we need to establish some code preliminaries.

- (i) Log in to your account in one of MIT's Athena Clusters, and go to the MITES 2018 Physics III course website.
- (ii) Download the code spin_model_simulation.nb from the course webpage and open it in *Mathematica*.
- (iii) Select a block of code and run it by pressing Shift+Enter.

Now we begin the problem

- (a) Run each line of the "Function Definitions" section of the code to determine what the associated function does
- (b) Low Temperature Simulation

Fill in all the "..." in the code. In particular define/fill in the following:

- Parameter definition $\mu = 1.0$, H = 1.0, $k_B T = 0.1$
- There are $N_{\text{spins}} = 1000$ in the system
- There are $N_{\text{steps}} = 10,000$
- An initial spin configuration with N_{spins} all of which have value -1;

• Energy of microstate is given by

$$E(\{s_i\}) = -\mu H \sum_{i=1}^{N} s_i$$
(2)

• The ratio between Boltzmann factors for final and initial spin configurations

$$\exp\left[-\frac{1}{k_BT}\left(E(\{s_i\}_{\text{final}}) - E(\{s_i\}_{\text{initial}})\right)\right]$$
(3)

- The theoretical value of $\langle s \rangle$ obtained from Eq.(1)
- (c) Low Temperature Simulation

Perform a new simulation (i.e., all the parts of (b)) at the higher temperature $k_B T = 10.0$. *Hint: Copy and paste is your friend here.*

- (d) In the simulations, why do we only take the final 100 time steps, to compute the average spin?
- (e) In (b) and (c), how do the theoretical values of $\langle s \rangle$ compare with the computational results? How might we improve the correspondence between the two.

2. Optional Bonus: MCMC of Mean-Field Ising Model

Simulate the Mean-Field Ising Model for J = 1.0 and *computationally* compute average spin at the two temperature $k_BT = 0.5$ and $k_BT = 1.5$. Compare these computational results with Figure 12(b) in Lecture Notes 05.

Review Problems

(Not to submit)

Preface: This assignment contains a list of review problems to help you prepare for the final exam.

1. Protein Expression and Probability

For a particular model of a gene in a cell, the probability density that said gene produces a concentration of x proteins during the cell cycle is given by

$$p(x) = A\left(\frac{x}{b}\right)^N e^{-x/b},\tag{4}$$

where *b* is a biological constant with units of concentration and *A* is a normalization parameter.

- (a) The concentration of proteins that can be produced ranges from zero to infinite. What must *A* be in order for Eq.(4) to be normalized?
- (b) What is the mean of the normalized probability density?
- (c) What is the standard deviation of the normalized probability density?

2. Adsorption



Figure 2: Particular microstate of system. The possible surface states are shown in yellow. The possible bulk states are the remaining states.

A gas is in a container, and some of the atoms are stuck (i.e., "adsorped") to the surface of the container, and others are free to move in the volume. We will study this system using a discrete model of the positions the atoms can take on.

Let *M* be the number of positions an atom can occupy while in the bulk, and let *N* be the number of positions an atom can occupy while on the surface. We have *N* atoms so that the surface states can be completely filled. Let $k \le N$ be the number of atoms that are on the surface. A *single* atom has energy $-\varepsilon$ while it is on the surface and energy 0 while it is in the bulk. *M*, *N*, and ε are all constants, and *k* is a random variable. *M*, *N*, and *k* are all much greater than 1.

- (a) How many microstates does the system have for a particular value of *k*?
- (b) What is the free energy of this system? You can write this free energy in terms of factorials.
- (c) Use Stirling's approximation $\ln N! \simeq N \ln N N$ to approximate the result in (b). Derive **one** condition that *k* must satisfy in order to define the thermal equilibrium of the system. We often state two conditions associated with thermal equilibrium. Only apply the first-derivative condition.

