Ordinary Differential Equations from Physical Models

In engineering and physics, modeling physical phenomena is the means by which technological and natural phenomena are understood and predicted. A model is an abstraction of a physical system, often with simplifying assumptions, into a mathematical framework. Every model should be verifiable by an experiment that, to the greatest extent possible, satisfies the approximations that were used to obtain the model.

In the context of modeling, differential equations appear frequently. Learning how to model new and interesting systems is a learned skill—it is best to learn by following a few examples. Grain growth provides some interesting modeling examples that result in first-order ODES.

Grain Growth

In materials science and engineering, a grain usually refers to a single element in an ensemble that comprises a polycrystal. In a single phase polycrystal, a grain is a contiguous region of material with the same crystallographic orientation. It is separated from other grains by grain boundaries where the crystallographic orientation changes abruptly.

A grain boundary contributes extra free energy to the entire system that is proportional to the grain boundary area. Thus, if the boundary can move to reduce the free energy it will.

Consider simple, uniformly curved, isolated two- and three-dimensional grains.
Figure 20-1: Illustration of a two-dimensional isolated circular grain and a three-dimensional isolated spherical grain. Because there is an extra energy in the system $\Delta G_{2D} = 2\pi R \gamma_{gb}$ and $\Delta G_{3D} = 4\pi R^2 \gamma_{gb}$, there is a driving force to reduce the radius of the grain. A simple model for grain growth is that the velocity (normal to itself) of the grain boundary is $v_{gb} = M_{gb} \gamma_{gb} \kappa$ where $M_{gb}$ is the grain boundary mobility and $\kappa$ is the mean curvature of the boundary. The normal velocity $v_{gb}$ is towards the center of curvature.

A relevant question is “how fast will a grain change its size assuming that grain boundary migration velocity is proportional to curvature?”

For the two-dimensional case, the rate of change of area can be formulated by considering the following illustration.

Figure 20-2: A segment of a grain boundary moving with normal velocity $v_n$ will move a distance $v_n \Delta t$ in a short time $\Delta t$. The motion will result in a change of area $-\Delta A$ for the shrinking grain. Each segment, $ds$, of boundary contributes to the loss of area by $\Delta A = -v_n \Delta t ds$. 
Because for a circle, the curvature is the same at each location on the grain boundary, the curvature is uniform and \( v_n = M_{gb} \kappa_{gb} \gamma_{gb} = M_{gb} \gamma_{gb} / R \). Thus
\[
\frac{dA}{dt} = -M_{gb} \gamma_{gb} \frac{1}{R} 2\pi R = -2\pi M_{gb} \gamma_{gb}
\]
(20-1)

Thus, the area of a circular grain changes at a constant rate, the rate of change of radius is:
\[
\frac{dA}{dt} = \frac{d\pi R^2}{dt} = 2\pi R \frac{dR}{dt} = -2\pi M_{gb} \gamma_{gb}
\]
(20-2)

which is a first-order, separable ODE with solution:
\[
R^2(t) - R^2(t = 0) = -2M_{gb} \gamma_{gb} t
\]
(20-3)

For a spherical grain, the change in volume \( \Delta V \) due to the motion of a surface patch \( dS \) in a time \( \Delta t \) is \( \Delta V = v_n \Delta t dS \). The curvature of a sphere is
\[
\kappa_{sphere} = \left( \frac{1}{R} + \frac{1}{R} \right)
\]
(20-4)

Therefore the velocity of the interface is \( v_n = 2M_{gb} \gamma_{gb} / R \). The rate of change of volume due to the contributions of each surface patch is
\[
\frac{dV}{dt} = -M_{gb} \gamma_{gb} \frac{2}{R} 4\pi R^2 = -8\pi M_{gb} \gamma_{gb} R = -4(6\pi^2)^{1/3} M_{gb} \gamma_{gb} V^{1/3}
\]
(20-5)

which can be separated and integrated:
\[
V^{2/3}(t) - V^{2/3}(t = 0) = -\text{constant}_1 t
\]
(20-6)
or
\[
R^2(t) - R^2(t = 0) = -\text{constant}_2 t
\]
(20-7)

which is the same functional form as derived for two-dimensions.

The problem (and result) is more interesting if the grain doesn’t have uniform curvature.
Figure 20-3: For a two-dimensional grain with non-uniform curvature, the local normal velocity (assumed to be proportional to local curvature) varies along the grain boundary. Because the motion is in the direction of the center of curvature, the velocity can be such that its motion increases the area of the interior grain for some regions of grain boundary and decreases the area in other regions.

However, it can still be shown that, even for an irregularly shaped two-dimensional grain, \( A(t) - A(t = 0) = -(\text{const})t \).
Integrating Factors, Exact Forms

**Exact Differential Forms** .................................................................

In classical thermodynamics for simple fluids, expressions such as

\[
dU = TdS - PdV
\]

\[
= \left( \frac{\partial U}{\partial S} \right)_V dS + \left( \frac{\partial U}{\partial V} \right)_S dV
\]

\[
= \delta q + \delta w
\]

represent the differential form of the combined first and second laws of thermodynamics. If \(dU = 0\), meaning that the differential Eq. 20-8 is evaluated on a surface for which internal energy is constant, \(U(S, V) = \text{const}\), then the above equation becomes a *differential form*

\[
0 = \left( \frac{\partial U}{\partial S} \right)_V dS + \left( \frac{\partial U}{\partial V} \right)_S dV
\]

(20-9)

This equation expresses a relation between changes in \(S\) and changes in \(V\) that are necessary to remain on the surface \(U(S, V) = \text{const}\).