3. Nucleotide zipper



Figure 3: "Zipper" of nucleotides

A strand of DNA can be modeled as a zipper with links between N base-pairs. Each link has the energy 0 in the closed state and the energy ε the open state. Assume the zipper can only open from one end, so that the *n*th link can open only if all the links 1, 2, ..., n - 1 are open. Assume that the system is at a temperature T.

- (a) How many microstates does this system have? What are the energies of these microstates?
- (b) Find the partition function for a single DNA strand as a function of *T* and ε . *Hint: Evaluate the summation using the geometric series identity* $\sum_{n=0}^{N} x^n = (1 x^{N+1})(1 x)^{-1}$.
- (c) Assume $N \gg 1$. Find the average number of open links as a function of *T* and ε .

4. Entropy of an Ideal Gas

For a system in thermal equilibrium defined by the partition function Z, we determined that the Helmholtz free energy of the system is

$$F = -k_B T \ln Z = -k_B T \ln \sum_{\{i\}} e^{-\beta E_i},\tag{5}$$

where *i* denotes a microstate of the system, and E_i is the energy of that microstate.

(a) For a system in thermal equilibrium, the Gibbs entropy is given by

$$S = -k_B T \sum_{\{i\}} p_i \ln p_i = -k_B \sum_{\{i\}} \frac{e^{-\beta E_i}}{Z} \ln\left(\frac{e^{-\beta E_i}}{Z}\right).$$
 (6)

Given Eq.(5), show that

$$S = -\frac{\partial F}{\partial T}.$$
(7)

Note: $\beta = 1/k_BT$.

(b) Using Eq.(7), show that the entropy of an ideal gas for $N \gg 1$ is

$$S \simeq Nk_B \left[\ln\left(\frac{V}{N}\right) - \frac{3}{2}\ln\left(\frac{h^2}{2\pi mk_BT}\right) + \frac{5}{2} \right].$$
(8)

Note: You will need to use Stirling's approximation $\ln N! \simeq N \ln N - N$

5. **Partition Function as an Integral**

The partition function for a system of N spins is given by

$$Z(T) = \int_{-\infty}^{\infty} d\rho \, \exp\left[-\beta N f(\rho, T)\right],\tag{9}$$

where

$$f(\rho, T) = -a(T - T_c)\rho^2 + c\rho^4.$$
(10)

The parameters a, c, and T_c are constants in the model. T stands in for temperature.

Approximate the partition function using Laplace's method. For what values of *T* is the approximation invalid?

6. Transition probabilities and equilibrium probabilities

A person is in one of four rooms, and after each time step has a probability 1/(M + 1) of moving to an adjacent room where M is the number of adjacent rooms; The probability of remaining in place is 1/(M + 1).



Figure 4: Four connected rooms. The person has a probability 1/(M + 1) of remaining in place and a probability of 1/(M + 1) of moving to any one of the adjacent rooms.

(a) Letting $\pi_{i \to j}$ represent the probability of transitioning from room *i* to room *j* in a single time step, fill in the elements below

$\pi_{1 \to 1} =$	$\pi_{1\to 2} =$	$\pi_{1\to 3} =$	$\pi_{1 \rightarrow 4} =$	
$\pi_{2 \rightarrow 1} =$	$\pi_{2 \rightarrow 2} =$	$\pi_{2 \rightarrow 3} =$	$\pi_{2 \rightarrow 4} =$	(11)
$\pi_{3 \rightarrow 1} =$	$\pi_{3\rightarrow 2} =$	$\pi_{3 \rightarrow 3} =$	$\pi_{3 \rightarrow 4} =$	
$\pi_{4 \rightarrow 1} =$	$\pi_{4\rightarrow 2} =$	$\pi_{4\rightarrow 3} =$	$\pi_{4 \rightarrow 4} =$	

(b) Using the relationship between equilibrium probabilities and transition probabilities (along with the fact that the equilibrium probabilities are normalized), compute the equilibrium probabilities p_1^{eq} , p_2^{eq} , p_3^{eq} , and p_4^{eq} , to be in rooms 1, 2, 3, and 4.

Resource: Introduction to Statistical Physics Mobolaji Williams

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