Suppose the situation is turned around and you are given the first-order ODE

\[
\frac{dy}{dx} = \frac{M(x, y)}{N(x, y)}
\]

(20-10)

which can be written as the differential form

\[
0 = M(x, y)dx + N(x, y)dy
\]

(20-11)

Is there a function \(U(x, y) = \text{const}\) or, equivalently, is it possible to find a curve represented by \(U(x, y) = \text{const}\)?
If such a curve exists then it depends only on one parameter, such as arc-length, and on that curve \( dU(x, y) = 0 \).

The answer is, “Yes, such a function \( U(x, y) = \text{const} \) exists if and only if \( M(x, y) \) and \( N(x, y) \) satisfy the Maxwell relations"

\[
\frac{\partial M(x, y)}{\partial y} = \frac{\partial N(x, y)}{\partial x}
\]  

(20-12)

Then if Eq. 20-12 holds, the differential form Eq. 20-11 is called an exact differential and a \( U \) exists such that \( dU = 0 = M(x, y)dx + N(x, y)dy \).

For fixed number of moles of ideal gas, the internal energy is a function of the temperature only, \( U(T) - U(T_o) = C_V(T - T_o) \). Consider the heat that is transfered to a gas that changes it temperature and volume a very small amount:

\[
dU = C_V dT = \delta q + \delta w = \delta q - PdV
\]

\[
\delta q = C_V dT + PdV
\]

(20-13)

Can a Heat Function \( q(T, V) = \text{constant} \) be found?

To answer this, apply Maxwell’s relations.

A linear differential equation is one that does not contain any powers (greater than one) of the function or its derivatives. The most general form is:

\[
Q(x) \frac{dy}{dx} + P(x)y = R(x)
\]

(20-14)
Equation 20-15 can always be reduced to a simpler form by defining $p = P/Q$ and $r = R/Q$.

$$\frac{dy}{dx} + p(x)y = r(x)$$

(20-15)

If $r(x) = 0$, Eq. 20-15 is said to be a homogeneous linear first-order ODE; otherwise Eq. 20-15 is a heterogeneous linear first-order ODE.

The reason that the homogeneous equation is linear is because solutions can superimposed—that is, if $y_1(x)$ and $y_2(x)$ are solutions to Eq. 20-15, then $y_1(x) + y_2(x)$ is also a solution to Eq. 20-15. This is the case if the first derivative and the function are themselves linear. The heterogeneous equation is also called linear in this case, but it is important to remember that sums and/or multiples of heterogeneous solutions are also solutions to the heterogeneous equation.

The homogeneous equation has a solution of the form

$$y(x) = \text{const } e^{-\int p(x)dx}$$

(20-16)
### Mathematica Example: Lecture-20

#### DSolve in Homogeneous and Heterogeneous ODEs

1. Show how DSolve solves the homogeneous equation

   \[
   \frac{dy}{dx} + p(x)y = 0
   \]

2. Show how DSolve solves the heterogeneous equation

   \[
   \frac{dy}{dx} + p(x)y = r(x)
   \]

3. Note how the homogeneous solution is one of the terms in the sum for the heterogeneous solution.

A trick (or, an integrating factor which amounts to the same thing) can be employed to find the solution to the heterogeneous equation. Multiply both sides of the heterogeneous equation by \(e^{\int p(x)dx}\):\(^{10}\)

\[
e^{\int_{a}^{x} p(z)dz} \frac{dy}{dx} + e^{\int_{a}^{x} p(z)dz} p(x)y = e^{\int_{a}^{x} p(z)dz} r(x)
\]  

Notice that the left-hand-side can be written as a derivative of a simple expression

\[
e^{\int_{a}^{x} p(z)dz} \frac{dy}{dx} + e^{\int_{a}^{x} p(z)dz} p(x)y = \frac{d}{dx} \left(e^{\int_{a}^{x} p(z)dz} y\right) \therefore \]

\[
\frac{d}{dx} \left(e^{\int_{a}^{x} p(z)dz} y\right) = e^{\int_{a}^{x} p(z)dz} r(x)
\]  

which can be integrated and then multiplied on both sides by \(e^{-\int_{a}^{x} p(z)dz}\):

\[
y(x) = e^{-\int_{a}^{x} p(z)dz} \left[ \int_{b}^{x} r(z) \left(e^{\int_{a}^{z} p(\eta)d\eta}\right) dz\right]
\]  

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**Example: The Bernoulli Equation**

The linear first-order ODEs always have a closed form solution in terms of integrals. In general non-linear ODEs do not have a general expression for their solution. However, there are some non-linear equations that can be reduced to a linear form; one such case is the Bernoulli equation:

\[
\frac{dy}{dx} + p(x) y = r(x) y^a
\]  

Reduction relies on a clever change-of-variable, let \(u(x) = [y(x)]^{1-a}\), then Eq. 20-21 becomes

\[
\frac{du}{dx} + (1 - a)p(x) u = (1 - a) r(x)
\]  

which is a linear heterogeneous first-order ODE and has a closed-form solution.

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\(^{10}\)The statistical definition of entropy is \(S(T, V) = k \log \Omega(U(T, V))\) or \(\Omega(U(T, V)) = \exp(S/k)\). Entropy plays the role of integrating factor.
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*Using NDSolve*